

SOLVAY WORKSHOP

"Beyond the Standard model with Neutrinos and Nuclear Physics"

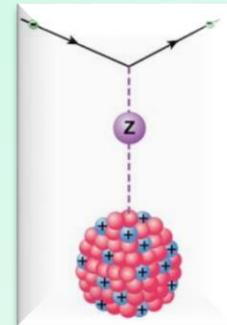
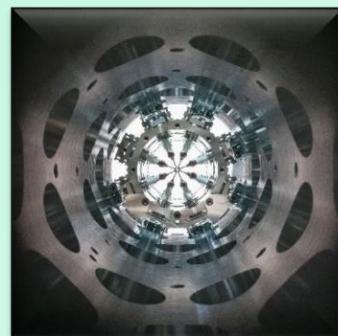
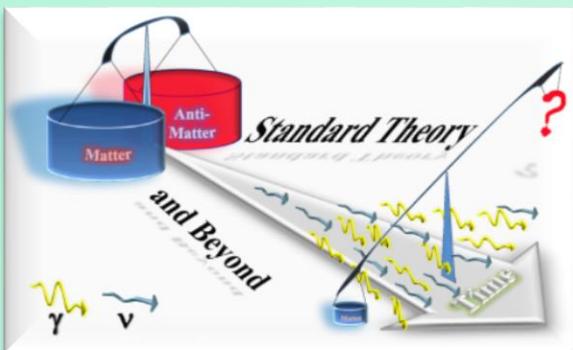
Brussels, November 29th - December 1st, 2017

Instable Particles as Probes for New Physics – Searches for APV and EDMs

Klaus Jungmann

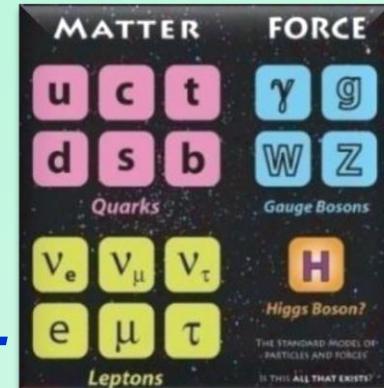
Van Swinderen Institute, University of Groningen

Searches for Electric Dipole Moments/Parity Violations



Standard Model Tests

- Standard Model (SM) of particle physics is *Best Theory we have*
- Still large number of open questions e.g. *particle masses, origin of parity violation,*



Direct:
Searches for New Particles



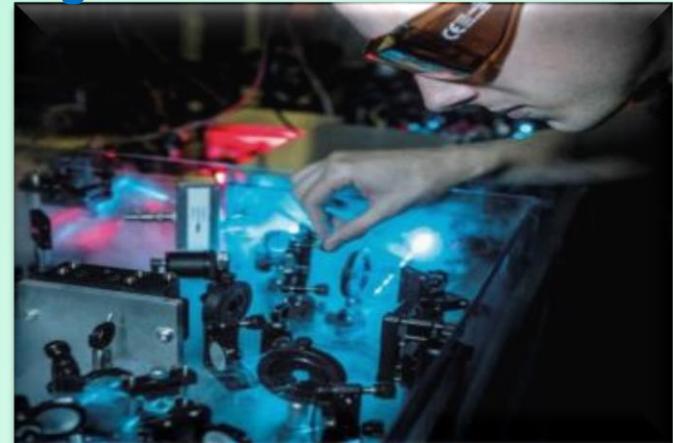
CERN e.g. LHC

e.g. Discovery of Higgs boson,...
also: Difference Matter-Antimatter ...



Equivalent
Approaches

Indirect:
High Precision Measurements

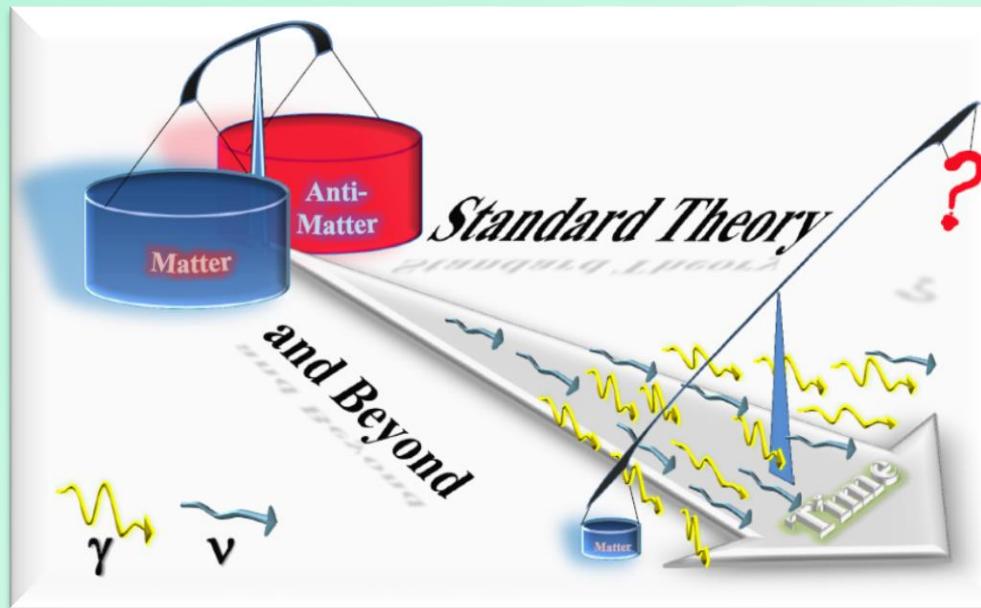


Small institutes e.g. VSI ..

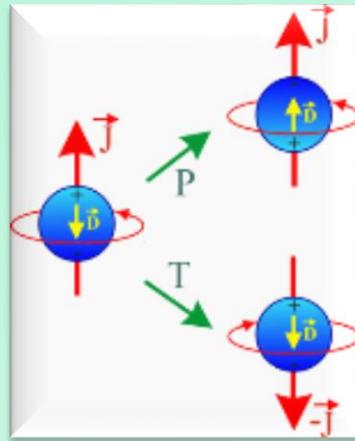
e.g. Atomic Parity Violation,
EDM searches,

Discrete Symmetries

C,P,T,CP,CPT



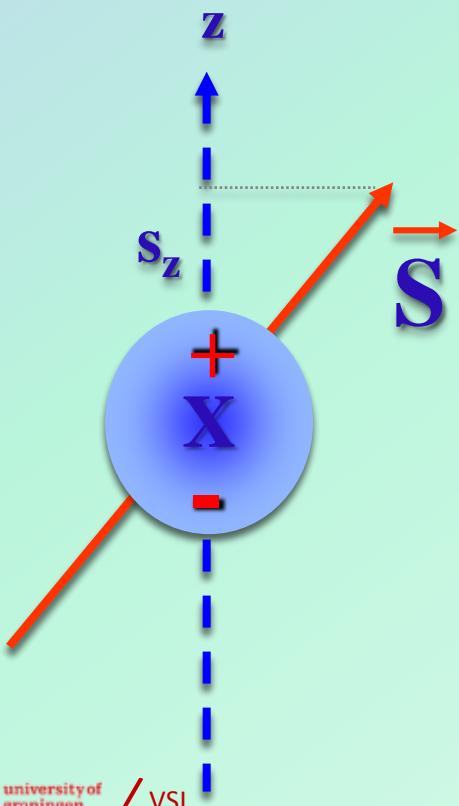
An EDM Violates P,T



and with CPT also CP

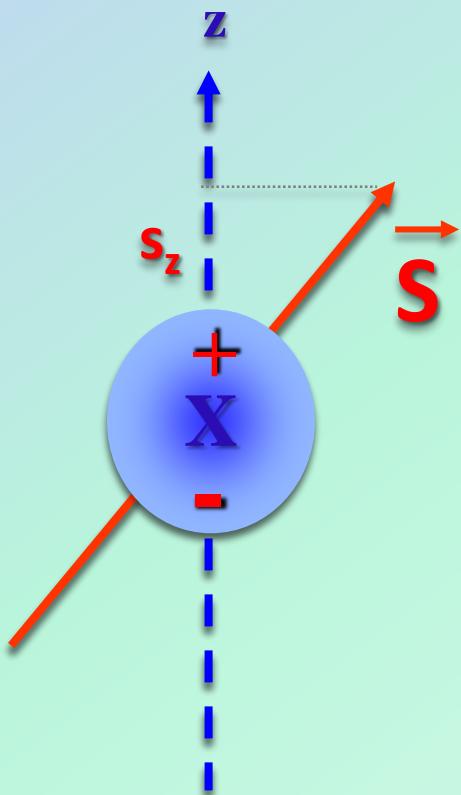
Permanent

Electric Dipole Moments



- Electron: clean and ready for New Physics
- Hadrons: depend on θ_{QCD} in Standard Model

Spin of Fundamental Particles



\vec{S} is the only vector characterizing a non-degenerate quantum state

magnetic moment:

$$\vec{\mu}_x = 2(1+a_x) \mu_{0x} c^{-1} \vec{S}$$

electric dipole moment:

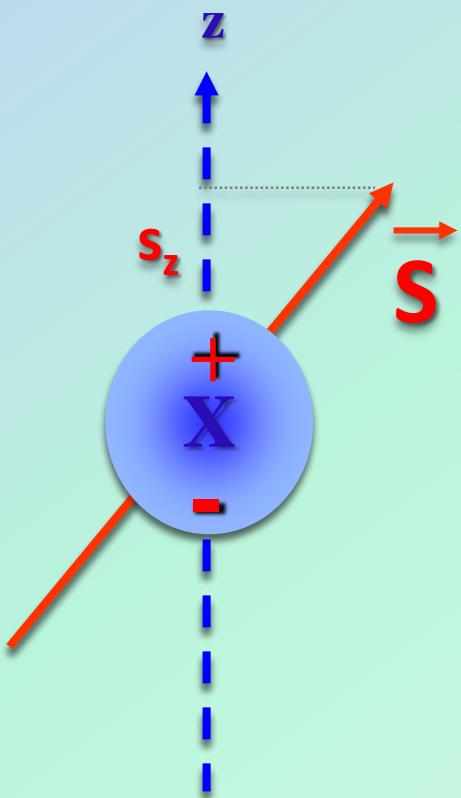
$$\vec{d}_x = \eta \mu_{0x} c^{-1} \vec{S}$$

magneton:

$$\mu_{0x} = e\hbar / (2m_x)$$

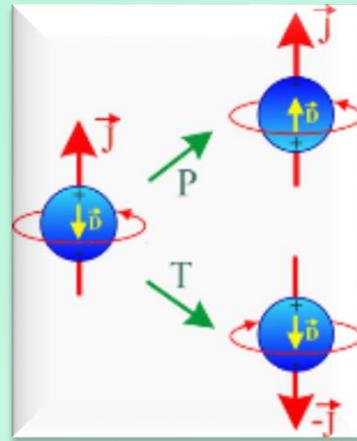
$$\mu_{0x} c^{-1} S = \begin{cases} 9.7 \cdot 10^{-12} \text{ e cm} & (\text{electron}) \\ 4.6 \cdot 10^{-14} \text{ e cm} & (\text{muon}) \\ 5.3 \cdot 10^{-15} \text{ e cm} & (\text{nucleon}) \end{cases}$$

Instable Particle EDMs



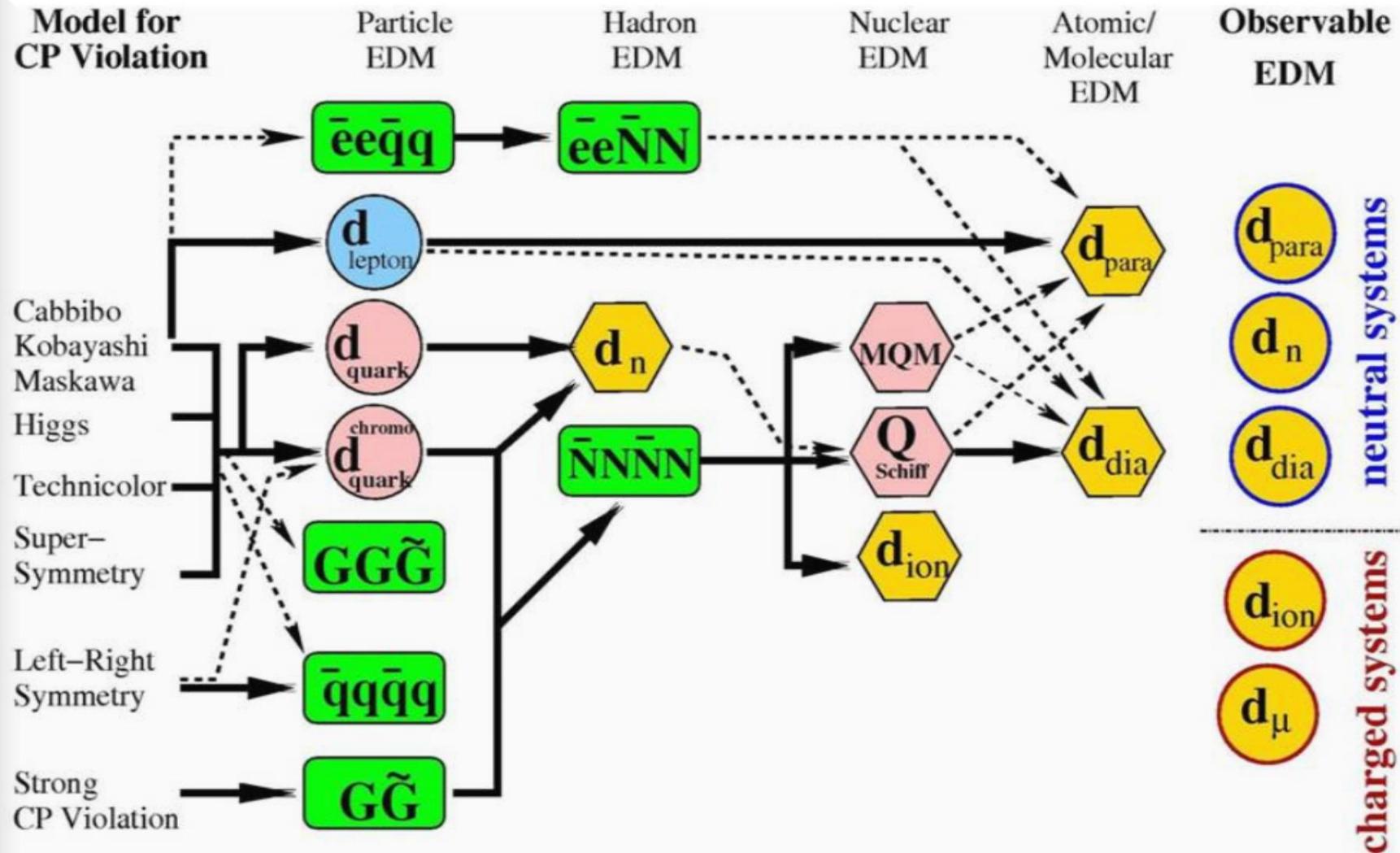
- In principle EDM not forbidden in instable states
→ e.g. transition dipoles exist
- Heavy (therefore instable) atoms have general advantage
 - deformed nuclei
 - Z^x enhancement (x typically 2...3)
- Instable particles may have detection advantage
 - β – asymmetry
 - are there oscillations in EDMs ? (axions)

An EDM Violates P,T



and with CPT also CP

Possible Sources of EDMs



The numerically best experiment until now- ^{199}Hg @Seattle –
Leaves somewhat restricted room for SUSY ...



Lines of attack towards an EDM

Free Particles

neutron
muon
deuteron
bare nuclei ?
...

- particle EDM
- unique information
- new insights
- new techniques
- challenging technology

Hg Xe
Tl ⁸⁷Rb
Cs Rn
Ra ...
Fr

Atoms

- electron EDM
- p...

Electric Dipole

* Since θ_{QCD} limited by hadronic EDMs
→ only leptons have direct transformative potential

spectroscopic
data

BaF , YbF
PbO , WC
PbF , ThO
HfF⁺, ThF⁺
RaF, ...

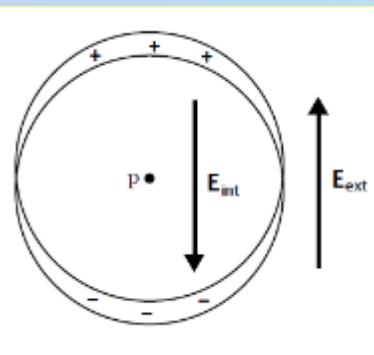
garnets
 $(\text{Gd}_3\text{Ga}_5\text{O}_{12})$
 $(\text{Gd}_3\text{Fe}_2\text{Fe}_3\text{O}_{12})$
solid He ?
liquid Xe

Condensed State

- electron EDM
- strong enhancements
- systematics ??



Enhancements of particle EDMs



$$\frac{d_{atom}}{d_e} \propto Z^3 \alpha^2 \chi$$

P. Sandars, 1968

$$d_{atom} = \sum_{n'} \frac{\langle n, l | -d_e(\beta - 1) \vec{\sigma} \cdot \vec{E} | n', l \pm 1 \rangle \langle n', l \pm 1 | -e \vec{r} | n, l \rangle}{E_{n,l} - E_{n',l\pm 1}} + h.c.$$

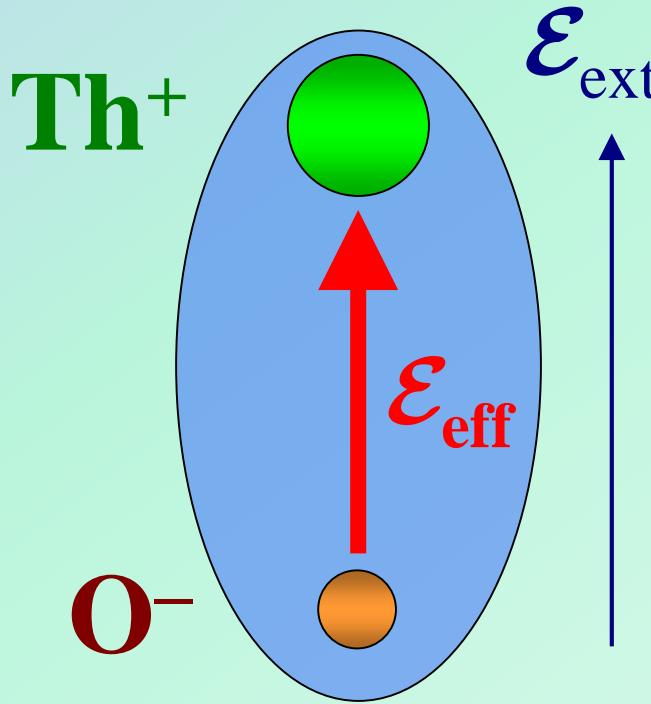
⇒ go for heavy systems, where $Z \gg 1$, e.g. Hg, Xe

⇒ take advantage of enhancements, e.g. Ra, Rn

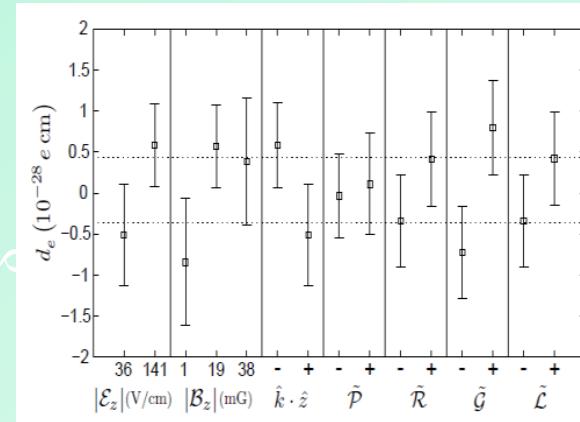
⇒ consider molecules such as YbF, RaF, ...



Highlight: ThO electron EDM experiment



$E_{\text{ext}} \sim 1 \text{ V/cm}$ enough for ThO



New limit for e-

$d_e < 8.7 * 10^{-29} \text{ ecm}$
(90% c.l.)

Doyle, Gabrielse , DeMille

Experiment presently taking further data

Atomic/Molecular Enhancement Factors

for Electron EDM

Particle	Rb	Cs	Tl	Fr	Ra
Enhancement	24	125	585	1 150	40 000

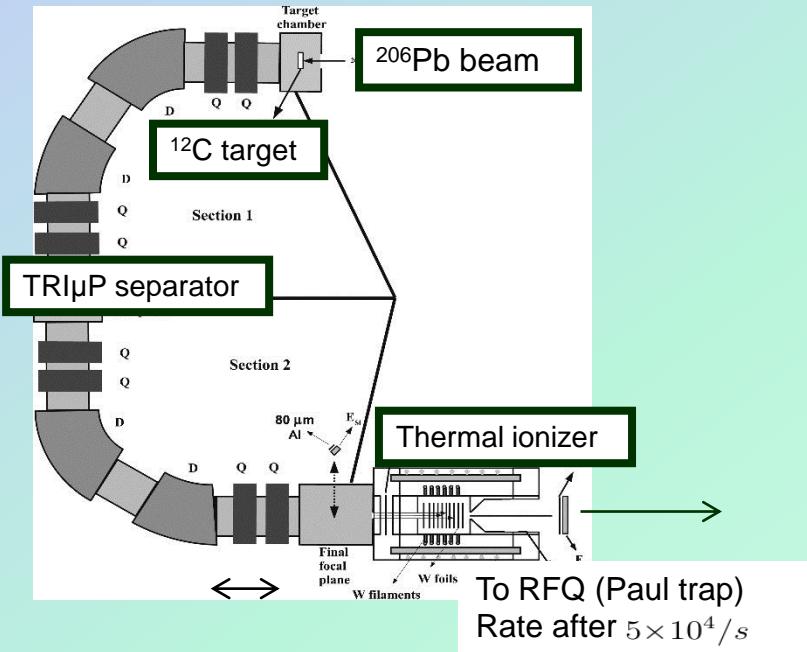
Flambaum, Dzuba, 2012

Particle	ThO	BaF	YbF	PbO
Enhancement	10^9	5×10^5	1.6×10^6	6×10^4

watch out:
Saturation

→ different theorists agree, typically at 30% level

Radium Isotopes



${}^{225}\text{Ra}$

extraction from ${}^{229}\text{Th}$ source ([ANL](#))

Long lived ${}^{229}\text{Th}$ source in an oven ([VSI](#))

Other Isotopes

Online production at accelerator facilities

e.g.

[TRIMuP@KVI](#) (flux $\sim 10^5/\text{s}$) (until 2013)

[ISOLDE, CERN](#) (flux $\sim 10^9/\text{s}$)

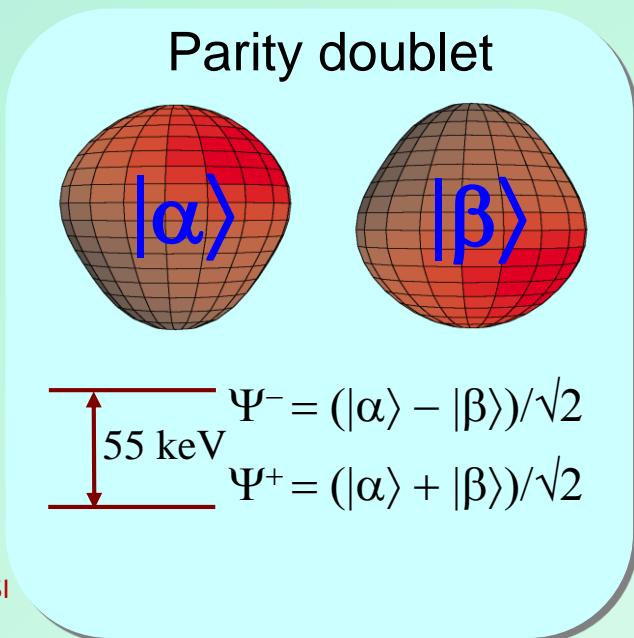
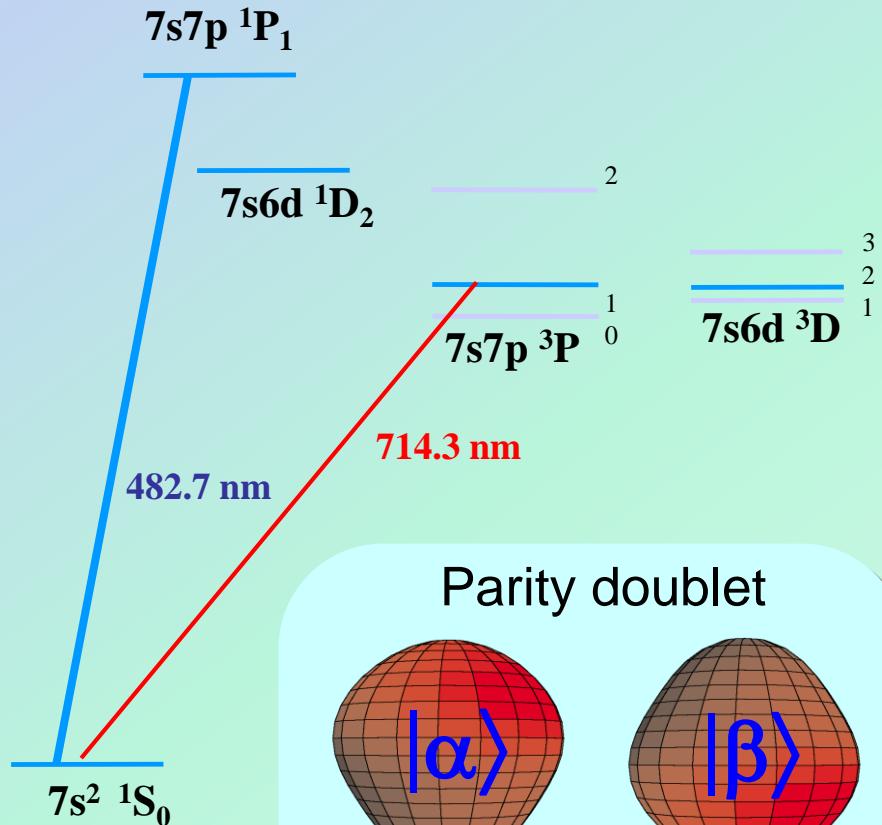
Sources or fragmentation

	Lifetime	Spin
209	4.6(2) s	5/2
211	13(2) s	1/2
212	13.0(2) s	
213	2.74(6) m	1/2
214	2.46(3) s	
221	28.2 s	5/2
223	11.43(5) d	3/2
224	3.6319(23) d	
225	14.9(2) d	1/2
226	1600 y	
227	42.2(5) m	3/2
229	4.0(2) m	5/2

ΔN > 14

Radium EDMs

Atomic energy level diagram of Ra



- Nearly degenerate opposite parity 3P_1 and 3D_2 , enhancement $>5000 \text{ e } \text{EDM}$

$$d = \frac{\langle {}^3D_1 | -er | {}^3P_1 \rangle \langle {}^3P_1 | H_{\text{EDM}} | {}^3D_1 \rangle}{E({}^3D_1) - E({}^3P_1)}$$

V. A. Dzuba et al. Phys. Rev. A, 61, 062509 (2000)

Density distribution of nuclear charge has mixed octupole and quadrupole deformation

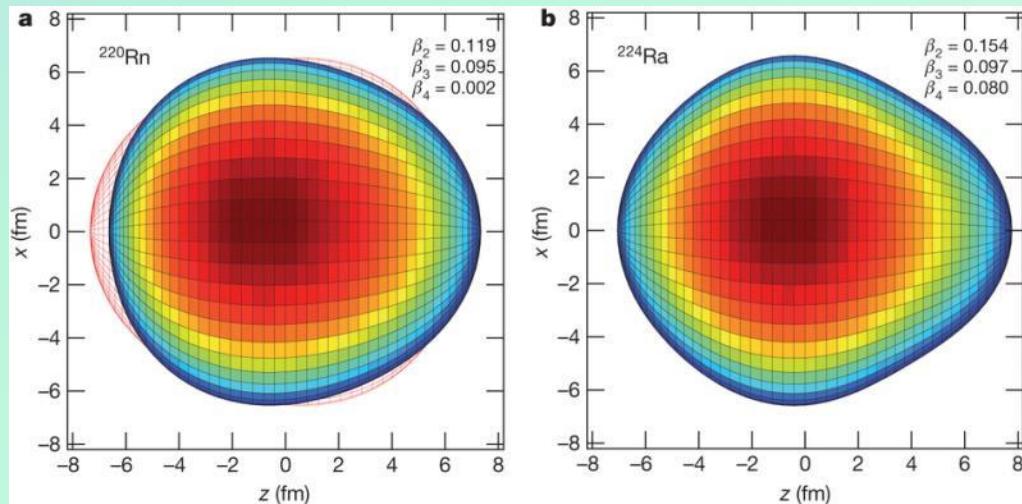
- Deformed charge distribution in some isotopes (${}^{225}\text{Ra}$)
- Nucleon EDM enhances $\approx 10^2$

Dobczewski, Engel, PRL (2005) & Phys. Rev. C (2010)

EDM Enhancement by Nuclear Deformation



L. P. Gaffney et al, **Nature** 497, 157 (2013)

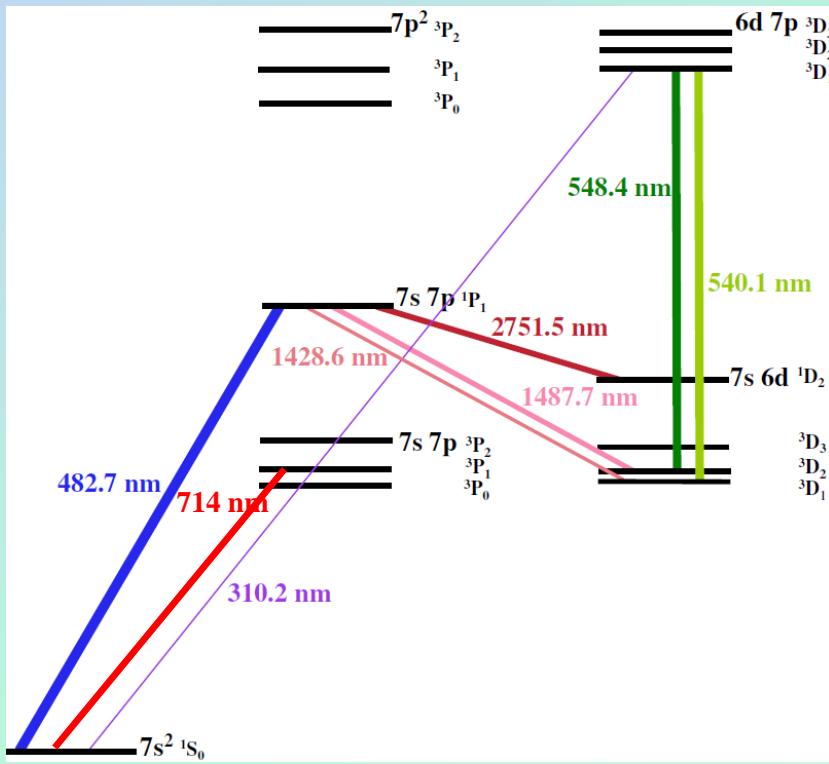


- If that were true also for odd spin isotopes there'd be a nucleon EDM enhancement by factor of some 200
- Need measurements for odd isotopes now !!

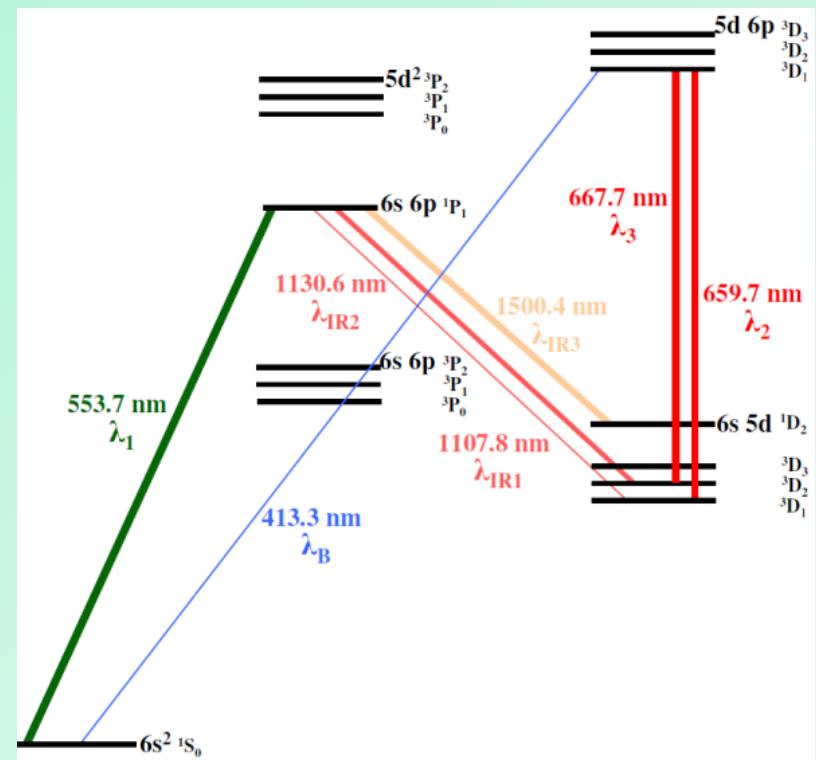
(see e.g. Y.K. Khriplovich)

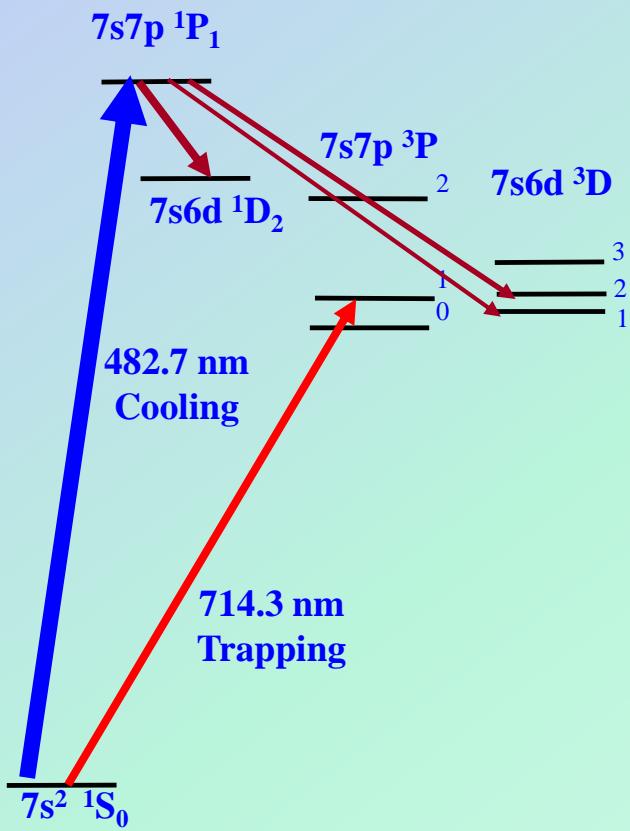
Radium and Barium

Radium



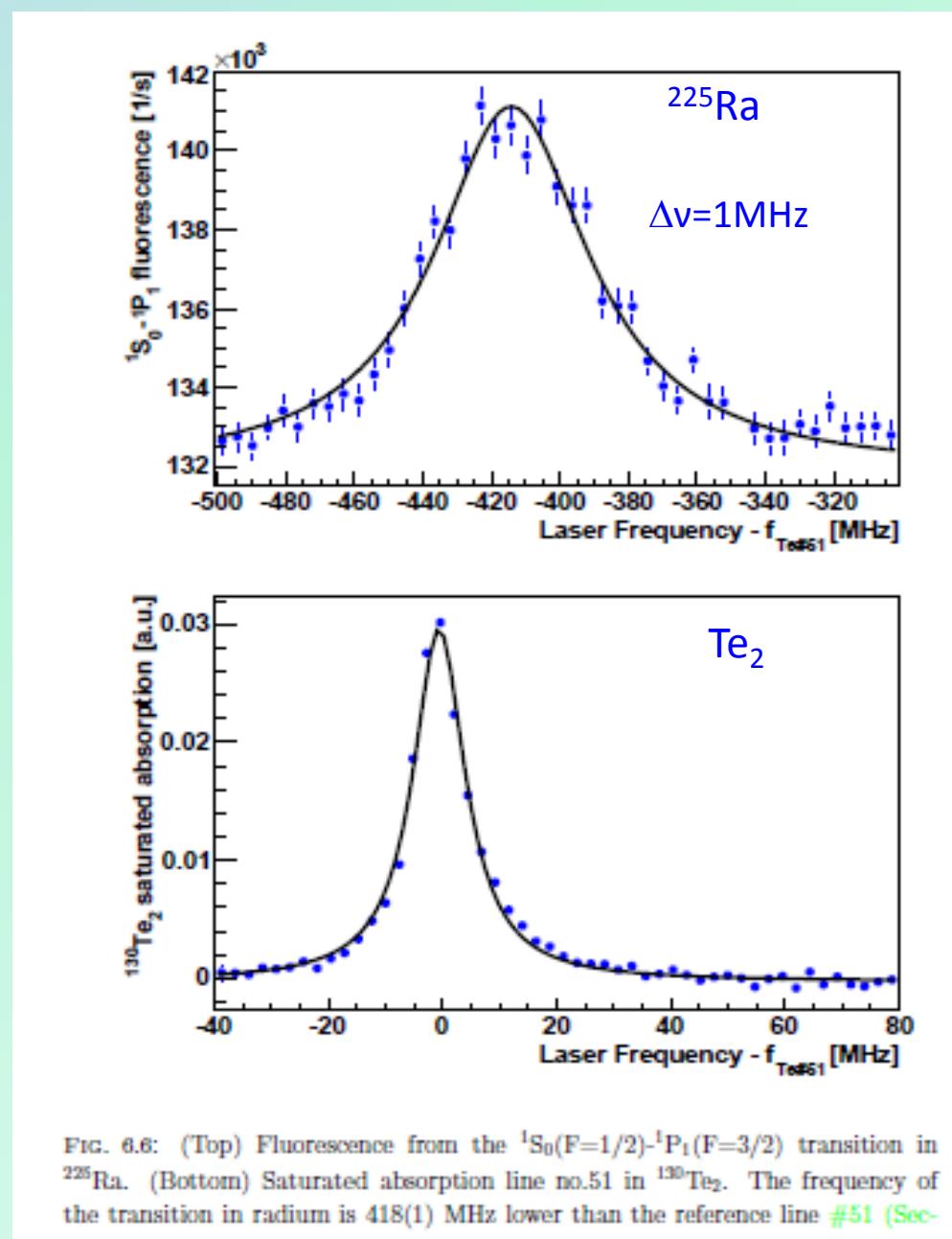
Barium





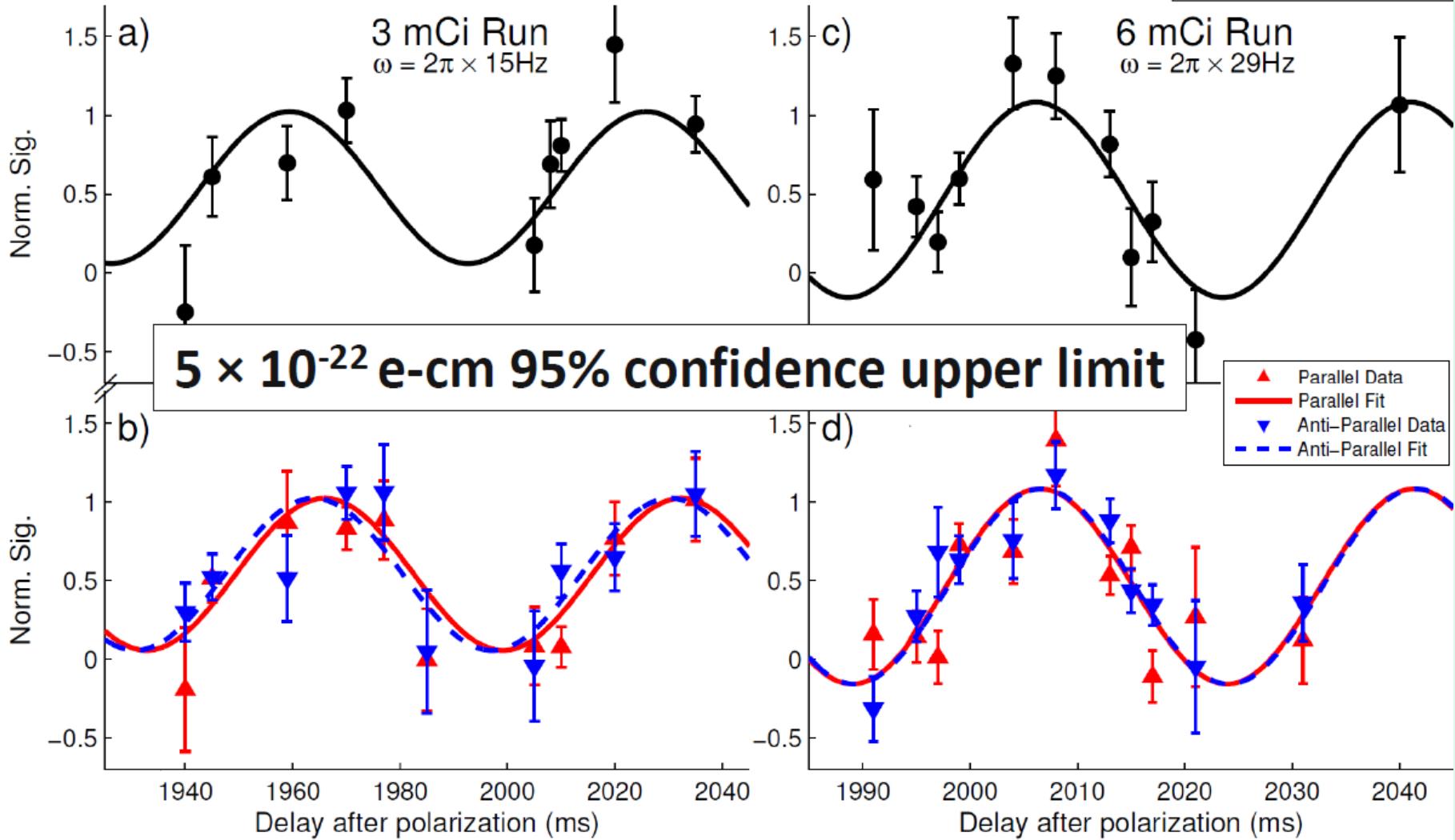
Isotope	Transition	Frequency [MHz]	Experiment by
^{226}Ra	${}^1\text{S}_0 - {}^1\text{P}_1$	621038489 (15)	This work
^{226}Ra	${}^1\text{S}_0 - {}^1\text{P}_1$	621038004 (180)	Trimble et al.
^{226}Ra	${}^1\text{S}_0 - {}^1\text{P}_1$	621041362 (1500)	Rasmussen

B. Santra et al, PRA (R) (2014)



Argonne ^{225}Ra Experiment

→ Near term goal $4 \times 10^{-25} \text{ ecm}$



Towards a Rn EDM Experiment at TRIUMF

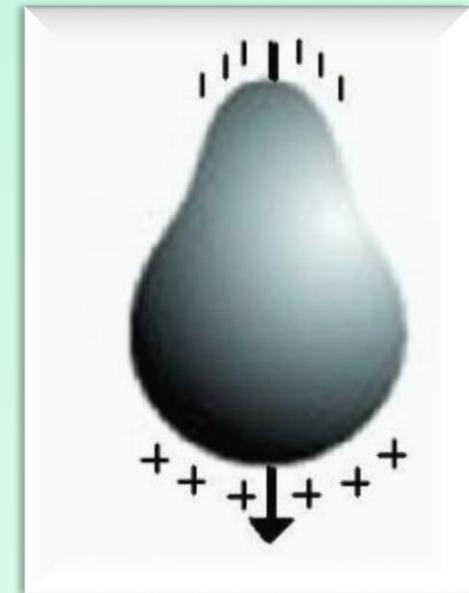
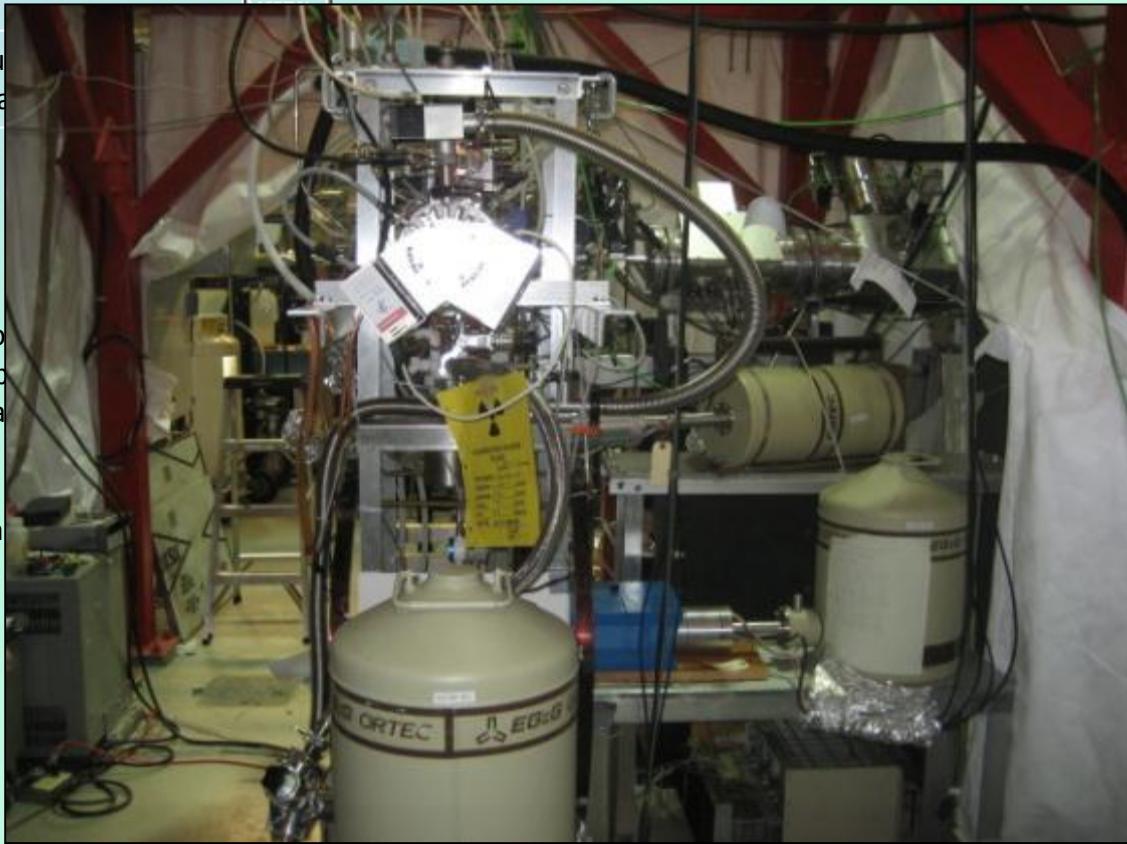
T. Chupp and C. Svensson

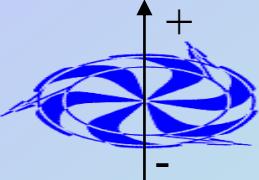
- Magnitude of EDM $\sim Z^3$
- Radon isotopes possibly octupole deformed
- Rn is predicted to be ~ 600 times more sensitive than ^{199}Hg

6. Pu
to tra

3. Co
temp
the a

4. Wa



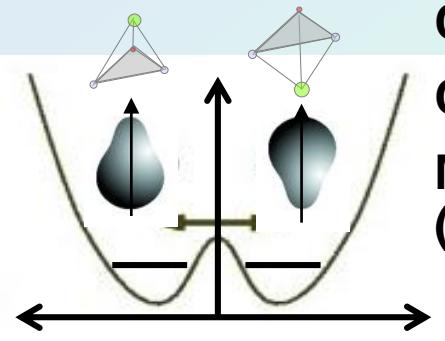


Radon-EDM Experiment

TRIUMF E929

T. Chupp (Michigan) & C. Svensson (Guelph)
Funding: NSF, DOE, NRC, NSERC

TRIUMF



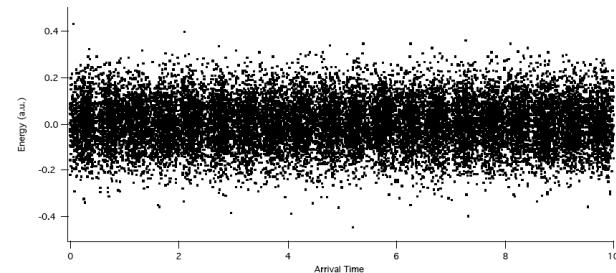
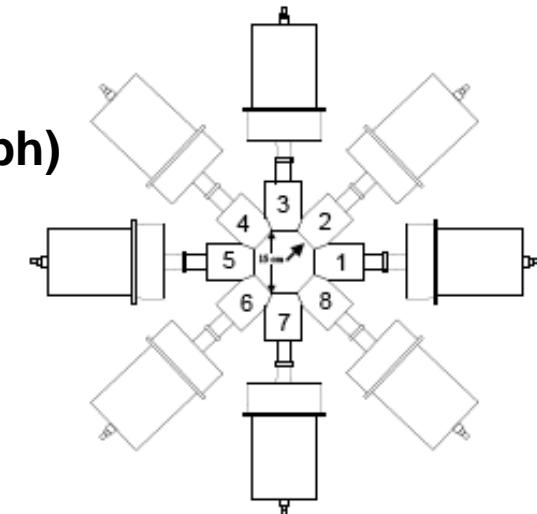
Produce rare ion radon beam

Collect in cell

Comagnetometer

Measure free precession
(γ anisotropy/ β asymmetry)

$$\sigma_d \approx \frac{\hbar}{AET_2\sqrt{N}}$$

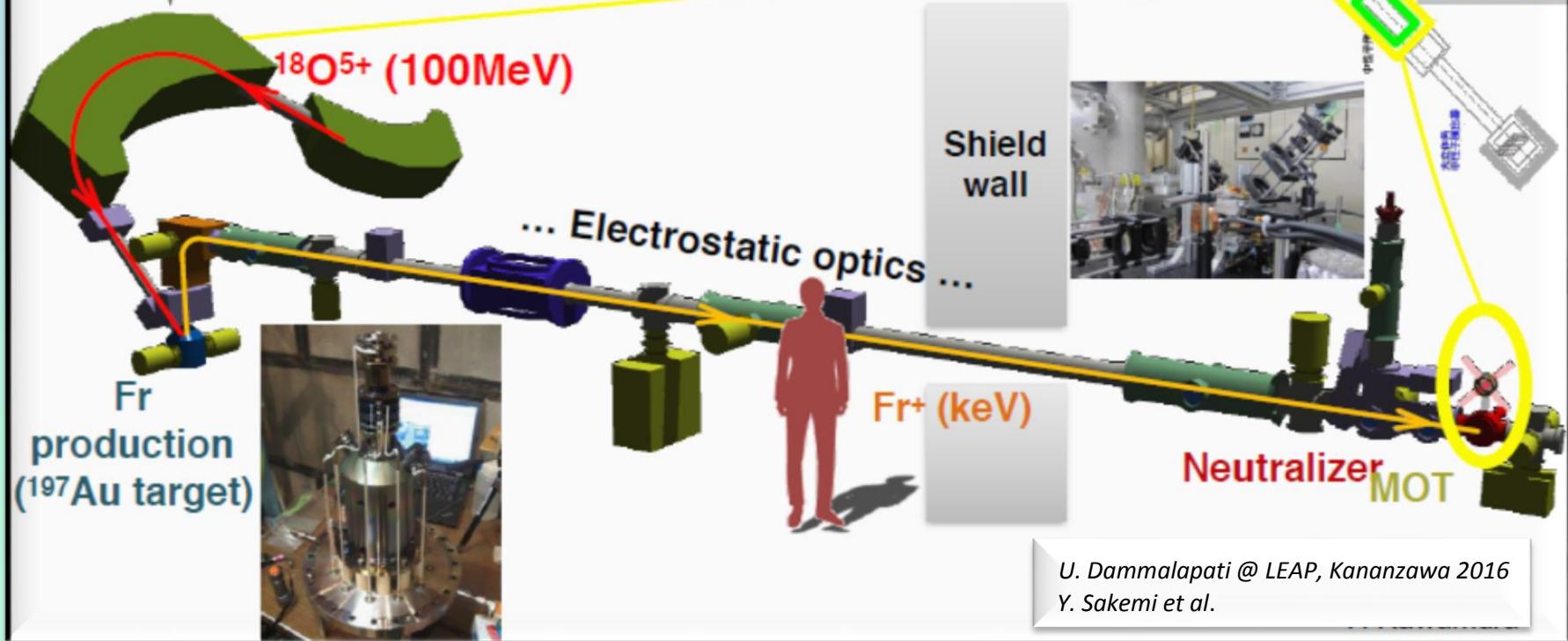
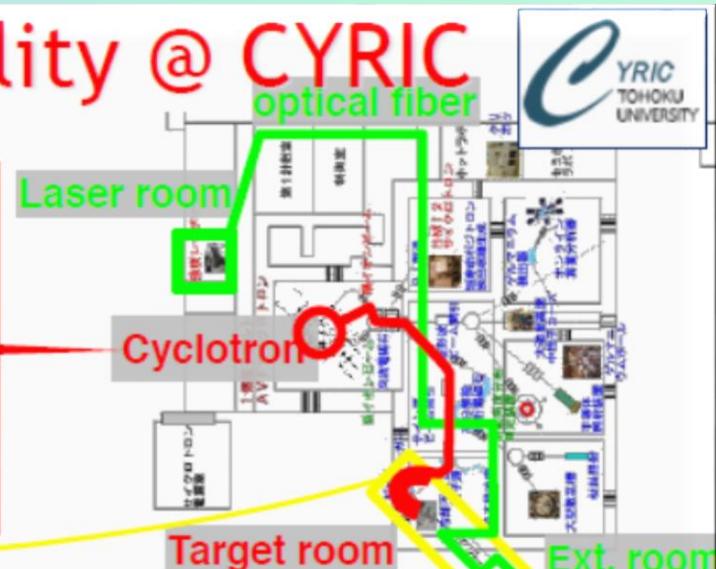
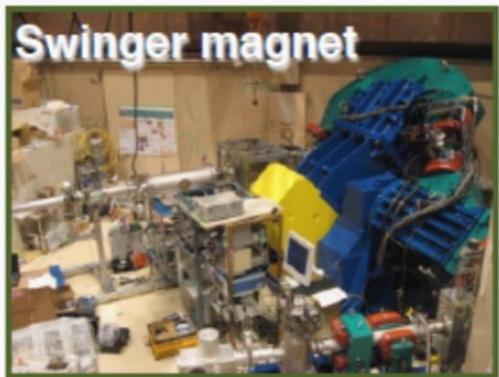


221/223Rn EDM projected sensitivity

Facility	Detection	S_d (100 d)
ISAC	γ anisotropy	2×10^{-26} e-cm
ISAC	β asymmetry	1×10^{-27} e-cm
FRIB	β asymmetry	2×10^{-28} e-cm

→ ~ 5×10^{-30} for ^{199}Hg

Francium project and facility @ CYRIC



U. Dammalapati @ LEAP, Kanazawa 2016
Y. Sakemi et al.

Experiment is on the move to U Tokyo / RIKEN

Generic EDM Figure of Merit

figure of merit

$$M = EP\varepsilon\sqrt{\tau TN} \text{ *enh}$$

electric field

polarization

efficiency

coherence time

enhancement
factor

particles
in experiment

total measurement time

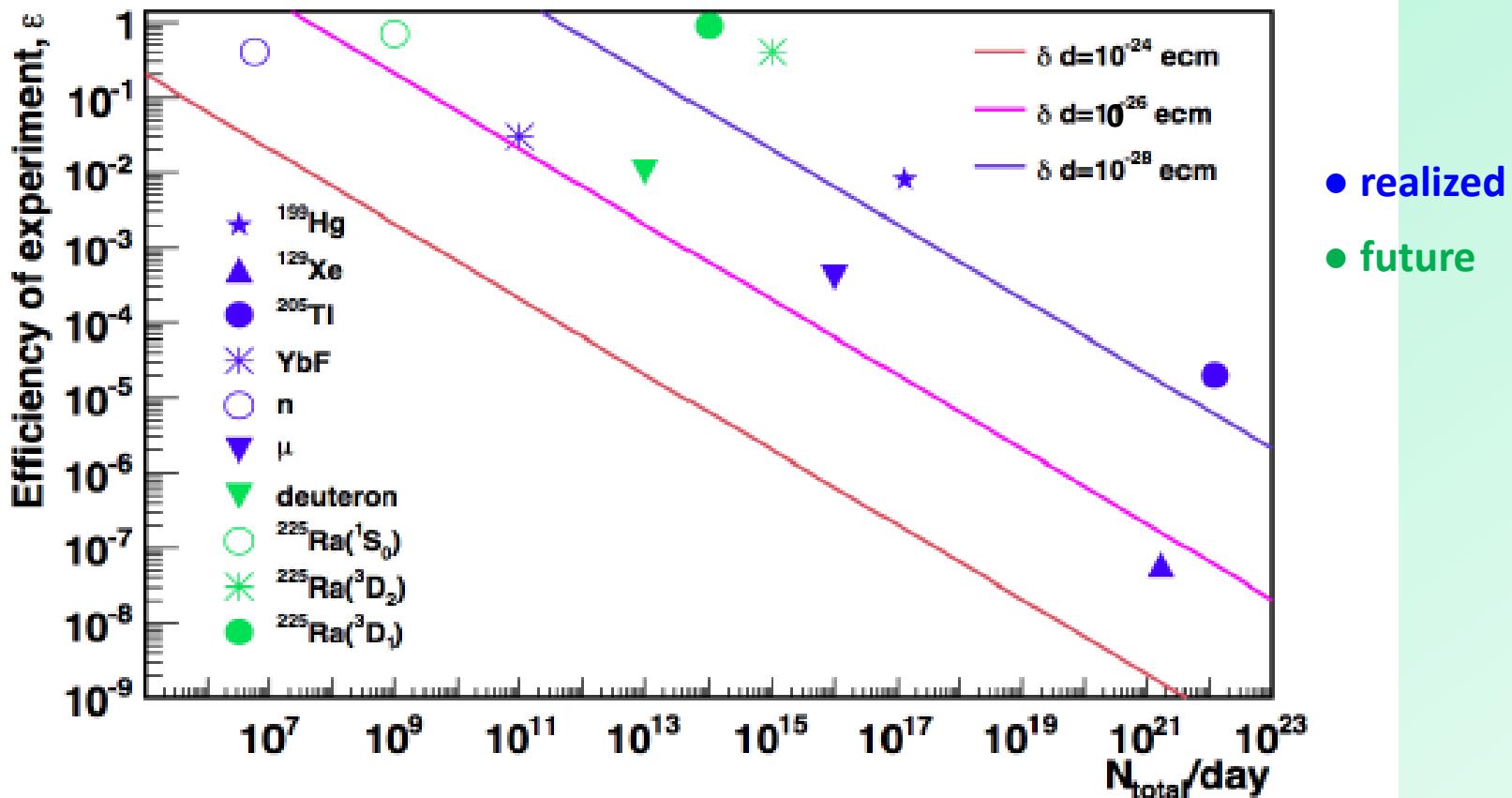
Preferred Systems

$$\delta d = \frac{\hbar}{EP\varepsilon\sqrt{\tau TN}} / \text{enh}$$

T measurement time
 P polarization
 enh enhancement

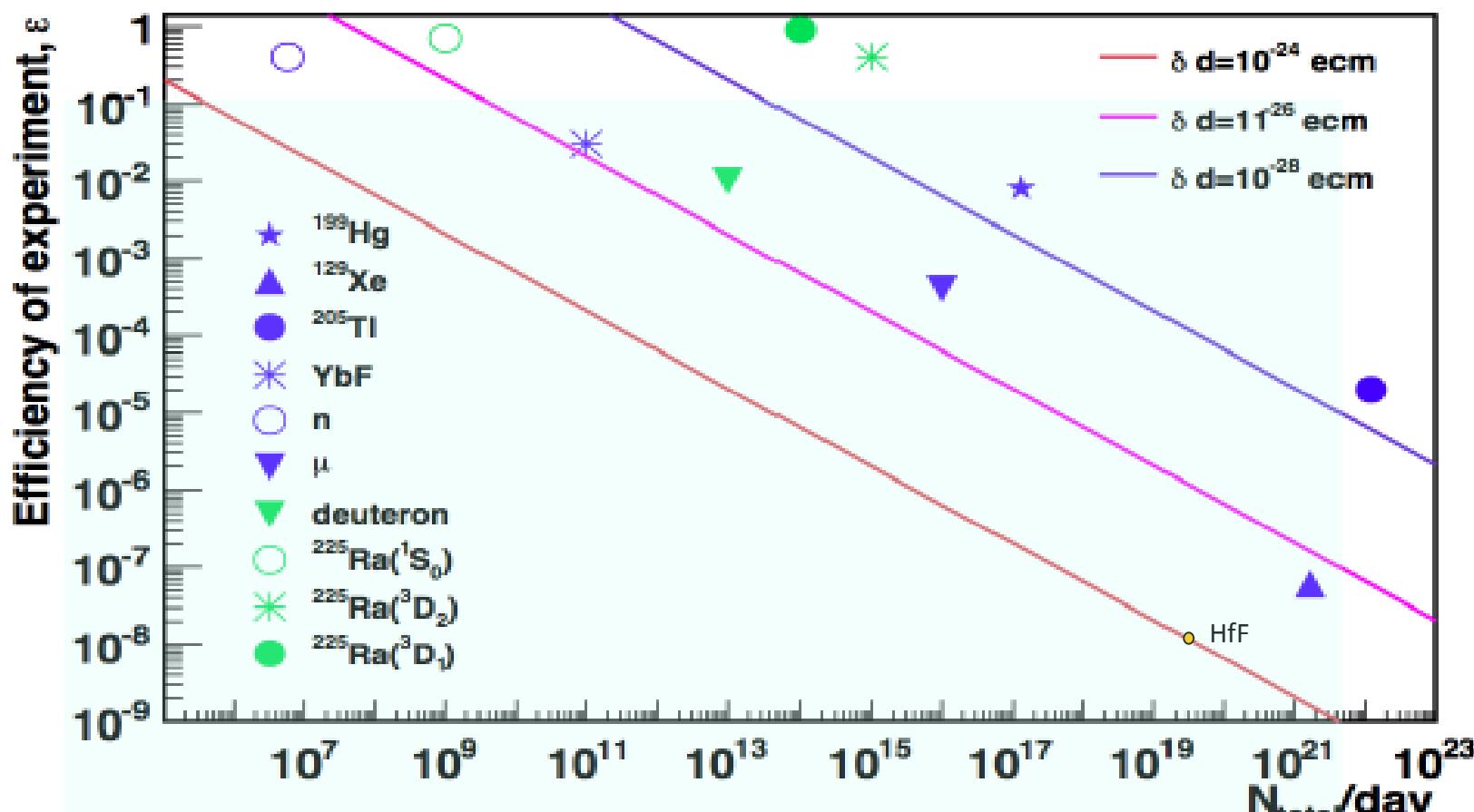
Particle	Number Particles N	Coherence Time τ [s]	Efficiency ε	Electric Field E [kV/cm]	Figure of Merrit
^{199}Hg	10^{14}	2×10^{-2}	8×10^{-3}	10	5×10^{13}
^{129}Xe	10^{22}	10^{-4}	9×10^{-9}	3.6	1×10^{14}
ThO	10^{11}	1.1×10^{-3}	2×10^{-2}	<0.1	2×10^{13}
YbF	10^5	1.5×10^{-3}	3×10^{-2}	10	1×10^{12}
BaF	10^{11}	10^{-1}	10^{-2}	10	5×10^{13}
^{225}Ra	10^3	4×10^1	7×10^{-5}	67	3×10^6

EDM Experiments: Efficiency



$$\delta d = \frac{\hbar |\vec{l}| / \text{enh}}{EP\epsilon\tau \sqrt{N_{\text{total}}}}$$

B. Santra, L. Willmann (2013)

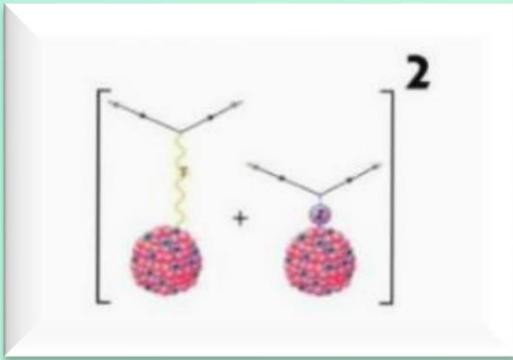


$$\delta d = \frac{\hbar|| / \text{enh}}{EP\epsilon\tau\sqrt{N_{\text{total}}}}$$

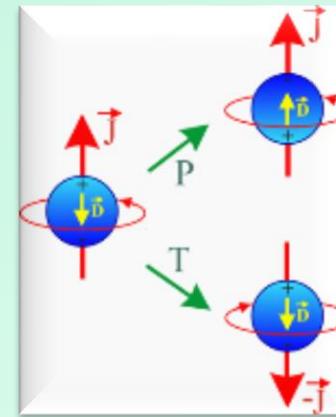
Cold Molecules

for

EDMs & Parity



SrF

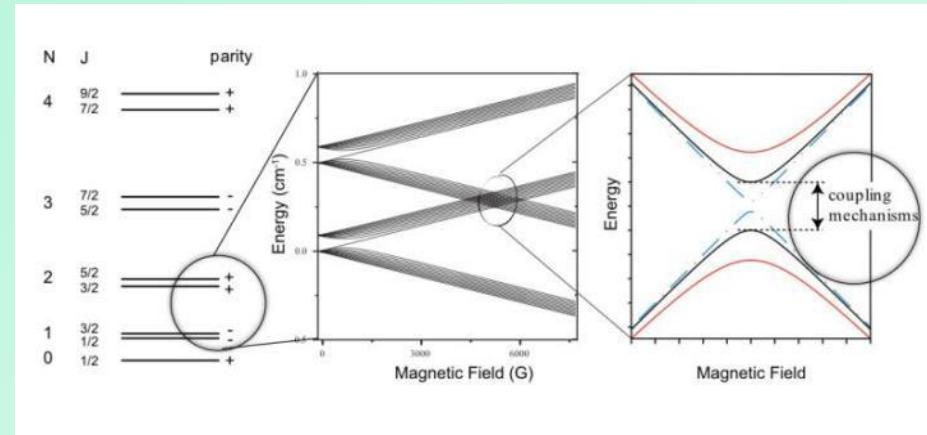


BaF

RaF

Precision Measurements with Molecules

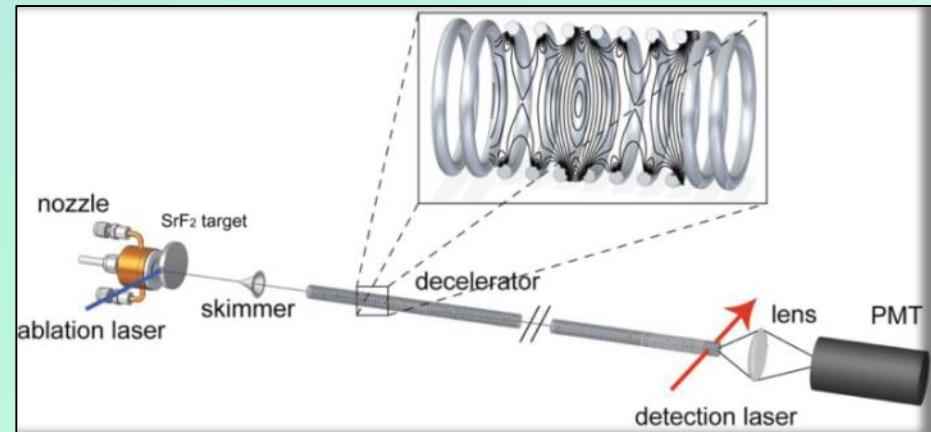
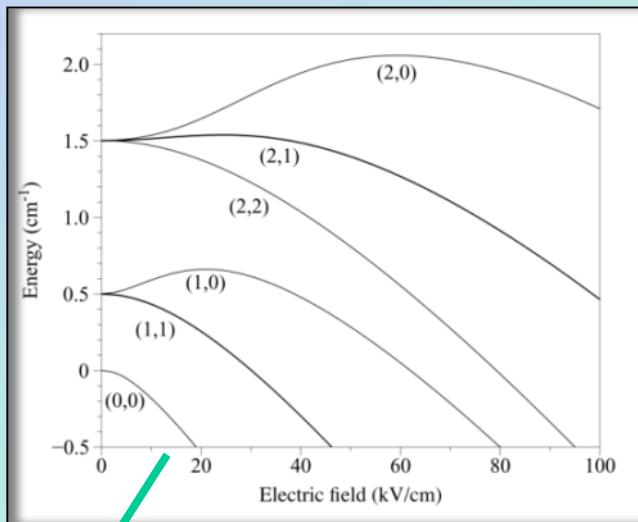
- Heavy diatomic molecules (SrF , RaF ,..) are suited for precision measurements (parity violation, eEDM)
- Large enhancement due to almost degenerate rotational levels



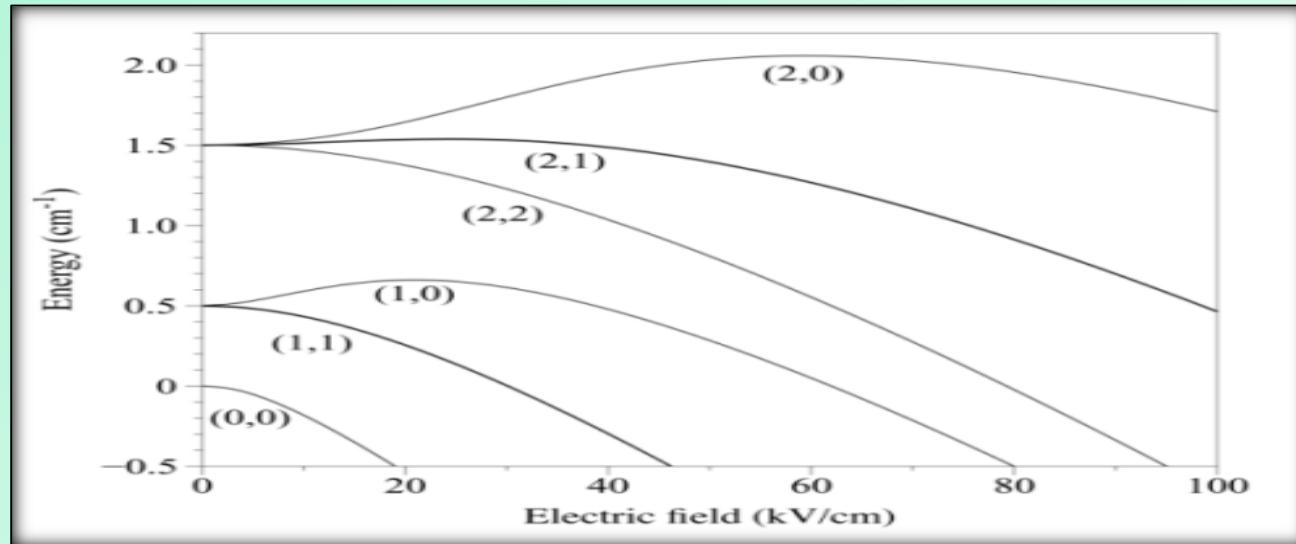
- Ultracold molecules by a traveling wave decelerator and laser cooling
- Benefit from the long interaction time provided by a cold, trapped sample

C. Meinema, J. v/d Berg, S. Hoekstra

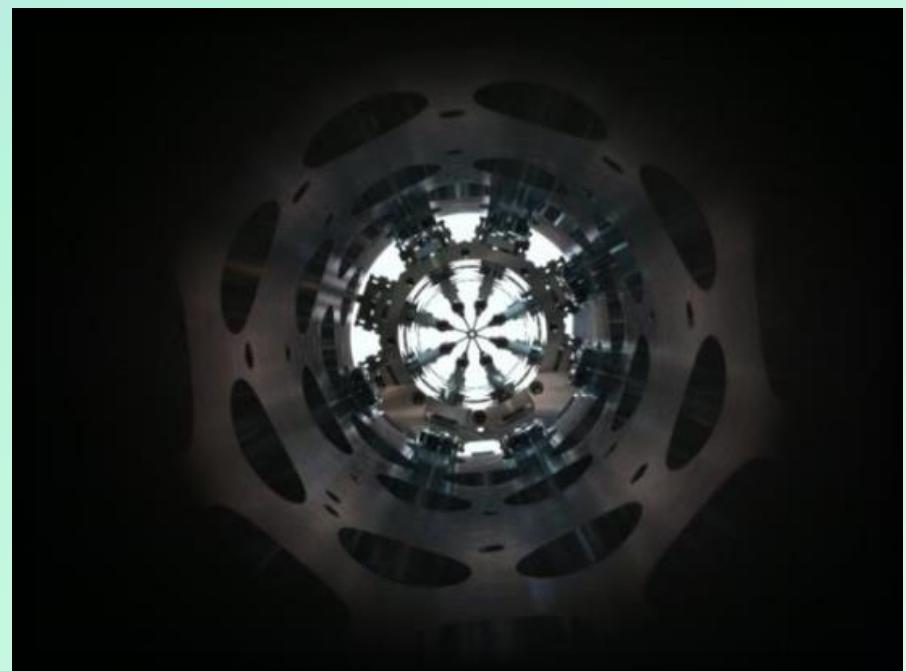
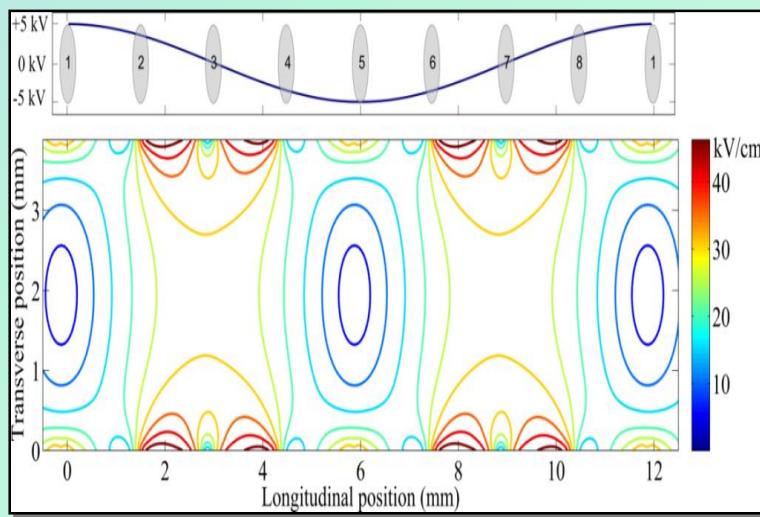
Traveling wave decelerator



SrF



Traveling wave decelerator



**5 m of decelerator
10 modules of 50 cm
3360 ring electrodes
diameter electrode: 4 mm**

C. Meinema, J. v/d Berg, S. Hoekstra

SrF Slowed Down

Signal and Simulations

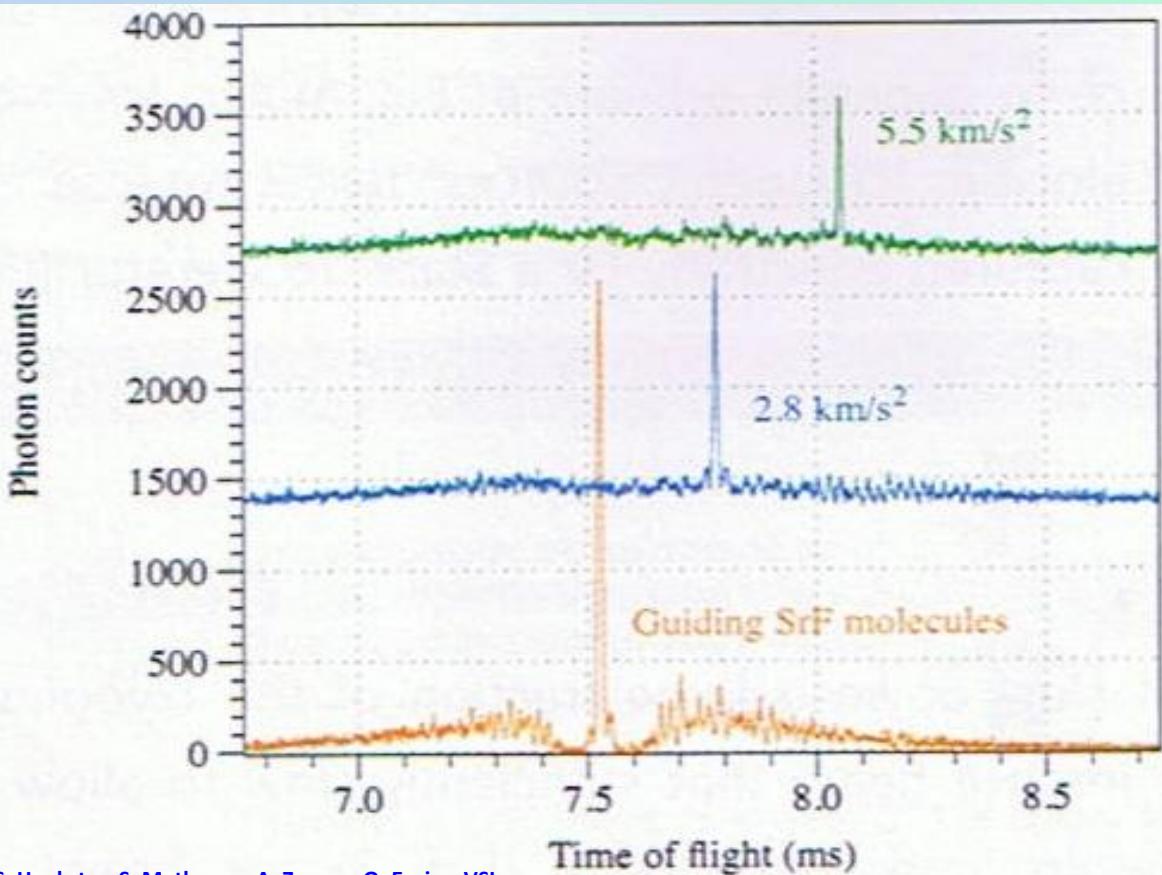
- 4 of 8 amplifiers
- 2 m machine



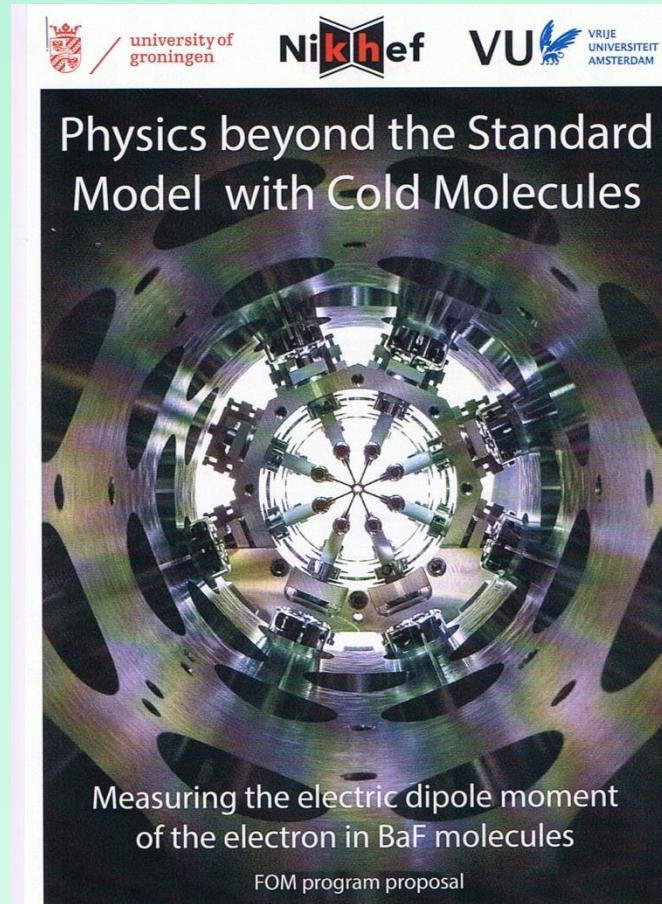
C. Meinema, J. v/d Berg, S. Hoekstra

SrF Slowed Down and Guided

- 8 of 8 amplifiers
- 4 m machine



The way to go for eEDM below 10^{-29} ecm

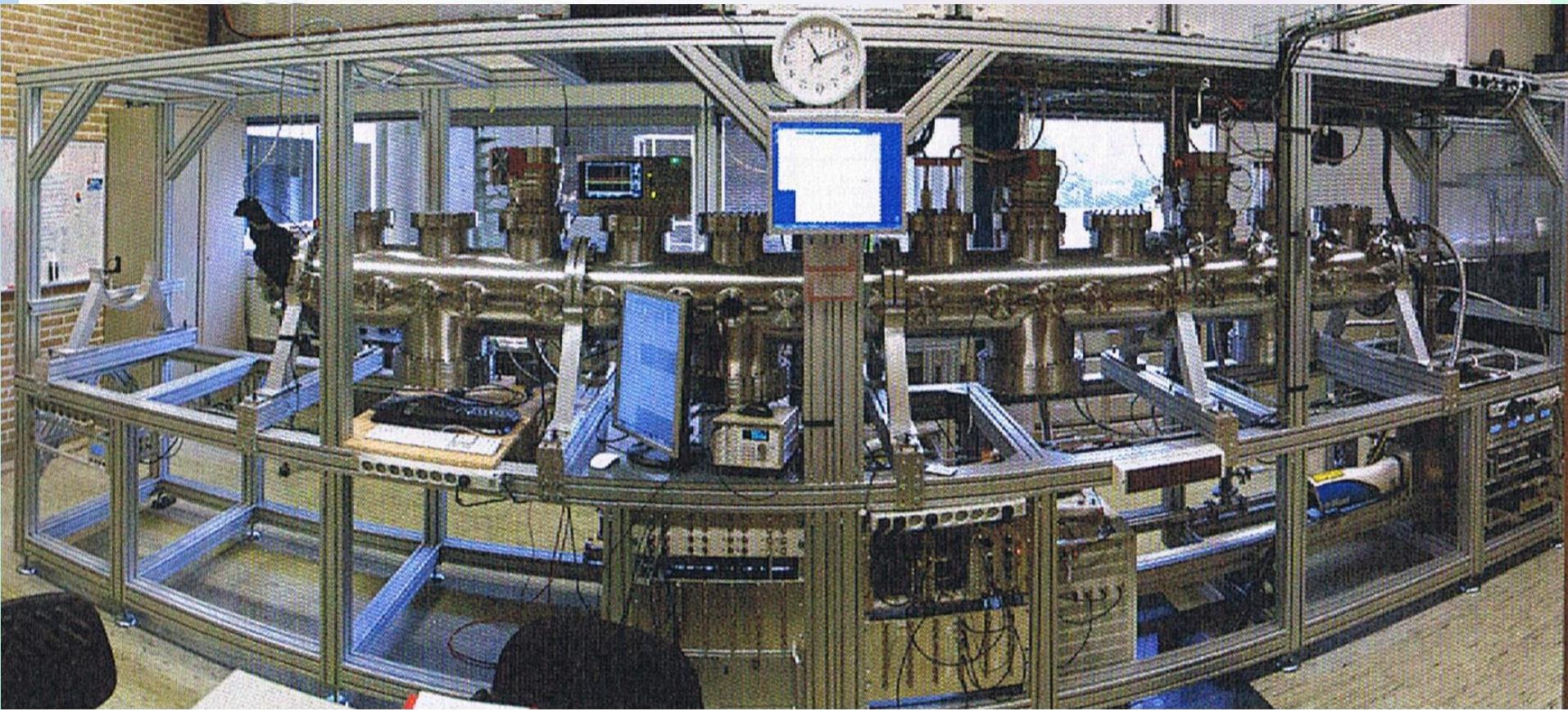


S. Hoekstra et al.

BaF eEDM

machine in statu nascendi

S.Hoekstra, A. Borschevsky, K. Jungmann, R.G.E. Timmermans, L. Willmann,
H. Bethlem, W. Ubachs et al. (FOM/NWO programme 2016-2022)

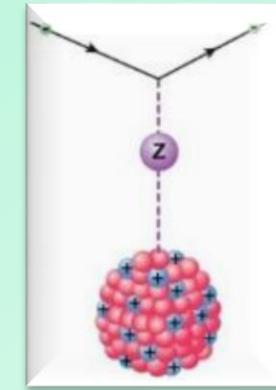
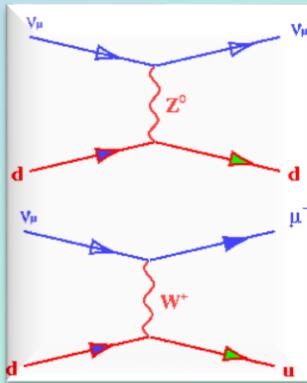


→ eEDM Collaboration Goal: Best EDM Limit on Electron

Time of flight (ms)

S. Hoekstra, S. Mathavan, A. Zapara, Q. Esajas, VSI

Parity



- *relatively large effects* in some atoms and molecules scaling with Z^3 or even stronger
- one valence electron atoms to extract precise constants
- more complex systems to study e.g. anapole moments

Atomic Parity Violation (APV)

Physics beyond the SM

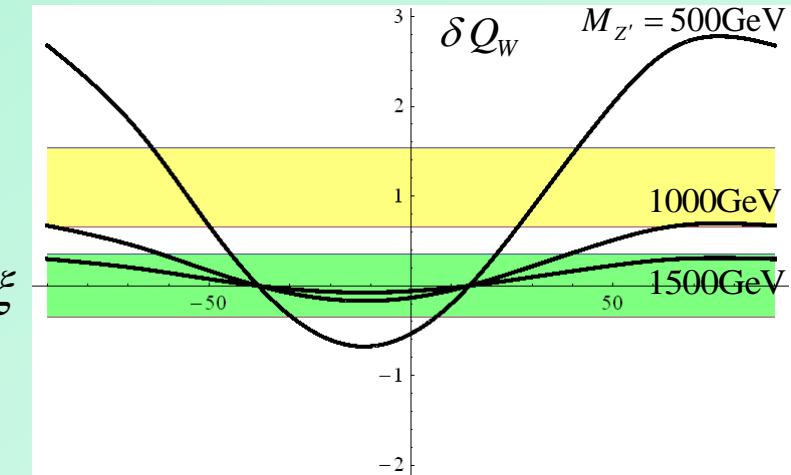
$$Q_W = -N + (1 - 4 \sin^2 \theta_W) Z + \text{rad. corr.} + \text{"new physics"}$$

Extra Z' boson in SO(10) GUTs:

- Additional U(1)' gauge symmetry
- Known Z and W unaffected
- No Z-Z' mixing

$$\delta Q_W \cong (2N + Z) a_e'(\xi) v_d'(\xi) \left[\frac{M_Z^2}{M_{Z'}^2} \right]$$

London en Rosner (1986), Marciano en Rosner (1990), Altarelli et al. (1991)



Bound on $M_{Z'}$ from cesium APV

(68% confidence level, $\xi = 52^\circ$) Wansbeek et al., PRA, (2010)

$$M_{Z'} > 1.2 \text{ TeV}/c^2$$

(Tevatron $M_{Z'} > 0.9 \text{ TeV}/c^2$)

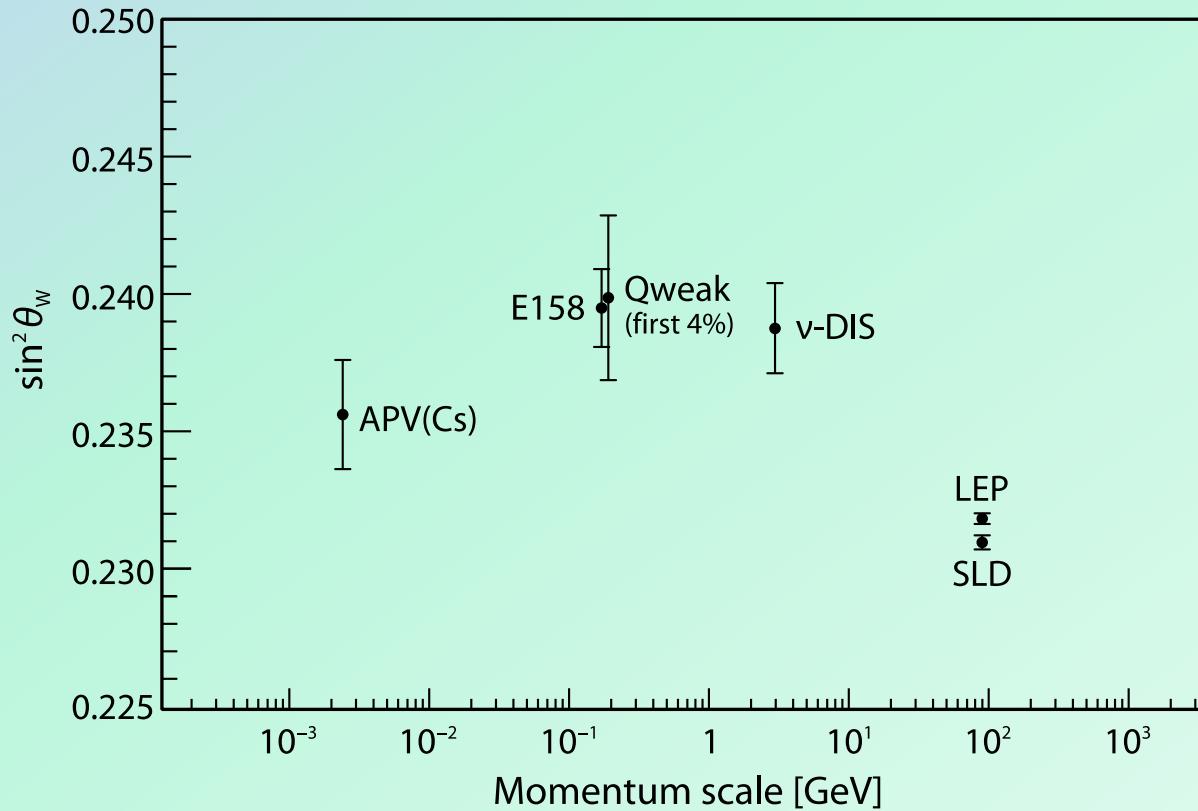
Bound (possible) on $M_{Z'}$ from Ra⁺ APV

$$M_{Z'} > 6 \text{ TeV}/c^2$$

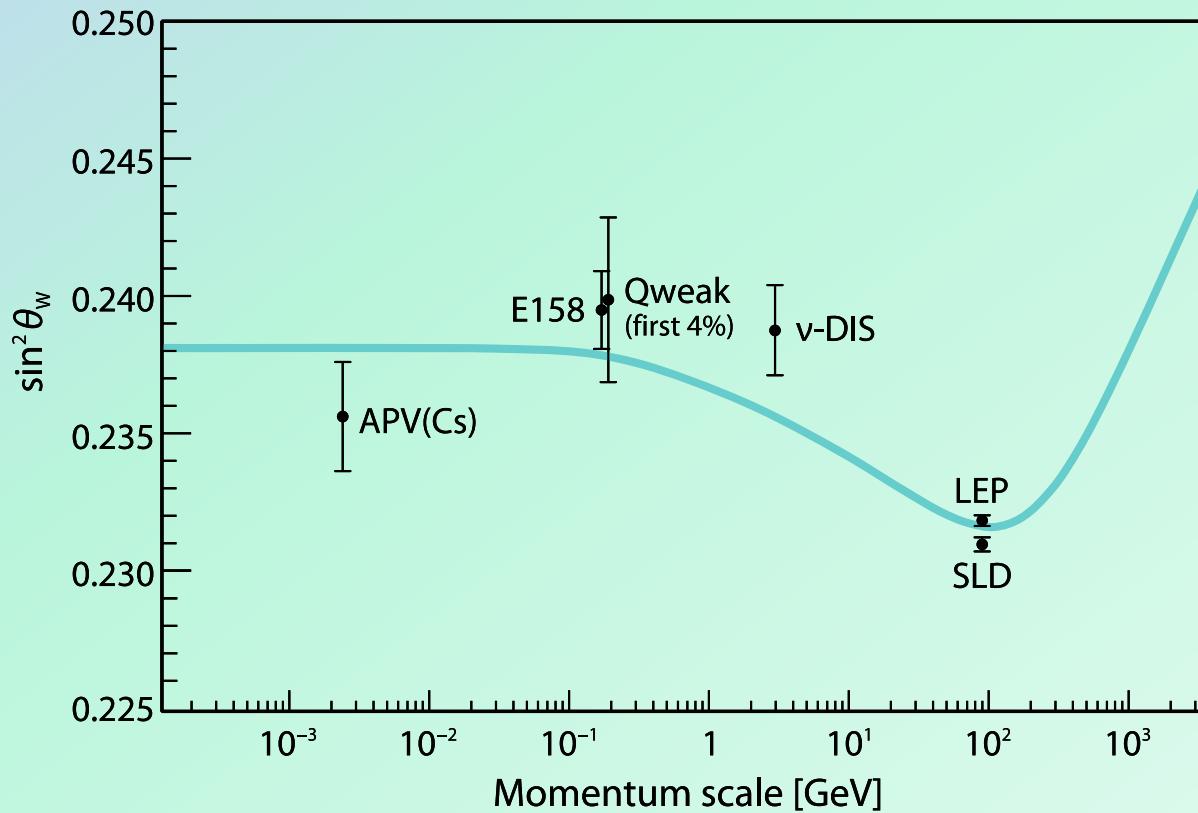
(full LHC $M_{Z'} \sim 4.5 \text{ TeV}/c^2$)

The way to go!

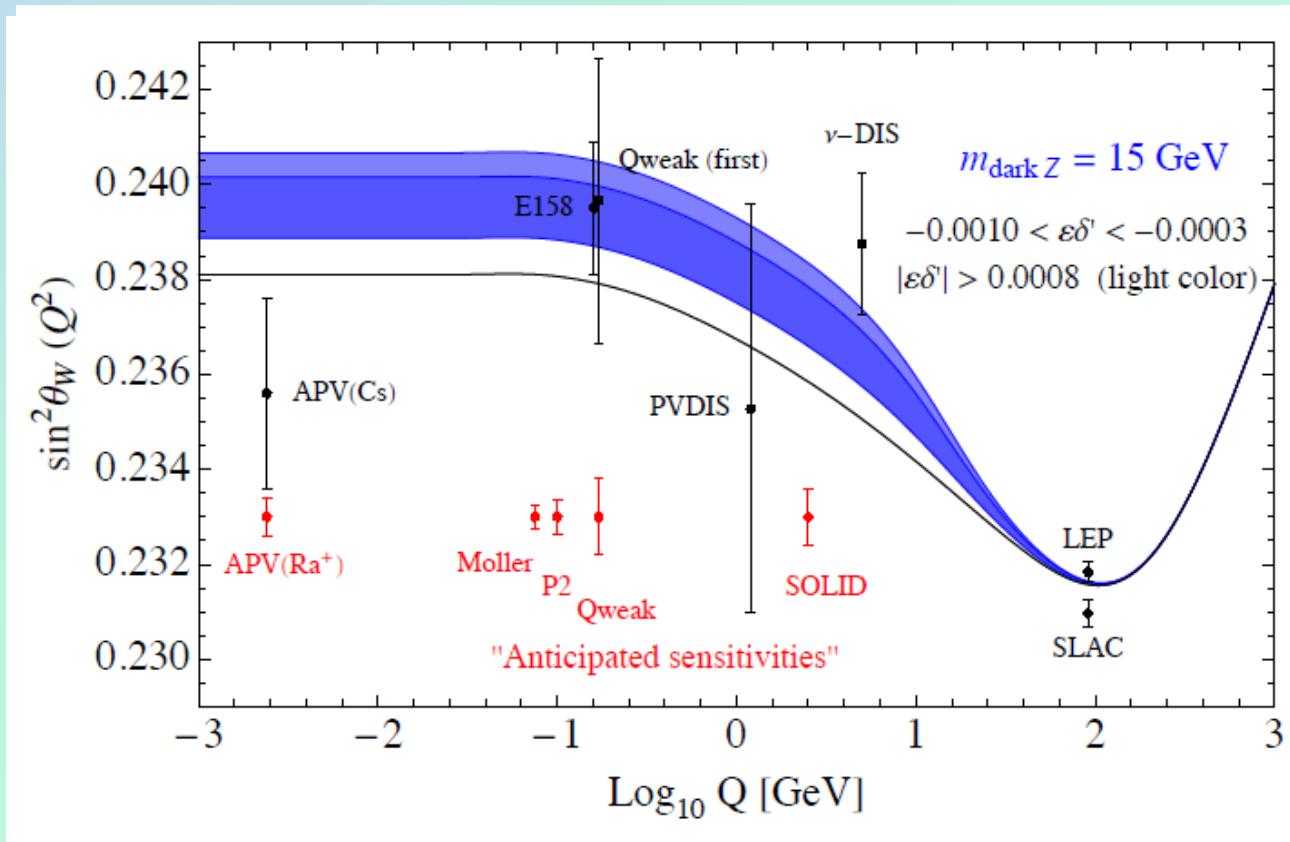
Test of Standard Model Electroweak Interaction



Test of Standard Model Electroweak Interaction



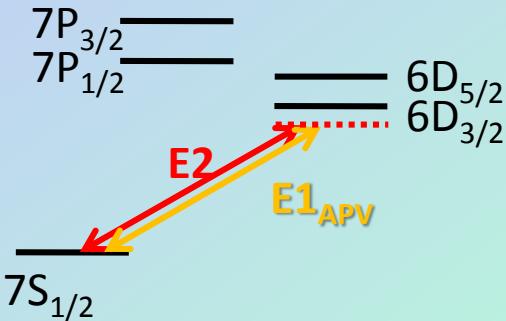
Test of Standard Model Electroweak Interaction



- S. Kumar, W. Marciano, Annu. Rev. of Nucl. Part. Sci. **63**, 237 (2013)
H. Davoudiasl, Hye-Sung Lee, W. Marciano, arxiv. 1402.3620 (2014)
H. Davoudiasl, H. S. Lee and W. J. Marciano, Phys. Rev. **D 92**, 055005 (2015)

Atomic Parity Violation

Ba⁺ and Ra⁺

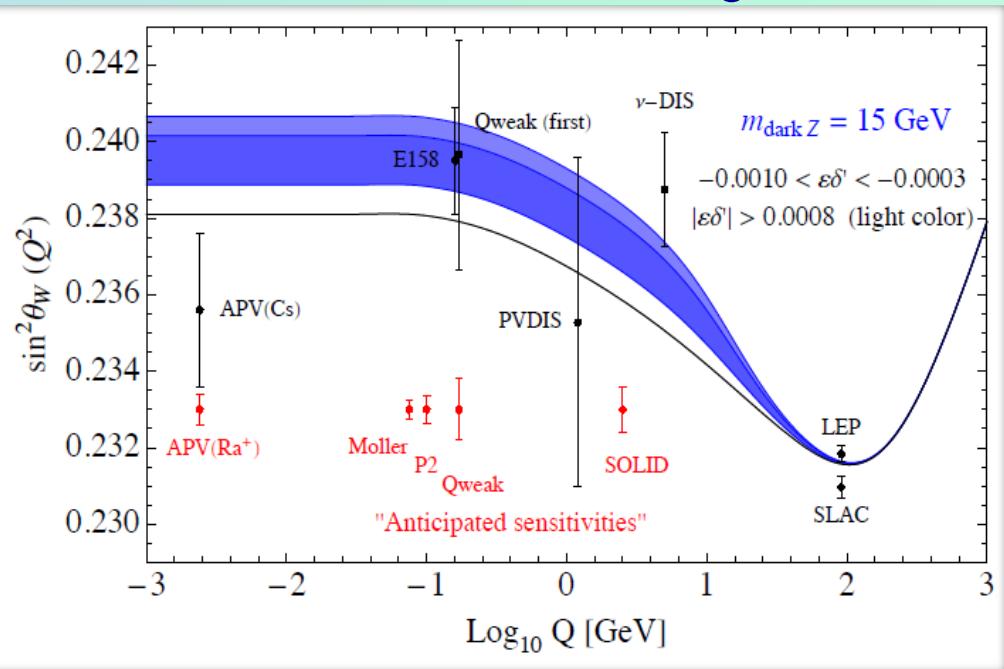


$$Q_W = \frac{E1_{APV}}{k}$$

Calculated from
atomic wavefunctions

Detailed calculations → stronger than Z³

S-S	S-D
Cs	Ba ⁺
0.9	2.2
Fr	Ra ⁺
14.2	46.4



Ra⁺ superior to measure APV ...
50x more sensitive to APV than current best measurement in Cs

Theory Calculations:

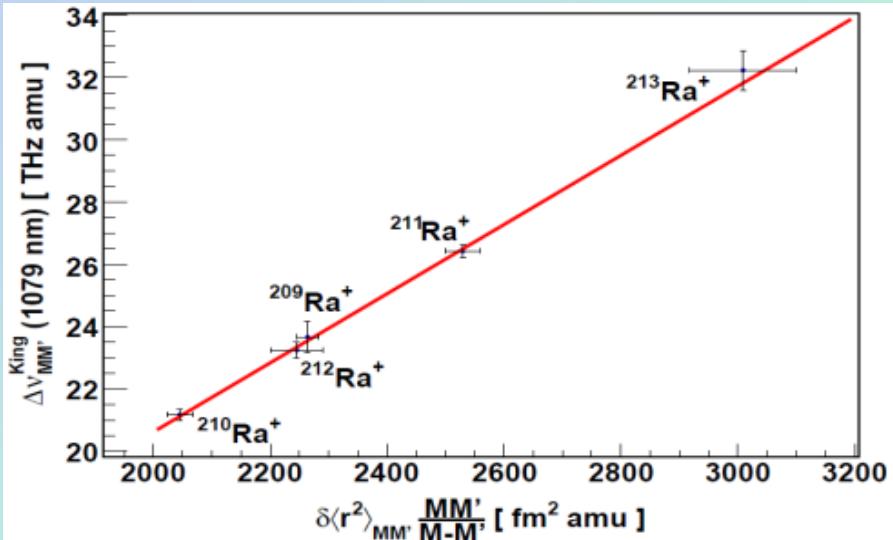
$$k_{Ra} = 46.4(1.4) \cdot 10^{-11} \text{ iea}_0 / N \quad *$$

$$k_{Cs} = 0.8906(26) \cdot 10^{-11} \text{ iea}_0 / N \quad **$$

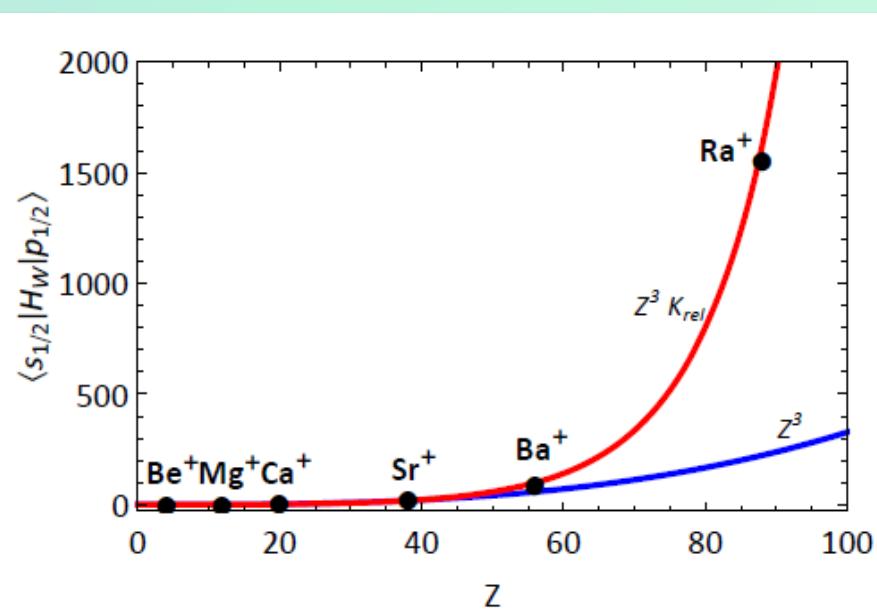
*L.W. Wansbeek et al., Phys. Rev. A 78, (2008)

**A. Derevianko et al., Phys. Rev. A 79, 013404 (2009)

Laser Spectroscopy in Ra⁺ ions



Probe of atomic theory & size and shape of the nucleus



Probe of atomic wave functions at the origin

Good agreement with theory at few % level
Theory improvement is in pipeline.

- O.O. Versolato et al., Phys. Lett. A 375, 3130 (2012)
O.O. Versolato et al., Phys. Rev. A 82, 010501(R) (2010)
G.S. Giri et al., Phys. Rev. A 84, 020503(R) (2011)

TABLE VI. The radial differences $\delta\langle r^2 \rangle = \langle r^2 \rangle_i - \langle r^2 \rangle_0$ in fm² extracted using the database of Table II, and corrected for nuclear deformation (column 8), compared with the values listed in the Atomic Data and Nuclear Data Tables [20] (column 9, originally from ISOLDE [25]; the square brackets indicate that the 10% error needs to be taken towards higher absolute values). The factor B^i of Eq. (4) is given in MHz in column 2. Column 3 gives the corresponding results of Sec. III for the sharp cutoff nucleus using coupled-cluster (CC) theory for the field shift and specific mass shift for the reference transition. Columns 4, 5, and 6 give the results of Sec. IV needed to correct for nuclear deformation, with a sharp cutoff and Fermi nucleus using the Dirac-Fock (DF) approach. “ Δ ” in column 7 gives the relative difference in % between the sharp cutoff model in column 4 and the Fermi model with deformation in column 6. We use Δ to correct the value of column 3, which gives column 8. Because for $A = 209$ and 232 the reference transition at 468 nm in Ra⁺ is not available, the effect of the deformation could not be calculated; for these two isotopes, the results for Δ (in italics) are estimated by extrapolation of the results of the neighboring isotopes.

A	B^i in MHz	$\delta\langle r^2 \rangle$ in fm ²				Δ in %	$\delta\langle r^2 \rangle$ in fm ²	
		Sharp (CC)	Sharp (DF)	Fermi	Fermi (def.)		This work	ISOLDE
208	11 860(15)	-0.298(15)	-0.357	-0.361	-0.362	1.3	-0.302(16)	-0.256[27]
209	11 630(15)	-0.292(15)	-	-	-	1.3	-0.296(15)	-0.253[25]
210	8393(11)	-0.211(11)	-0.252	-0.255	-0.256	1.3	-0.214(11)	-0.182[19]
211	7728.9(70)	-0.1941(98)	-0.2320	-0.2347	-0.2364	0.2	-0.1941(98)	-0.1941(98)
212	4554.9(44)	-0.1144(58)	-0.1369	-0.1386	-0.1396	0.1	-0.1144(58)	-0.1144(58)
213	3035.3(41)	-0.0762(39)	-0.0911	-0.0922	-0.0932	0.0	-0.0762(39)	-0.0762(39)
214	0	0	0	0	0	0	-0.0000(00)	-0.0000(00)
220	-30 731(18)	0.772(39)	0.920	0.946	0.959	0.9	0.772(39)	0.772(39)
221	-36 402(19)	0.914(47)	1.090	1.127	1.151	1.1	0.914(47)	0.914(47)
222	-40 444(21)	1.016(52)	1.211	1.269	1.311	1.2	1.016(52)	1.016(52)
223	-45 533(23)	1.144(58)	1.364	1.414	1.464	1.3	1.144(58)	1.144(58)
224	-49 274(24)	1.238(63)	1.476	1.535	1.585	1.5	1.238(63)	1.238(63)
225	-54 560(27)	1.370(69)	1.634	1.701	1.761	1.6	1.370(69)	1.370(69)
226	-57 692(29)	1.449(73)	1.728	1.805	1.771	2.5	1.486(75)	1.277[129]
227	-61 638(32)	1.548(78)	1.846	1.929	1.892	2.5	1.587(80)	1.365[138]
228	-65864(34)	1.654(84)	1.973	2.068	2.023	2.6	1.697(86)	1.459[148]
229	-70235(37)	1.764(89)	2.103	2.208	2.158	2.6	1.810(92)	1.556[158]
230	-75243(39)	1.890(95)	2.253	2.391	2.317	2.8	1.943(98)	1.667[169]
232	-83590(47)	2.10(11)	-	-	-	2.8	2.16(11)	1.854[188]

available $\delta(r^2)$ values:

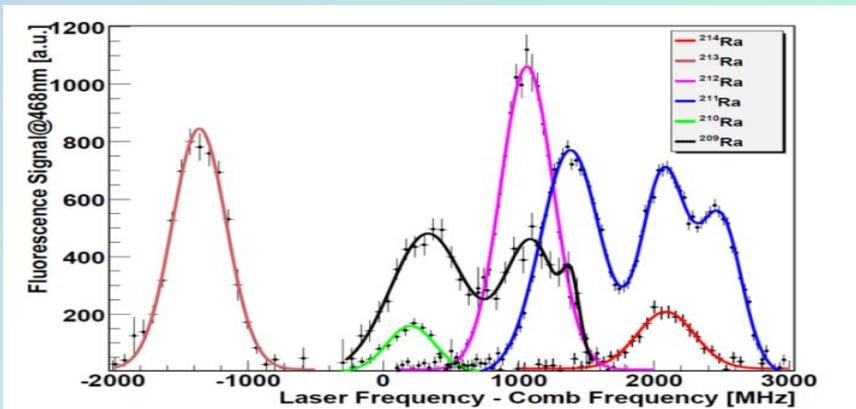
1.486(75) fm²

vs.

1.277(129) fm²

Single Ra⁺ and Ba⁺ Ions

$$Q_W = \frac{E1_{APV}}{k} \quad \text{To be measured}$$

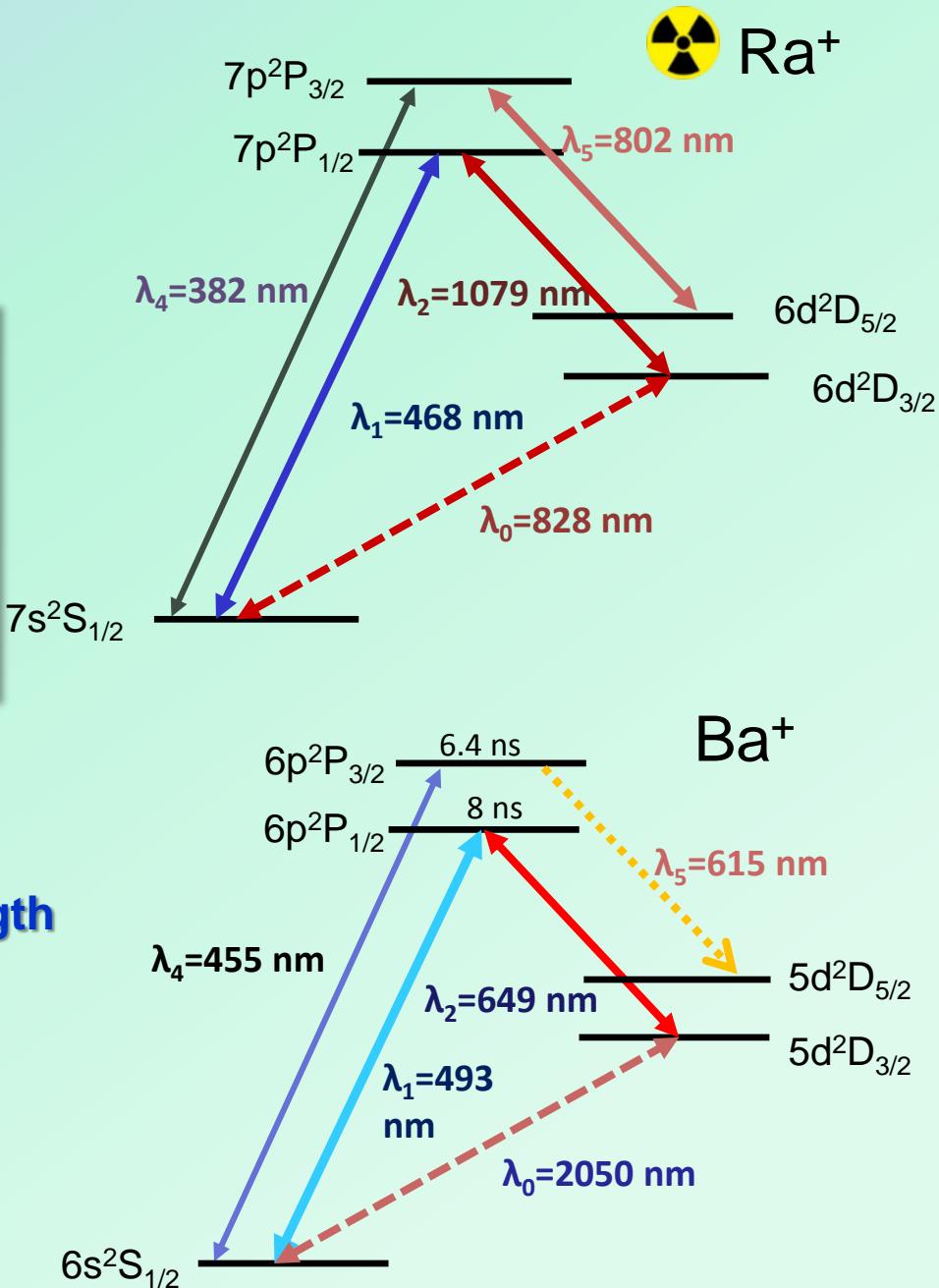


\leftrightarrow
5mm

Hyperbolic Paul Trap

- **localize one ion within one wavelength**
- **electron shelving**
- **large volume**

Ba⁺ : Precursor to Ra⁺



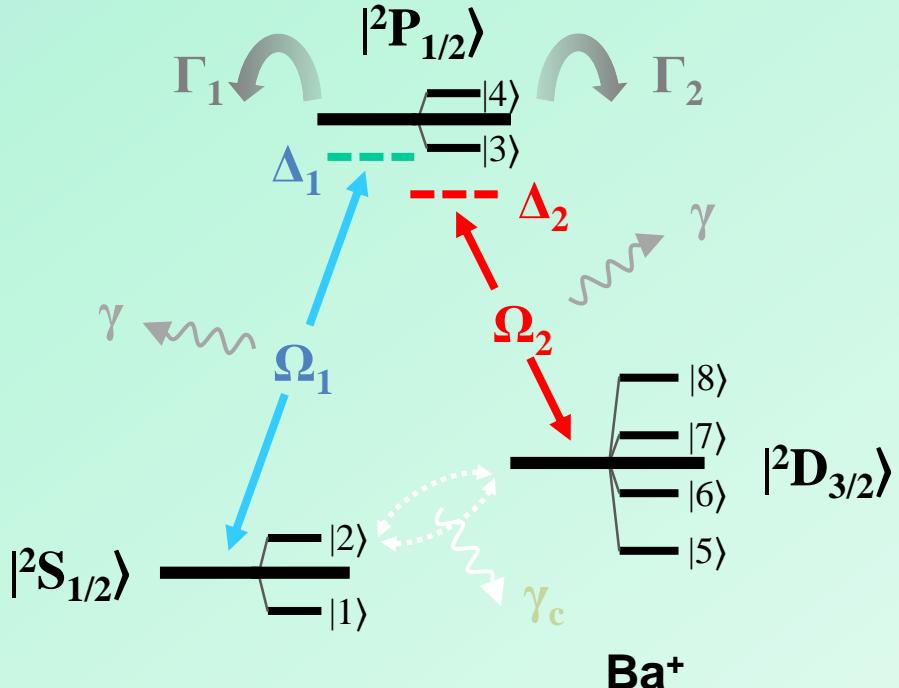
Modeling of Line Shape

- Optical Bloch equation
3 level example

$$\frac{d}{dt}\rho_{ij} = \frac{i}{\hbar} [H, \rho] + R(\rho)$$

$$H = \hbar \begin{pmatrix} \Delta_1 - \omega_B & 0 & -\frac{1}{\sqrt{2}}\Omega_1 & 0 & 0 & 0 & 0 & 0 \\ 0 & \Delta_1 + \omega_B & 0 & \frac{1}{\sqrt{2}}\Omega_1 & 0 & 0 & 0 & 0 \\ -\frac{1}{\sqrt{2}}\Omega_1 & 0 & -\frac{1}{\sqrt{2}}\omega_B & 0 & \frac{1}{\sqrt{2}}\Omega_2 & \frac{1}{\sqrt{2}}\Omega_2 & -\frac{1}{\sqrt{2}}\Omega_2 & 0 \\ 0 & \frac{1}{\sqrt{2}}\Omega_1 & 0 & \frac{1}{\sqrt{2}}\omega_B & 0 & \frac{1}{\sqrt{2}}\Omega_2 & \frac{1}{\sqrt{2}}\Omega_2 & -\frac{1}{\sqrt{2}}\Omega_2 \\ 0 & 0 & -\frac{1}{\sqrt{2}}\Omega_2 & 0 & \Delta_2 - \frac{5}{3}\omega_B & 0 & 0 & 0 \\ 0 & 0 & \frac{1}{\sqrt{2}}\Omega_2 & -\frac{1}{\sqrt{2}}\Omega_2 & 0 & \Delta_2 - \frac{1}{3}\omega_B & 0 & 0 \\ 0 & 0 & \frac{1}{\sqrt{2}}\Omega_2 & \frac{1}{\sqrt{2}}\Omega_2 & 0 & 0 & \Delta_2 + \frac{1}{3}\omega_B & 0 \\ 0 & 0 & 0 & \frac{1}{\sqrt{2}}\Omega_2 & 0 & 0 & 0 & \Delta_2 + \frac{6}{3}\omega_B \end{pmatrix}$$

$$R(\rho) = \begin{pmatrix} \Gamma_1\left(\frac{1}{2}\rho_{13} + \frac{1}{2}\rho_{14}\right) & -\Gamma_1\frac{1}{2}\rho_{14} & -\gamma'\rho_{13} & -\gamma'\rho_{14} & -\gamma_1\rho_{15} & -\gamma_1\rho_{17} & -\gamma_1\rho_{18} \\ -\Gamma_1\frac{1}{2}\rho_{13} & \Gamma_1\left(\frac{1}{2}\rho_{13} + \frac{1}{2}\rho_{14}\right) & -\gamma'\rho_{13} & -\gamma'\rho_{14} & -\gamma_1\rho_{25} & -\gamma_1\rho_{27} & -\gamma_1\rho_{28} \\ -\gamma'\rho_{13} & -\gamma'\rho_{12} & -\Gamma_2\rho_{13} & -\Gamma_2\rho_{14} & -\gamma'\rho_{13} & -\gamma'\rho_{16} & -\gamma'\rho_{17} \\ -\gamma'\rho_{14} & -\gamma'\rho_{12} & -\Gamma_2\rho_{14} & -\Gamma_2\rho_{13} & -\gamma'\rho_{14} & -\gamma'\rho_{16} & -\gamma'\rho_{18} \\ -\gamma_1\rho_{15} & -\gamma_1\rho_{12} & -\gamma'\rho_{13} & -\gamma'\rho_{14} & \Gamma_2\frac{1}{2}\rho_{13} & \Gamma_2\frac{1}{2}\rho_{14} & 0 \\ -\gamma_1\rho_{17} & -\gamma_1\rho_{12} & -\gamma'\rho_{13} & -\gamma'\rho_{14} & \Gamma_2\frac{1}{2}\rho_{13} & \Gamma_2\frac{1}{2}\rho_{14} & 0 \\ -\gamma_1\rho_{18} & -\gamma_1\rho_{12} & -\gamma'\rho_{13} & -\gamma'\rho_{14} & 0 & \Gamma_2\frac{1}{2}\rho_{14} & \Gamma_2\frac{1}{2}\rho_{14} \\ -\gamma_1\rho_{12} & -\gamma_1\rho_{12} & -\gamma'\rho_{13} & -\gamma'\rho_{14} & 0 & \Gamma_2\frac{1}{2}\rho_{14} & \Gamma_2\frac{1}{2}\rho_{14} \end{pmatrix}$$



Ω_1, Ω_2 Rabi frequencies
(laser power)

Δ_1, Δ_2 laser detunings

$\Gamma = \Gamma_1 + \Gamma_2$ relaxation rate
 $\gamma = \Gamma/2$ decoherence rate

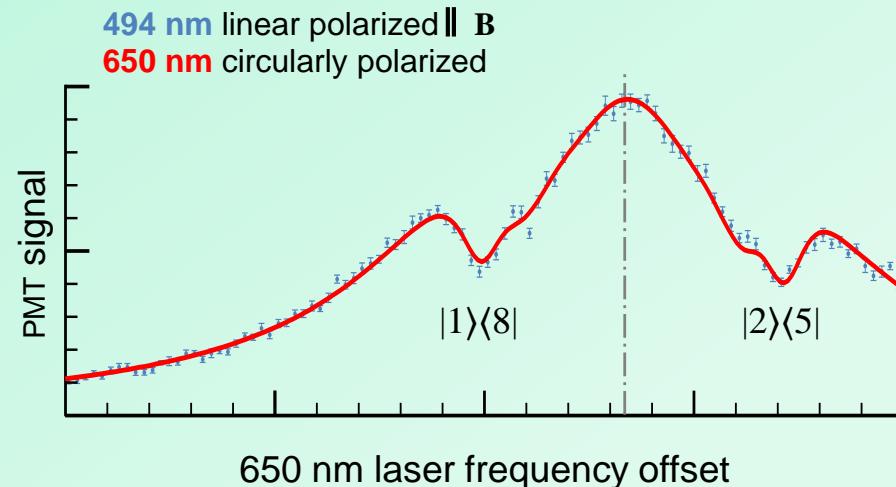
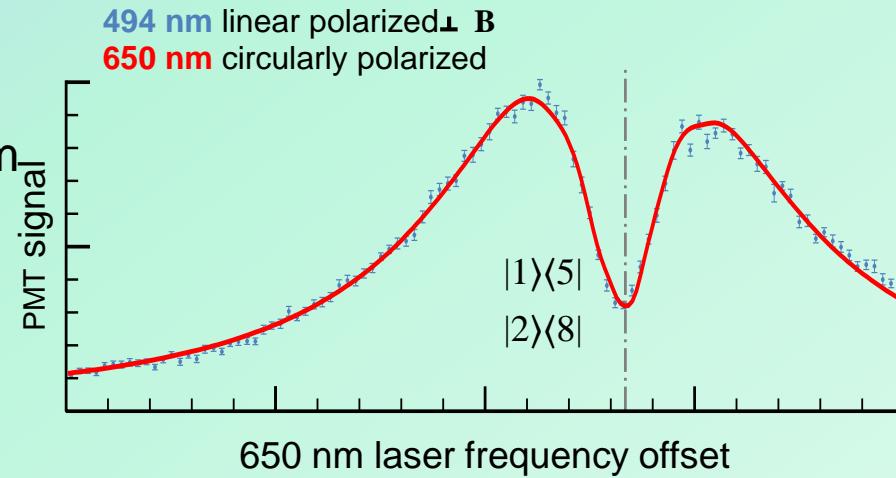
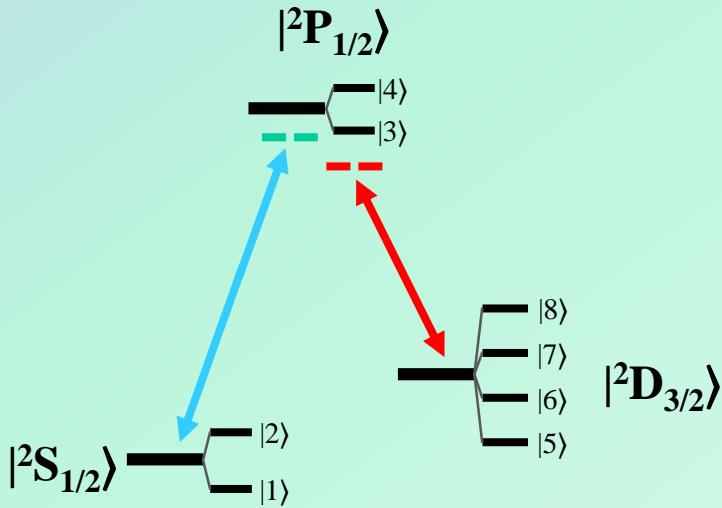
γ_c laser linewidth

Line Shapes and Polarization

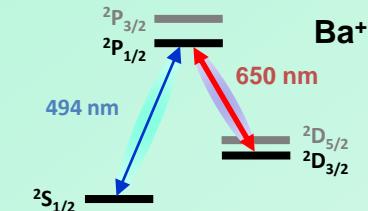
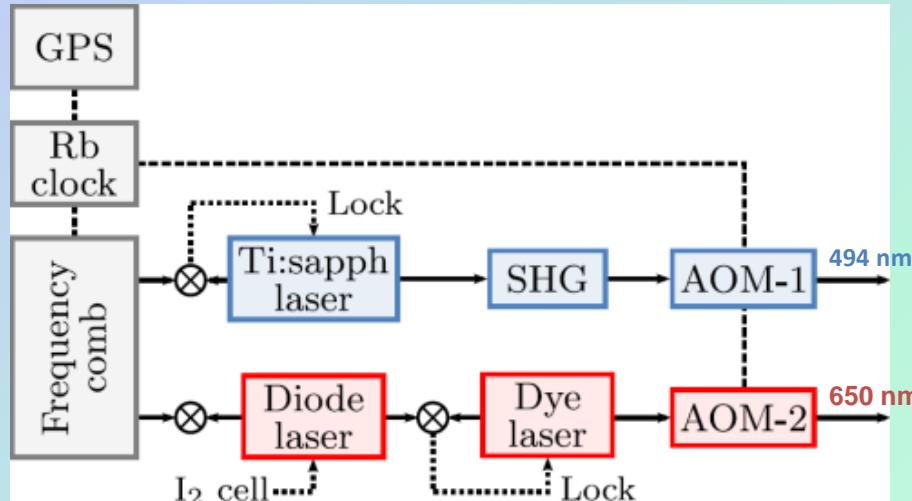
- Zeeman sublevels: 8 level system

Magnetic field **B**

Laser polarization

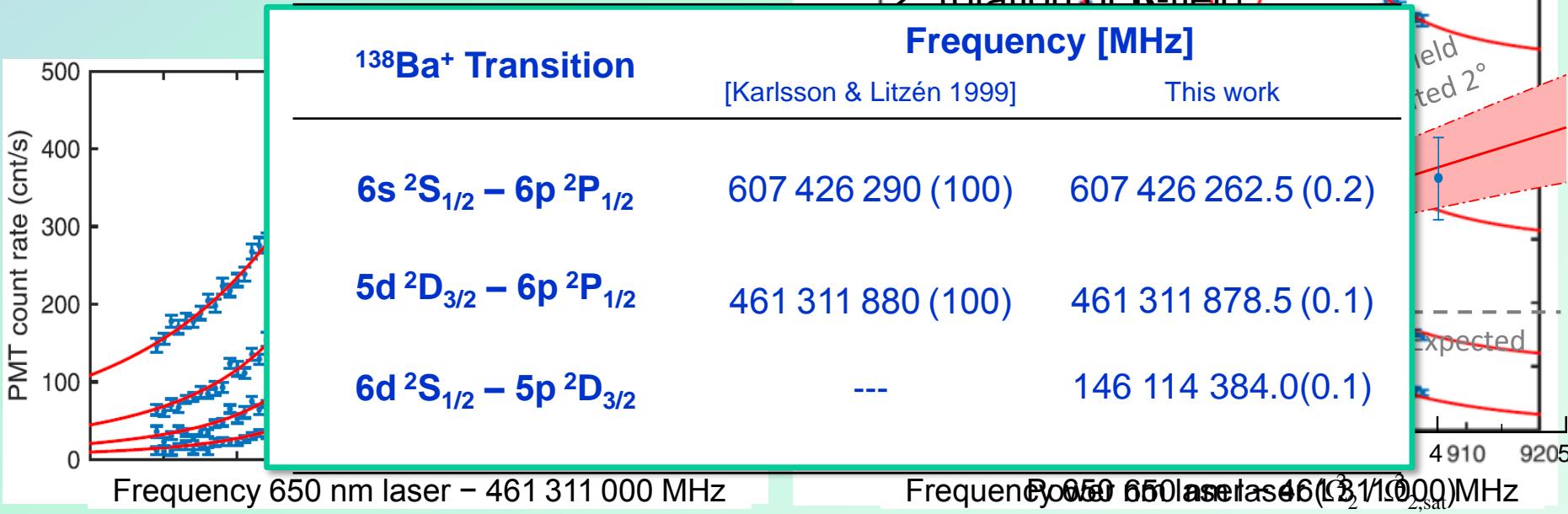


Transition Frequencies



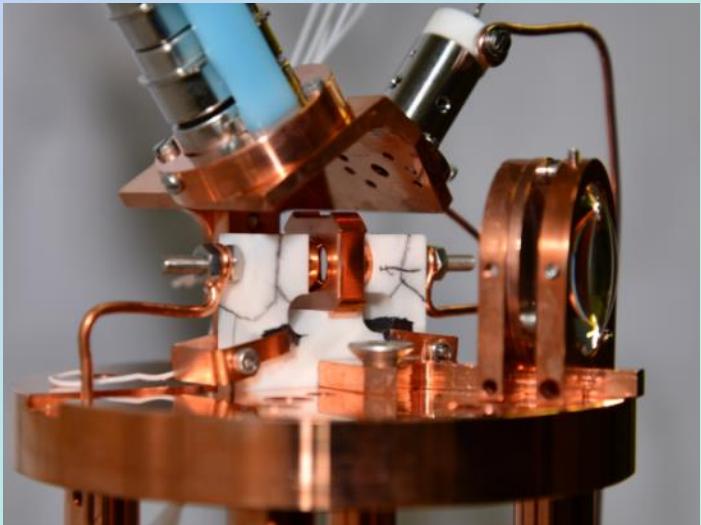
Light shift?

No, correction in transition frequencies for Ω_2 dependent shift consistent with 2° rotation of B-field

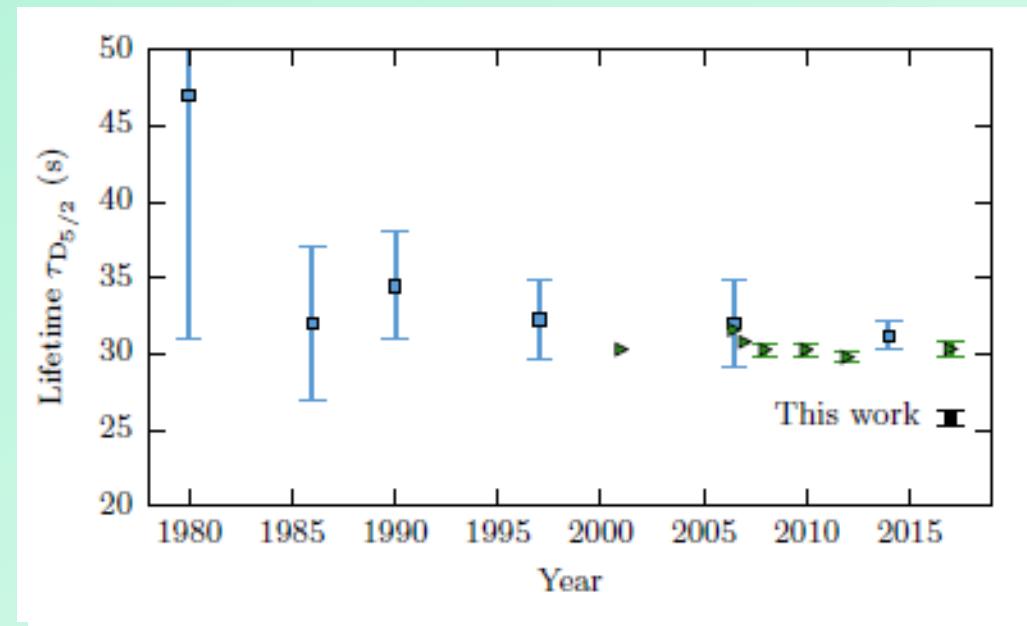
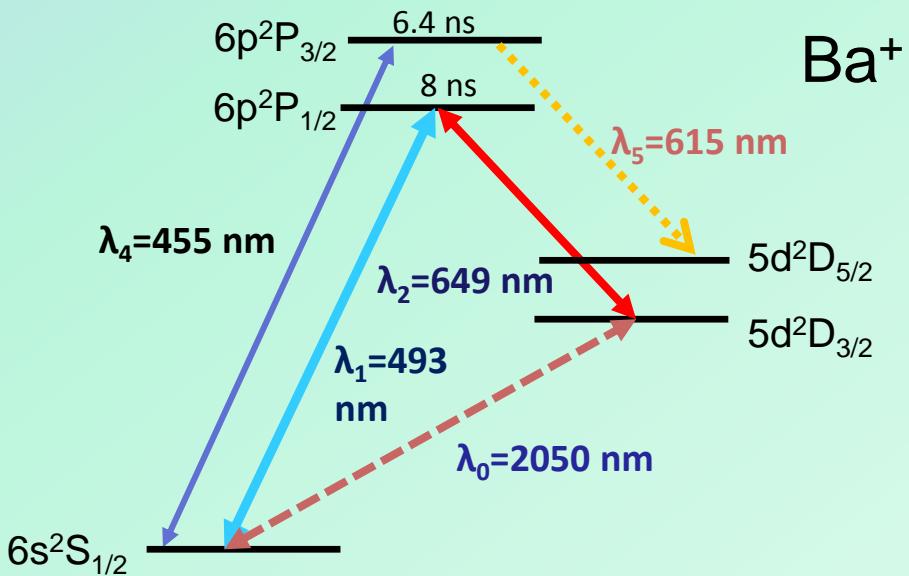


- Data fit to optical Bloch equation model
- Extract transition frequencies with 100 kHz accuracy

Ba⁺ Experiment : Lifetime D_{5/2}



Ba⁺ D_{5/2} state lifetime



$$\tau_{D5/2} = 25.8(5) \text{ s}$$

E.A. Dijck et al, submitted (2017)

Radium for APV

Accuracy of single ion Experiment

$$\frac{\mathcal{E}^{\text{PNC}}}{\delta \mathcal{E}^{\text{PNC}}} \cong \frac{\mathcal{E}^{\text{PNC}} E_0}{\hbar} f \sqrt{N \tau t}$$

E_0 = Light electric field amplitude, τ = Coherence time
 N = Number of ions = 1, t = Time of observation

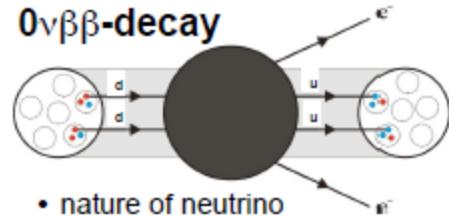
	Coherence Time	Projected Accuracy	Measurement Time
Ba ⁺	80 sec	0.2%	1.1 day
Ra ⁺	0.6 sec	0.2%	1.4 day

55	56
Cs	Ba
0.9	2.2
87	88
Fr	Ra
14.2	46.4

→ 10 days for 5 fold improvement over Cs

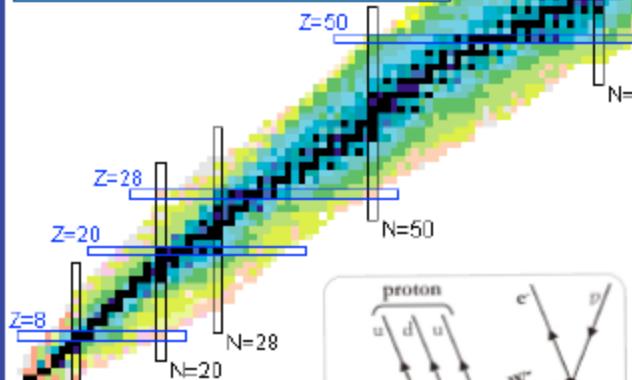
weak interaction studies in radionuclides

$0\nu\beta\beta$ -decay



- nature of neutrino
- nuclear matrix elements
- effective Majorana mass

F. T. Avignone et al., Rev. Mod. Phys. 80, 481 (2008)



β decays

- V_{ud} of CKM matrix
- CKM unitarity test
- limits on scalar & tensor currents
- ...

J. C. Hardy and I. S. Towner, Phys. Rev. C 91, 025501 (2015)
N. Severijns and O. Naviliat-Cuncic, Phys. Scr. T152, 014018 (2013)
V. Cirigliano et al., Prog. Part. Nucl. Phys. 71, 93 (2013)
O. Naviliat-Cuncic and M. González-Alonso, Ann. Phys. 525, 600 (2013)
K. K. Vos et al., Rev. Mod. Phys. 87, 1483 (2015)

nuclear masses for neutrino physics

- for Δm^2 mass measurements
- identify best cases for $0\nu\beta\beta$

S. Eliseev et al., Ann. Phys. 525, 707 (2013)

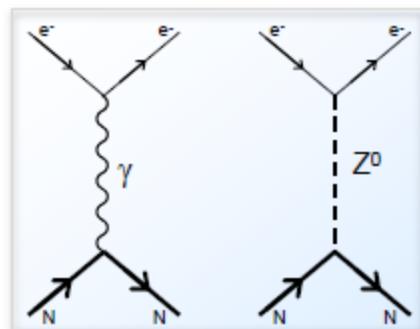
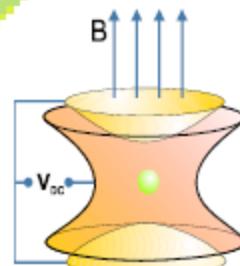
2



Octupole enhanced atomic EDMs

- in 'pear shaped' nuclei

L. P. Gaffney, Nature 497, 199–204 (2013)
J. Dobaczewski and J. Engel, PRL 94, 232502 (2005)



atomic parity violation

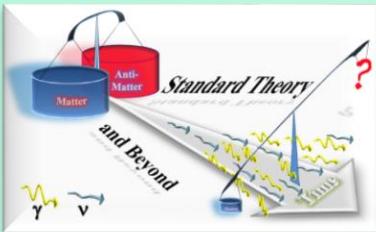
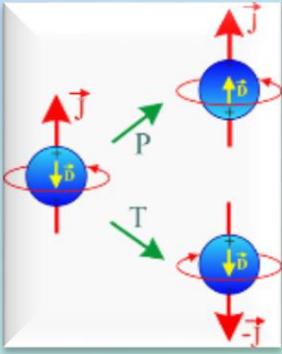
- access $\sin^2(\Theta_W)$ at low energy
- strong enhancement in (radioactive) Fr or Ra⁺

S. Aubin et al., Hyp. Int. 214, 163 (2013).
L. Willmann et al., CERN-INTC-2017-066 / INTC-I-196

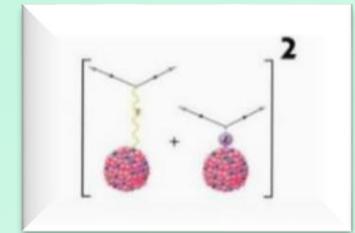
ISOLDE

SUMMARY

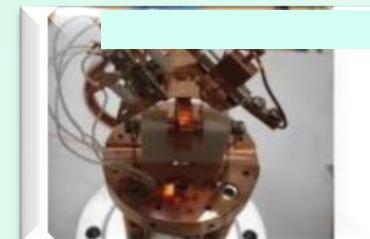
Precision Tests of Discrete Symmetries at Low Energies



- A few Selected Experiments
 - Focus on Transformativity
- C, P, CP, CPT Tests
 - Precision Test of Standard Model
- Experiment & Theory Hand in Hand
 - Atomic Parity Violation and EDM to search for New Physics
- Search for permanent Electric Dipole Moments
 - Atoms & Molecules with Enhancement
 - Electron and Nucleon EDMs
 - Some Radioactive Species may have Advantages
 - No particular advantage from Radioactivity per se
 - Central Goal: Challenge New Physics Models



THANK YOU !



Spares