

g_A, g_S, g_T from Lattice QCD

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Acknowledgements

Received Data From

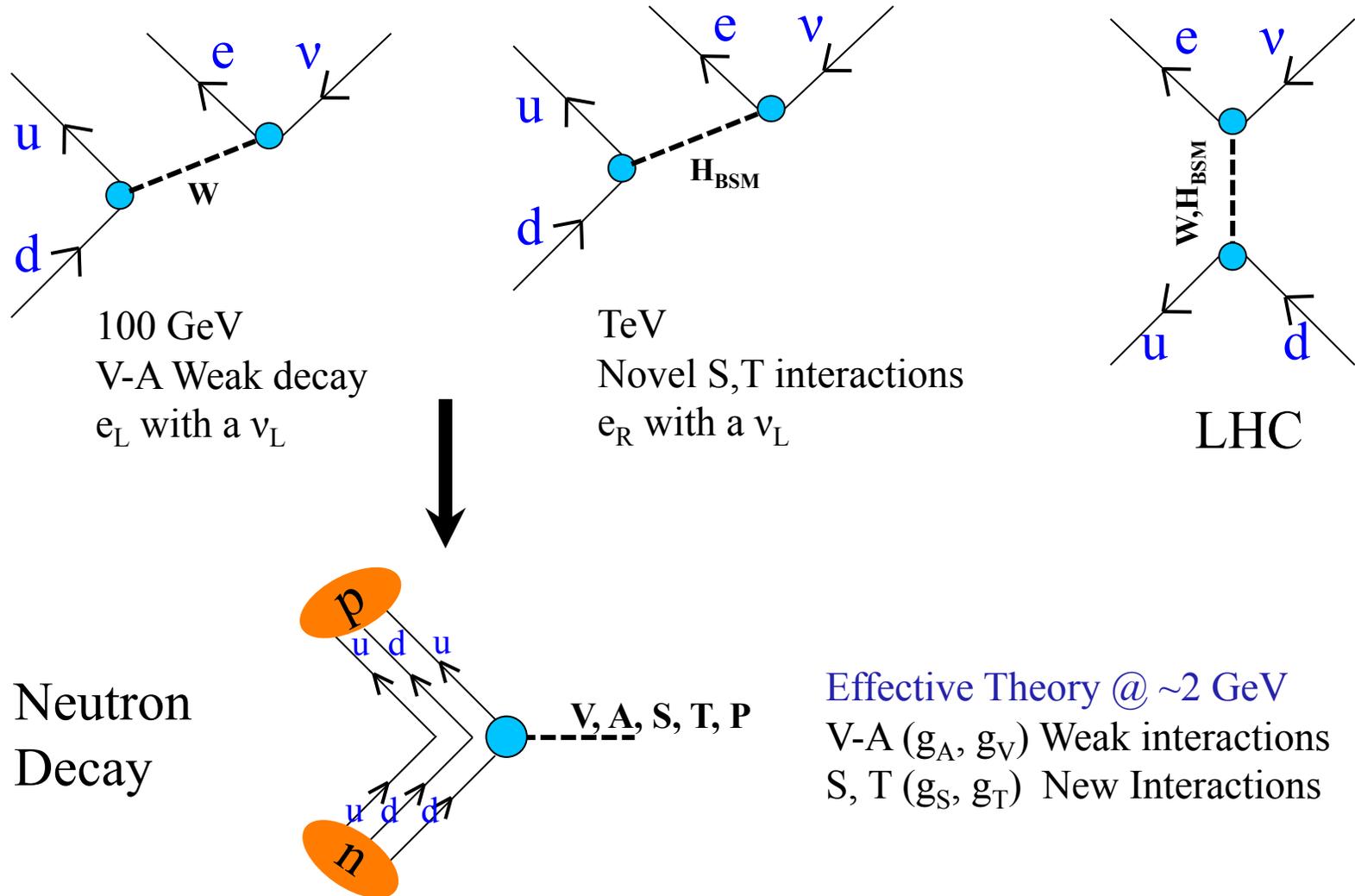
- Gunnar Bali (RQCD -- Regensburg)
- Martha Constantinou (ETMC)
- Jeremy Green (LHPC)
- Keh-Fei Liu (χ QCD)
- Shigemi Ohta (RBC/UKQCD)
- Hartmut Wittig (Mainz)

Discussions with

- Boram Yoon, T. Bhattacharya, V. Cirigliano (LANL)
- Huey-Wen Lin, Saul Cohen (U. Wash.)

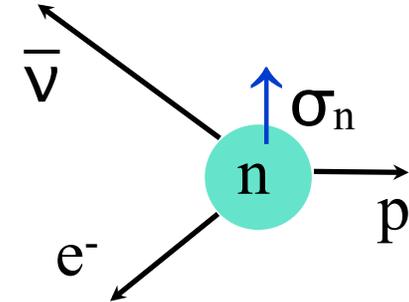
Probing New Interactions: $M_{\text{BSM}} \gg M_W \gg 1 \text{ GeV}$

Many BSM possibilities for novel Scalar & Tensor interactions: Higgs-like, leptoquark, loop effects, ...



[Ultra] Cold Neutron Decay: Terms sensitive to new physics

Neutron decay can be parameterized as



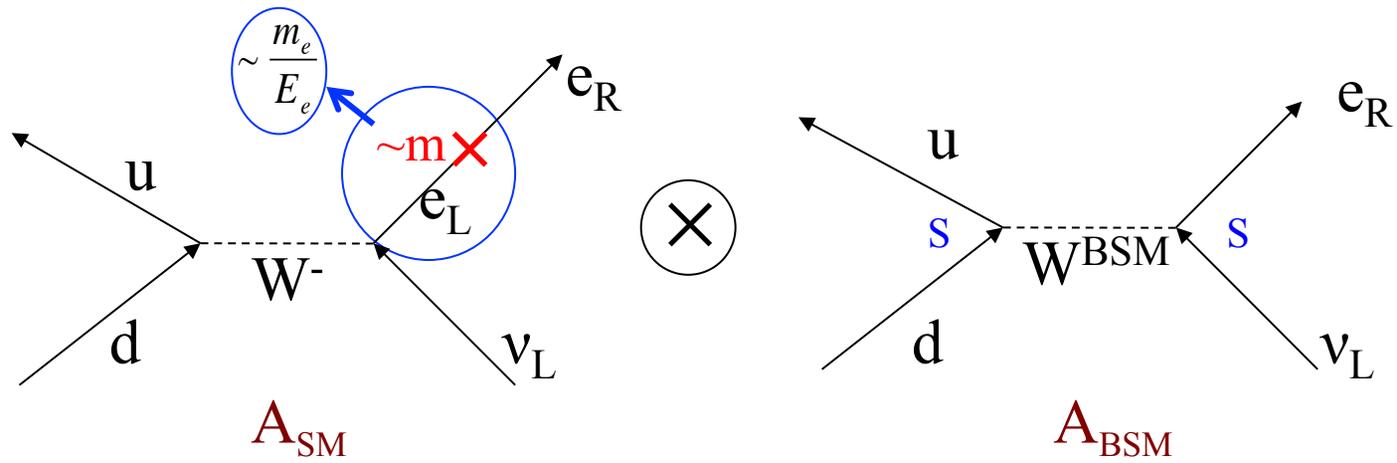
$$d\Gamma \propto F(E_e) \left[1 + b \frac{m_e}{E_e} + \left(B_0 + B_1 \frac{m_e}{E_e} \right) \frac{\vec{\sigma}_n \cdot \vec{p}_\nu}{E_\nu} + \dots \right]$$

- b : Deviations from the leading order electron spectrum:
Fierz interference term
- B_1 : Energy dependent part of antineutrino correlation
with neutron spin

Novel S and T interactions at TeV scale

→ Effective Theory at low energy

At leading order, contributions from BSM physics arise due to interference of A_{SM} and A_{BSM} and contribute to b and B_1 only through ϵ_S and ϵ_T



$$H_{\text{eff}} = G_F \left[J_{V-A}^{\text{lept}} \times J_{V-A}^{\text{quark}} + \sum_{n=1,10} \epsilon_n^{\text{BSM}} \hat{O}_n \right]$$

$\epsilon_S \bar{u}d \times \bar{e}(1 - \gamma_5)\nu_e$
 $\epsilon_T \bar{u}\sigma_{\mu\nu}d \times \bar{e}\sigma^{\mu\nu}(1 - \gamma_5)\nu_e$

Physics Case: (BSM/SM) $\sim O(1)$

- Couplings $\varepsilon_{P,S,T} \sim (\Lambda_{\text{BSM}})^2/G_F \sim (v/\Lambda_{\text{BSM}})^2 \sim 10^{-3}$
- Recoil corrections: $q/M_N \sim 10^{-3}$
- Radiative corrections: $\alpha_{\text{em}}/\pi \sim 10^{-3}$
- Isospin-breaking: $(M_N - M_P)/M_N \sim q/M_N \sim 10^{-3}$
- UCN: small Doppler broadening of e spectrum
- SM contribution is $O(10^{-3})$ and known to ($\sim 10^{-5}$): It is controlled by 2 small parameters $(M_n - M_p)/M_n$ and α_{em}/π
- **Unique:** scalar and tensor BSM interactions involve helicity-flip (m_e/E_e suppression) and are hard to detect in high energy experiments

Relating b , B_1 to $g_{S,T}$ & BSM couplings $\varepsilon_{S,T}$

$$H_{eff} \supset G_F \left[\varepsilon_S \boxed{\bar{u}d} \times \bar{e}(1-\gamma_5)\nu_e + \varepsilon_T \boxed{\bar{u}\sigma_{\mu\nu}d} \times \bar{e}\sigma^{\mu\nu}(1-\gamma_5)\nu_e \right]$$

$$g_S = Z_S \langle p | \bar{u}d | n \rangle \quad g_T = Z_T \langle p | \bar{u}\sigma_{\mu\nu}d | n \rangle \quad \boxed{\text{Lattice QCD}}$$

Linear order relations from $n \rightarrow p e \bar{\nu}$ decay

$$b^{BSM} \approx 0.34 g_S \varepsilon_S - 5.22 g_T \varepsilon_T$$

$$b_\nu^{BSM} \equiv B_1^{BSM} = E_e \frac{\partial B^{BSM}(E_e)}{\partial m_e} \approx 0.44 g_S \varepsilon_S - 4.85 g_T \varepsilon_T$$

What we know

- Experiment

- $g_A = 1.2701(25)$

Neutron decay

- Phenomenology: CVC

$$\frac{g_S}{g_V} = \frac{(M_N - M_P)^{QCD}}{(m_d - m_u)^{QCD}} = 1.02(8)(7)$$

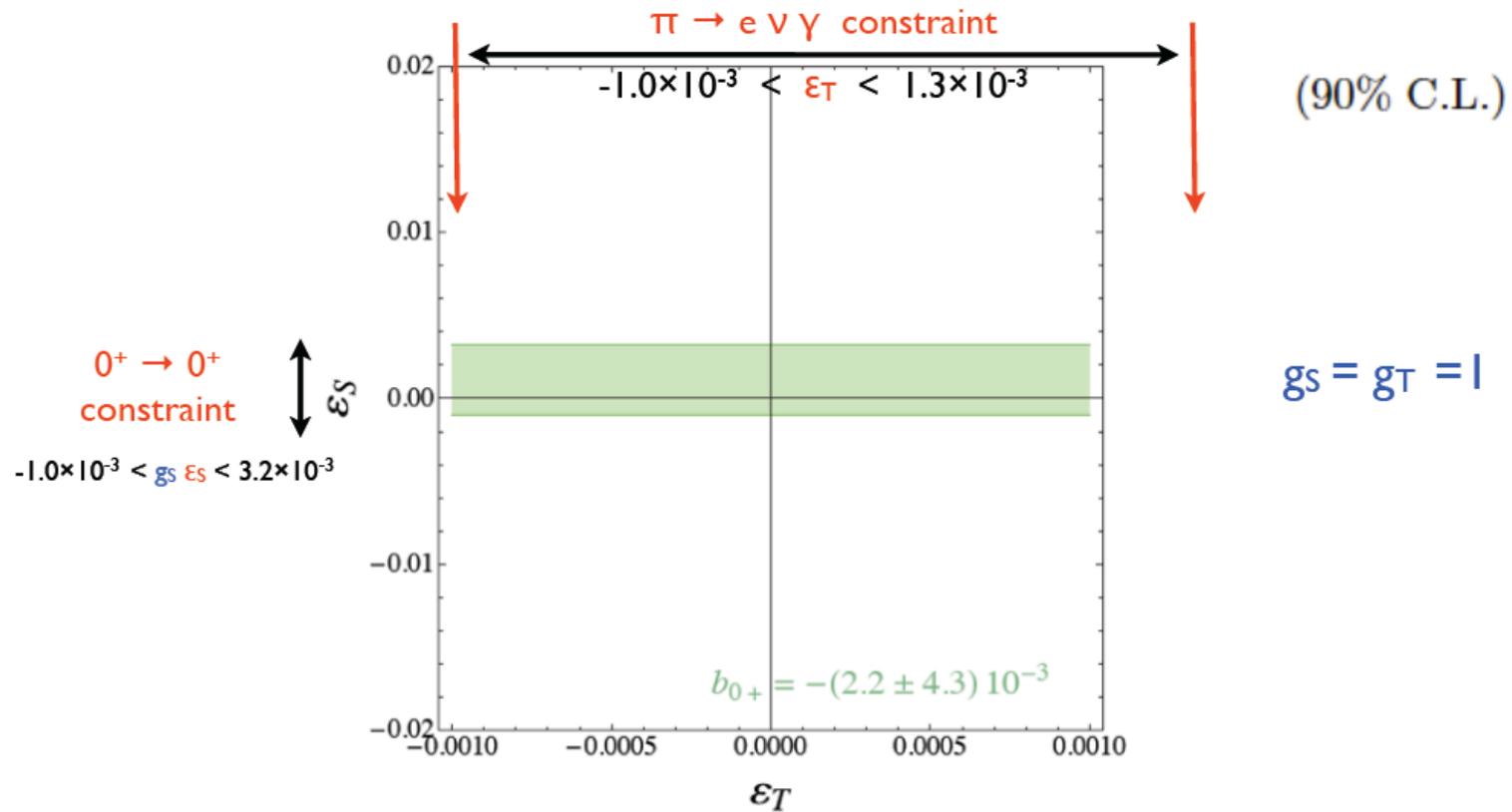
Gonzalez-Alonso & Camalich
Phy. Rev. Lett. 112 (2014) 042501

Lattice QCD can provide precise estimates of nucleon structure

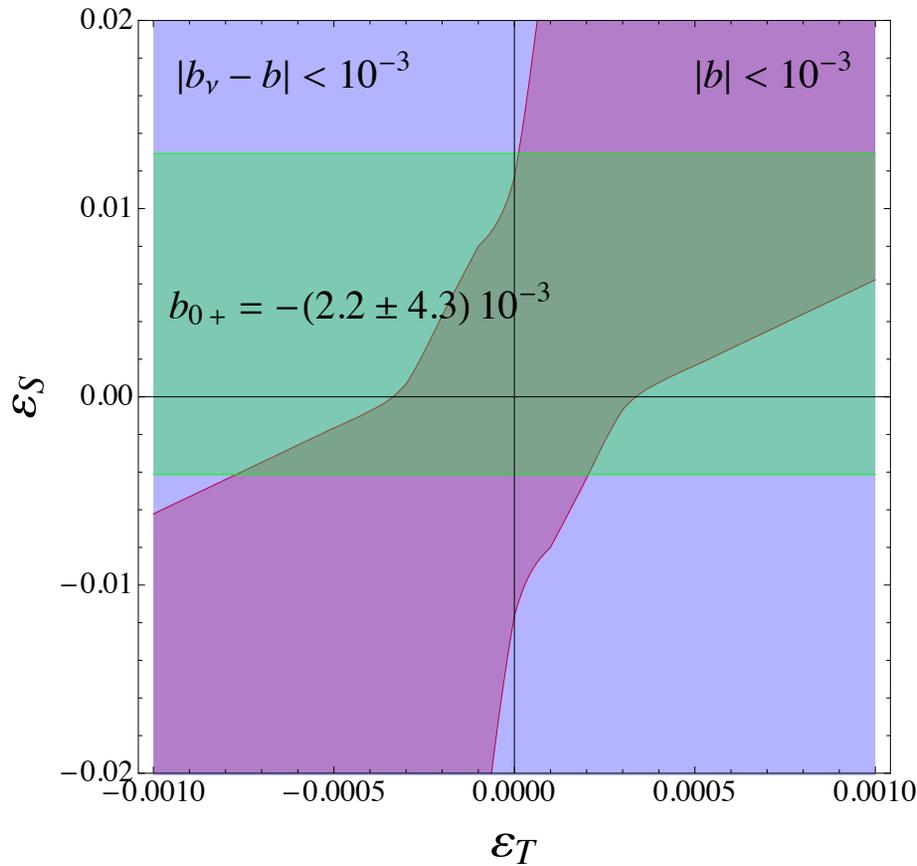
- Charges (g_A, g_S, g_T)
- Vector and axial form factor
- Generalized Parton Distribution functions

Low energy constraints

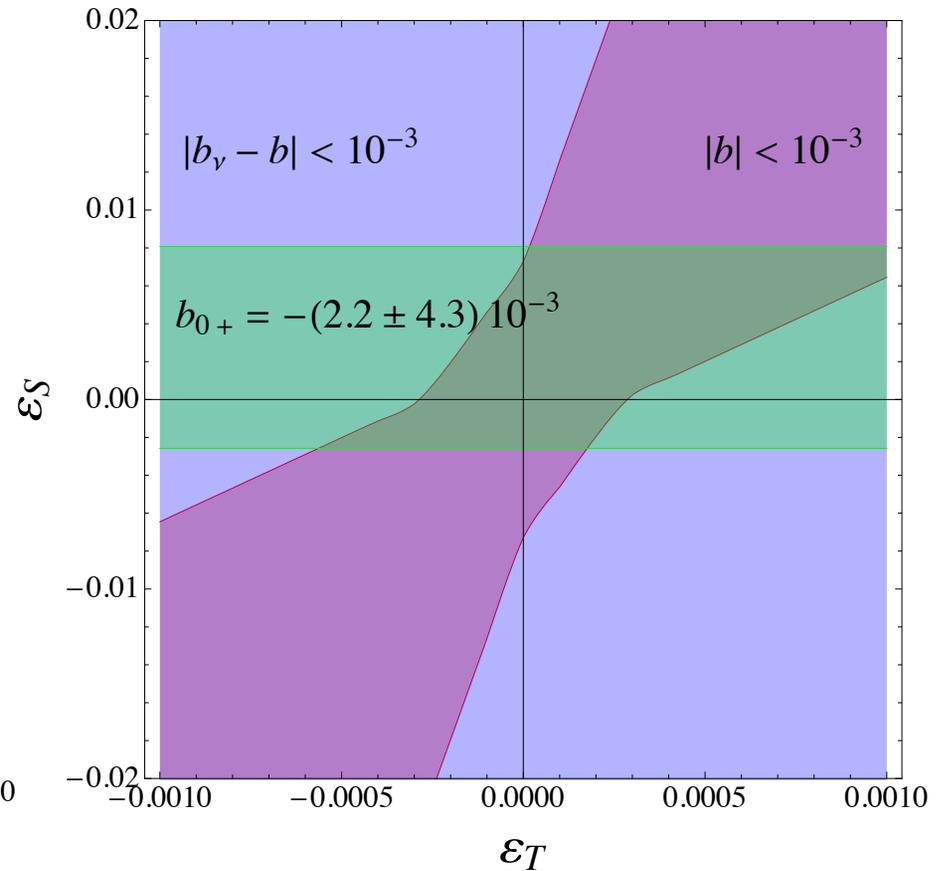
- **Current:** $0^+ \rightarrow 0^+$ and $\pi \rightarrow e \nu \gamma$



Allowed region in ε_S and ε_T are being
constrained as estimates of g_S and g_T improve

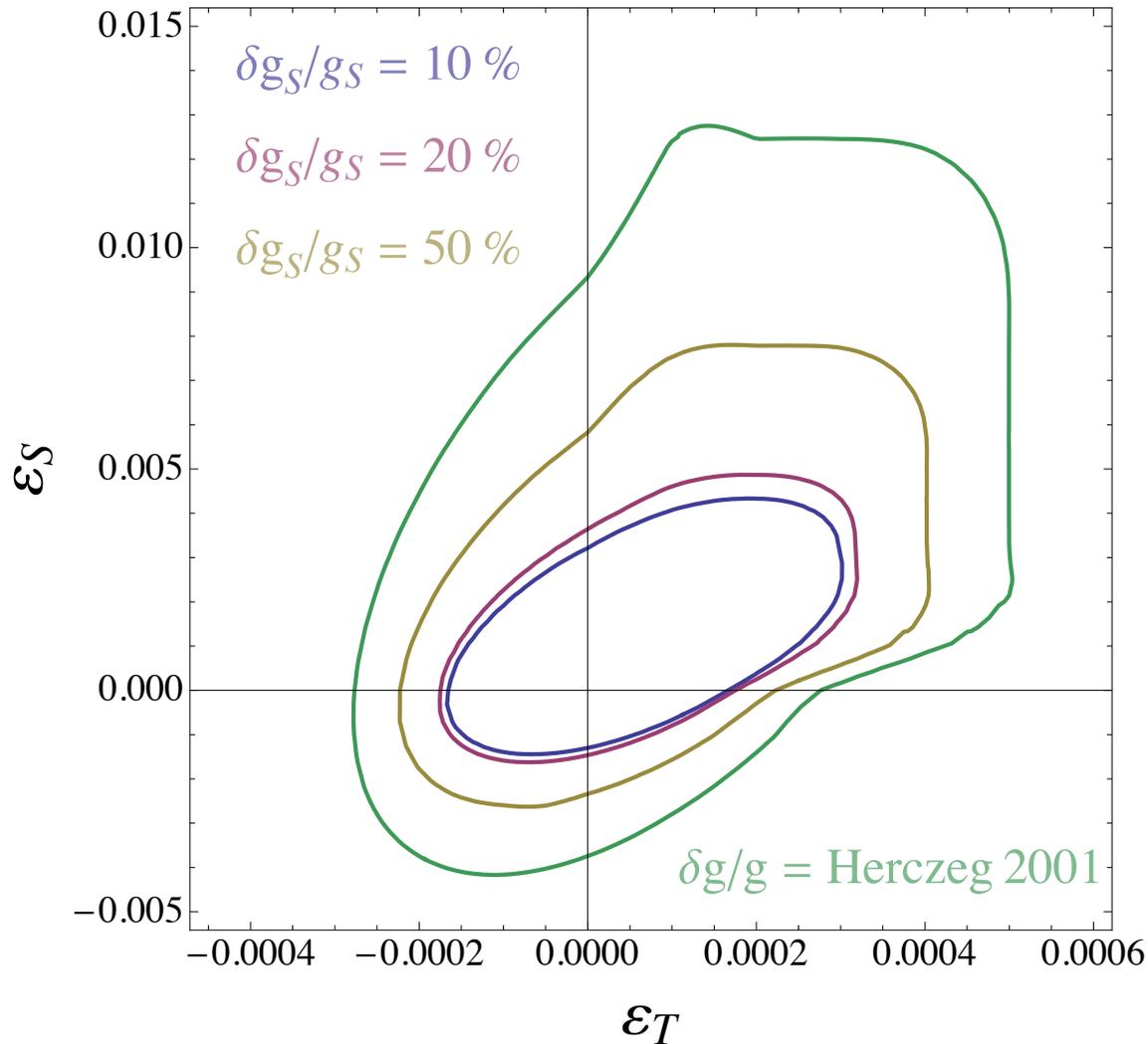


Herczeg-2001
 $0.25 < g_S < 1.0$
 $0.6 < g_T < 2.3$



Lattice-2011
 $g_S = 0.8(4)$
 $g_T = 1.05(35)$

Target Precision for g_S, g_T : 10-20%



Expt. input

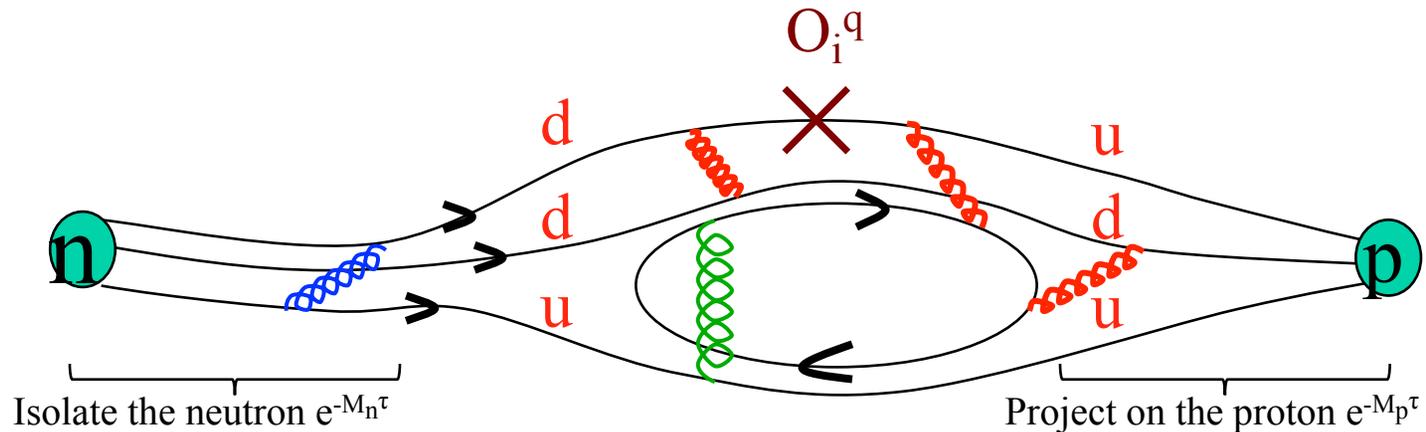
$$|B_1 - b| < 10^{-3}$$

$$|b| < 10^{-3}$$

$$b_{0+} = 2.2 (4.3) * 10^{-3}$$

Allowed region in $[\varepsilon_S, \varepsilon_T]$ (90% contours)

Precision Lattice QCD calculations: $\langle p | \bar{u} \Gamma d | n \rangle$

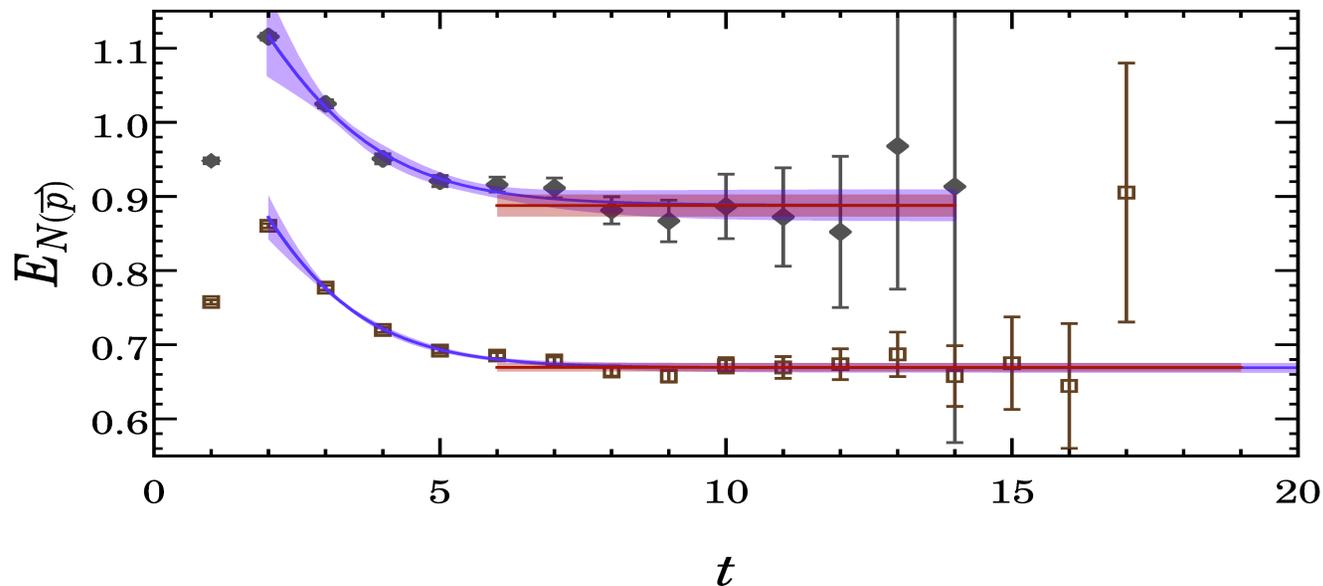


- **Achieving 10-20% uncertainty is a realistic goal but requires:**
 - High Statistics ($O(15000)$ measurements)
 - Controlling all Systematic Errors:
 - Contamination from excited states
 - Non-perturbative renormalization of bilinears (RI_{smom} scheme)
 - Finite volume effects
 - Chiral extrapolation to physical m_u and m_d (simulate at physical point)
 - Extrapolation to the continuum limit (lattice spacing $a \rightarrow 0$)

Reducing excited state contamination

2-point correlation function $\rightarrow M_N$

$$\Gamma^2(t_f, t_i) = A_0 e^{-M_0 \Delta t} + A_1 e^{-M_1 \Delta t} + \dots$$



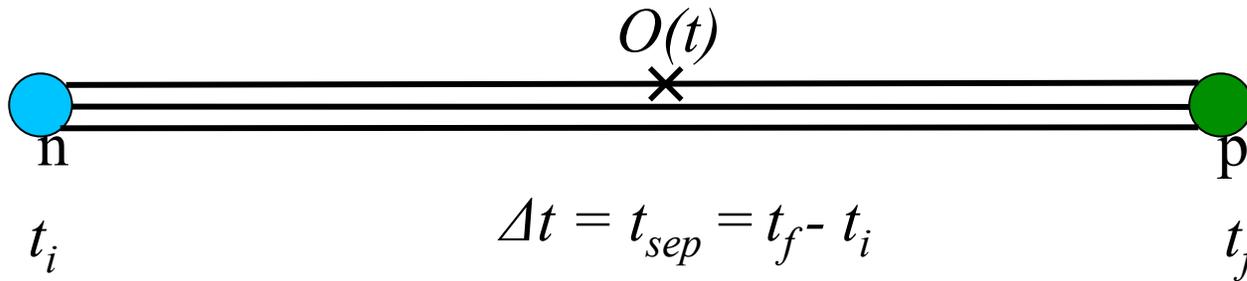
Current data are fit by including 1 “excited” state

Reducing excited state contamination: 3-pt fn.

Assuming 1 excited state, the 3-point function is given by

$$\Gamma^3(t_f, t, t_i) = |A_0|^2 \langle 0|O|0\rangle e^{-M_0 \Delta t} + |A_1|^2 \langle 1|O|1\rangle e^{-M_1 \Delta t} + \\ A_0 A_1^* \langle 0|O|1\rangle e^{-M_0 \Delta t} e^{-M_1(\Delta t - t)} + A_0^* A_1 \langle 1|O|0\rangle e^{-M_1 \Delta t} e^{-M_0(\Delta t - t)}$$

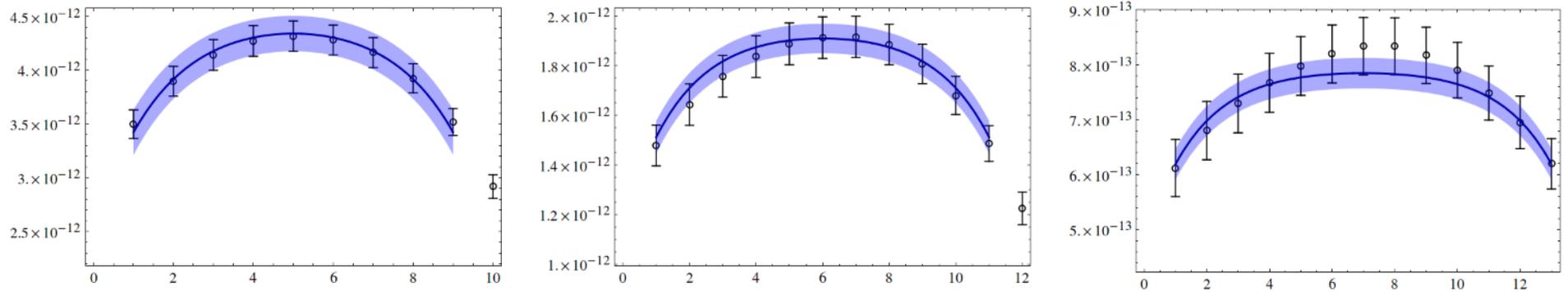
Where M_0 and M_1 are the masses of the ground & excited state and A_0 and A_1 are the corresponding amplitudes.



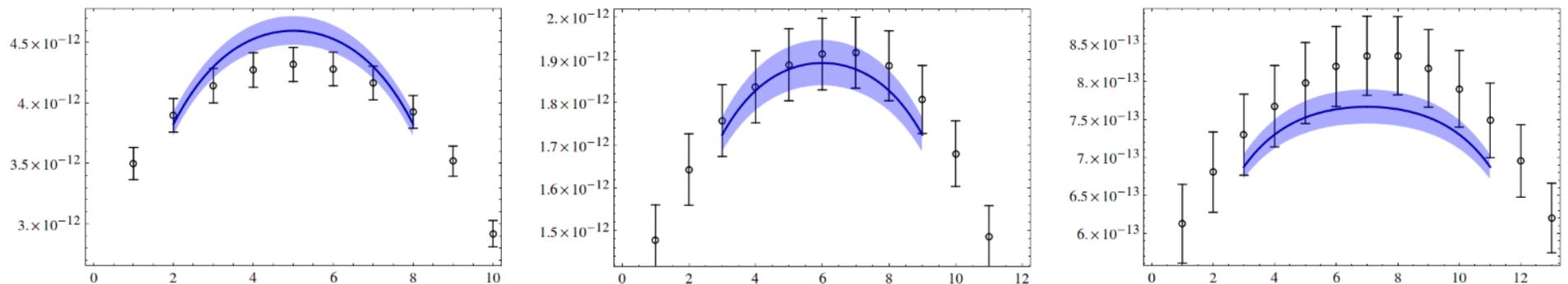
Need simultaneous fit to data at multiple $\Delta t = t_{sep} = t_f - t_i$

Simultaneous fit to multiple t_{sep}

Data for g_S on the $M_\pi=220$ MeV ensemble at $a=0.09$ fm



Excluding $\langle 1|O|1\rangle$ Term



$t_{sep}=10$

$t_{sep}=12$

$t_{sep}=14$

Renormalization of bilinear operators

- Non-perturbative renormalization factors Z_Γ using the RI-sMOM scheme ($p_1^2 = p_2^2 = q^2$)
 - Need quark propagator in momentum space

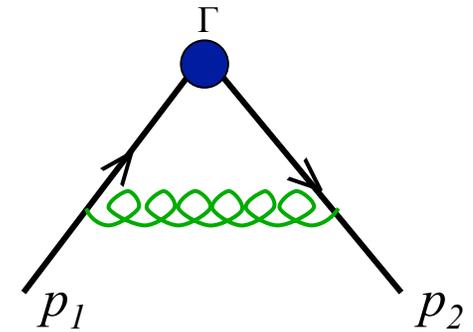
- Basic Assumption: there exists a window

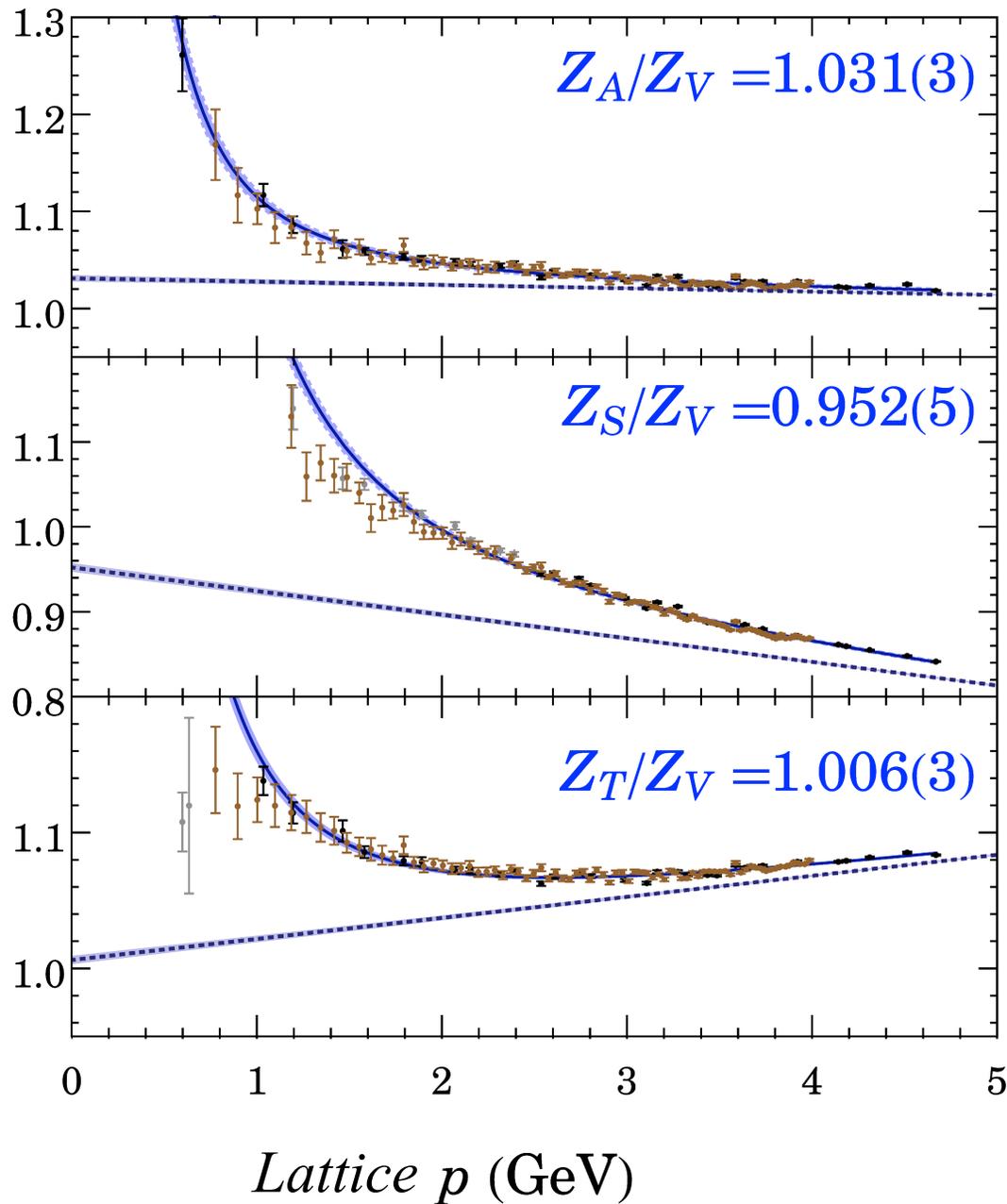
$$\Lambda_{QCD} \ll p \ll \pi/a$$

- HYP Smearing introduces artifacts

- Gluon momentum above ($\sim 1/a$) are averaged out
- $\Lambda_{QCD} \ll p \ll \pi/a$ window may not exist on coarse lattices

- No detectable dependence of Z 's on m_q





$$\frac{Z_{A,S,T}}{Z_V} \left(\overline{MS}, 2 \text{ GeV} \right)$$

*Fit data to: $A/p + Z + Cp$
in the range $\{1 < p < 4 \text{ GeV}\}$
Or
Choose Z at $p^2 = 5 \text{ GeV}^2$
& errors from $\{4 < p^2 < 6 \text{ GeV}^2\}$*

Renormalized Charges

$$\frac{Z_{A,S,T}}{Z_V} \times \frac{g_{A,S,T}}{g_V}$$

Ward identity: $Z_V g_V = 1$

Observations and Lessons Learned

- For given statistics: $\sigma(g_S) \sim 5 \sigma(g_A)$ [or $5 \sigma(g_T)$]
Need $O(15000)$ independent measurement (Configs \times Sources)
- Excited state contamination is significant but controlled
Need: data at multiple t_{sep} with good signal for $t_{sep} > 1.2\text{fm}$
fits including at least one excited state to data $t_{sep} > 1.0\text{fm}$
- Renormalization (RI-sMOM): Smearing introduces artifacts
Impose a prescription with a well-defined continuum limit

Analyzing lattice data: Extrapolations in a, M_π^2, L

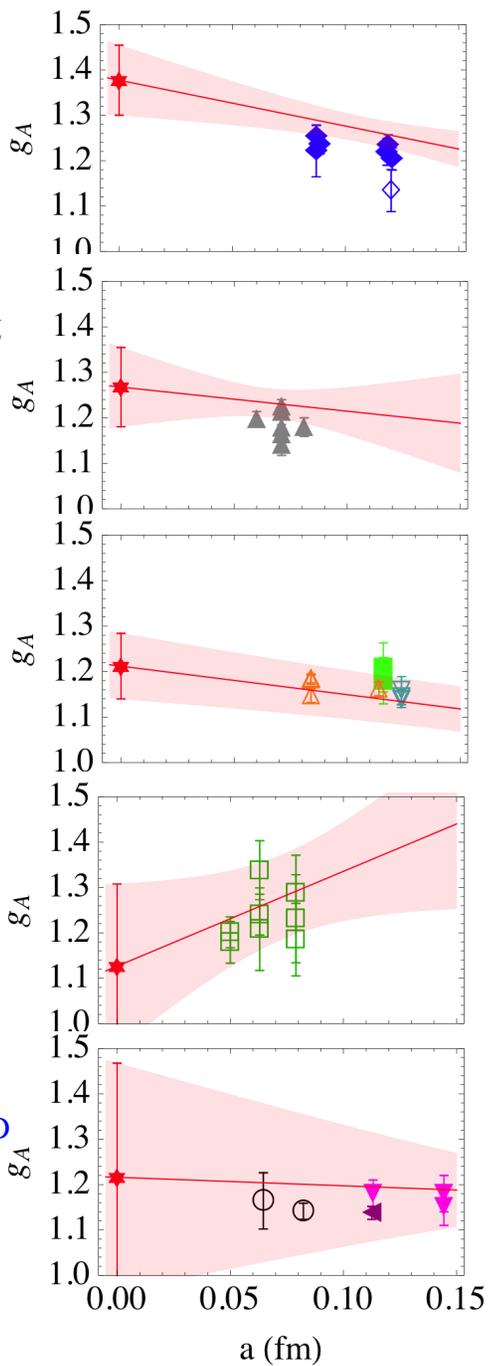
Using lowest order corrections when fitting lattice data w.r.t.

- Lattice spacing a
- Dependence on quark mass $m_q \sim M_\pi^2$
- Finite volume $M_\pi L$

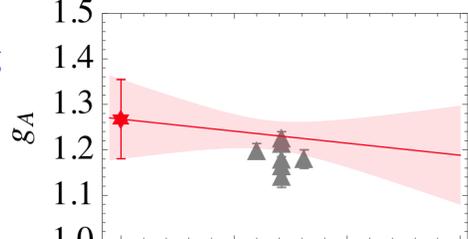
$$g(a, M_\pi, L) = g + A a + B M_\pi^2 + C e^{-M_\pi L} + \dots$$

Lattice QCD Calculations
are ongoing and collaborations
are addressing all sources of
systematic errors

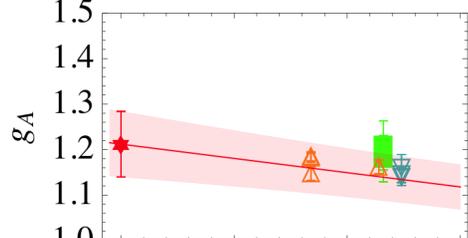
PNDME



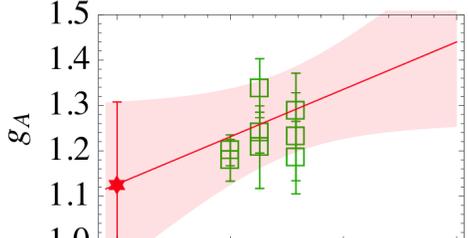
Regensburg



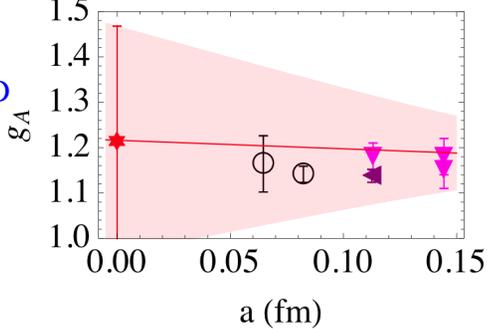
LHPC



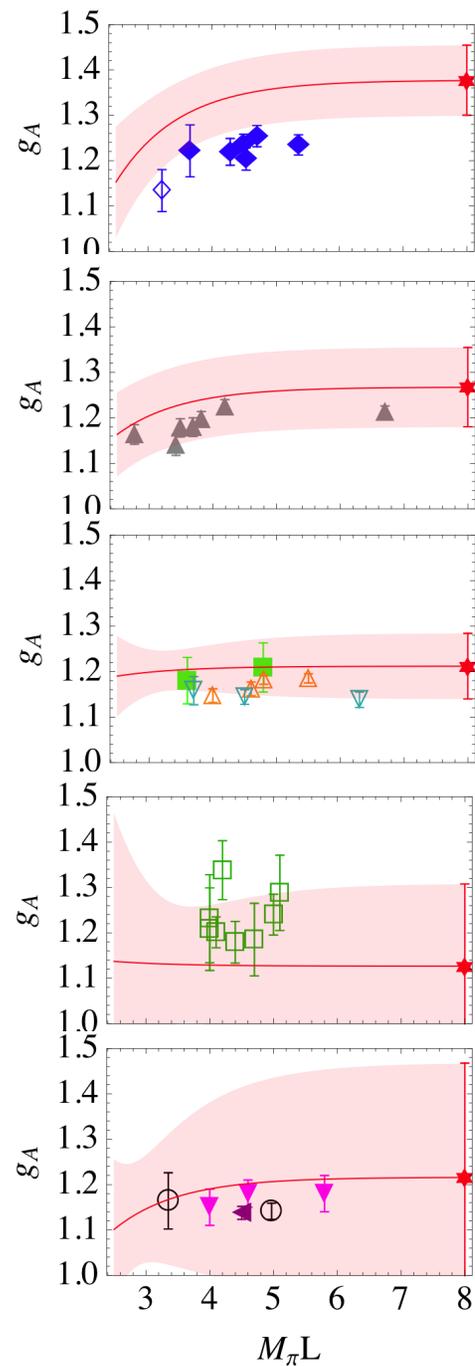
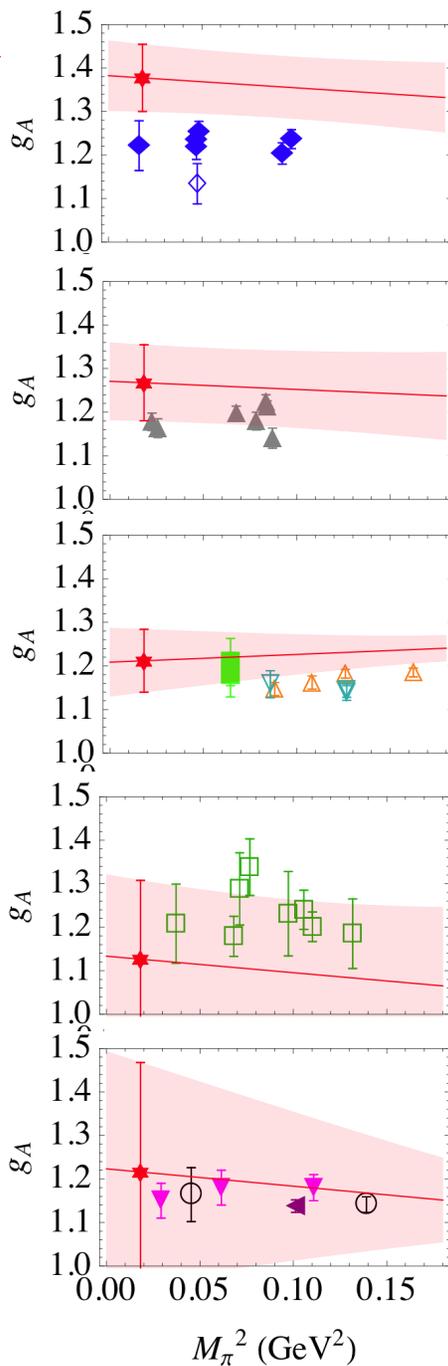
Mainz



ETMC
RBC/UKQCD
 χ QCD

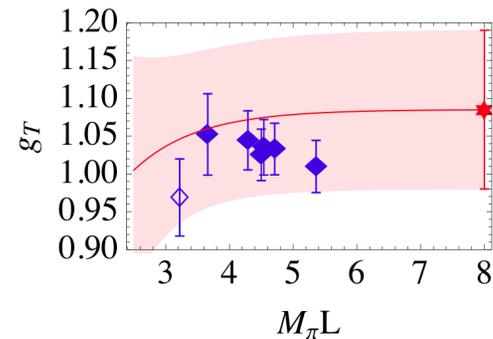
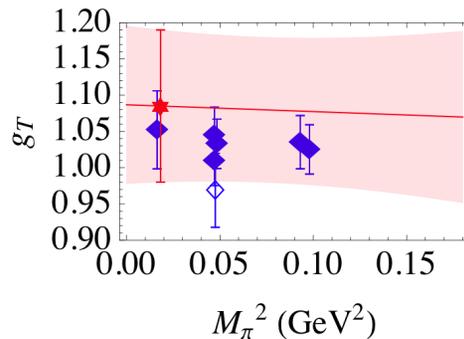
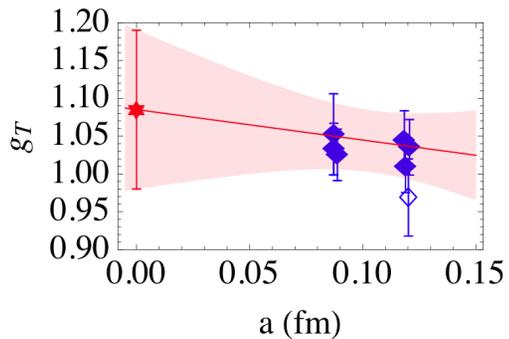


g_A

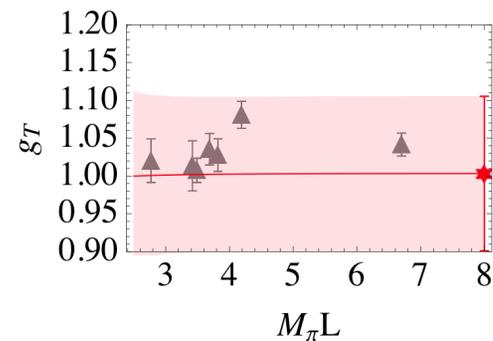
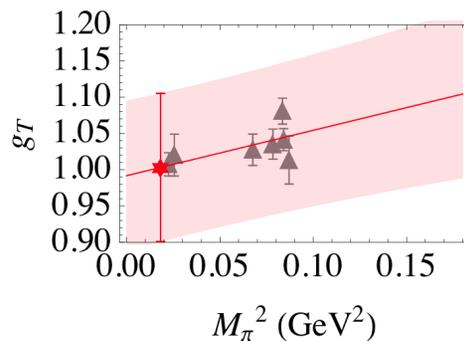
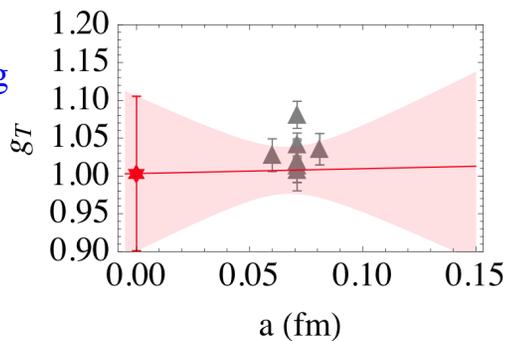


g_T

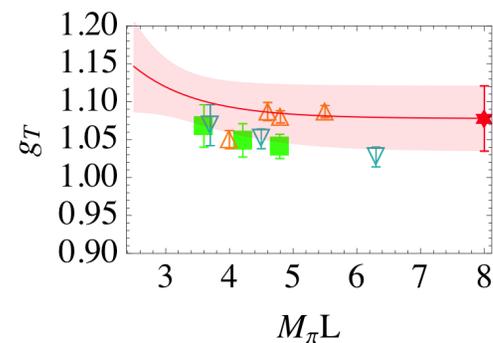
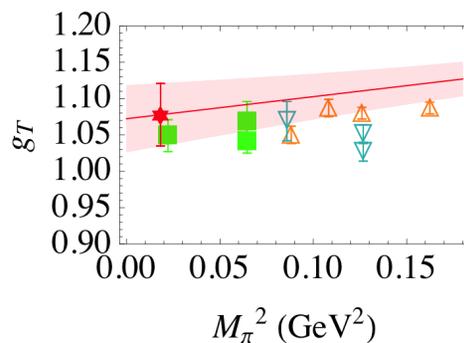
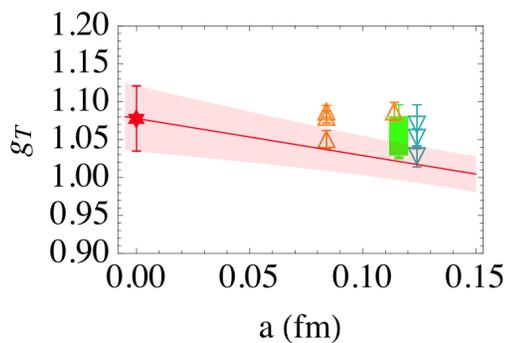
PNDME



Regensburg

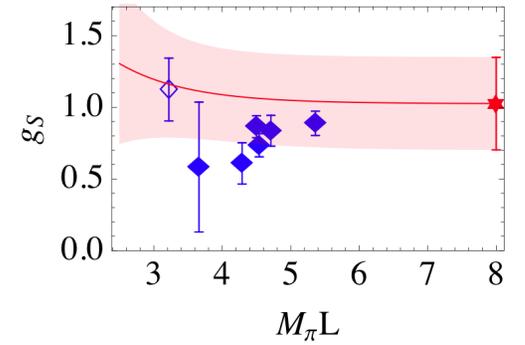
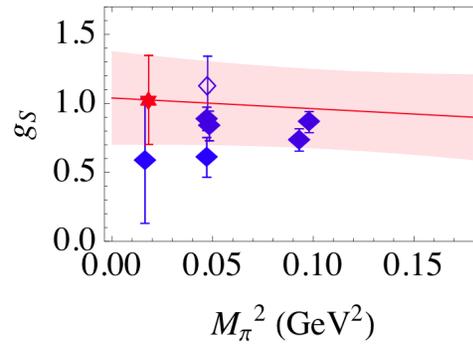
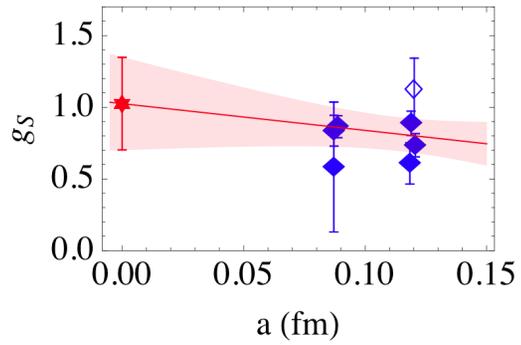


LHPC

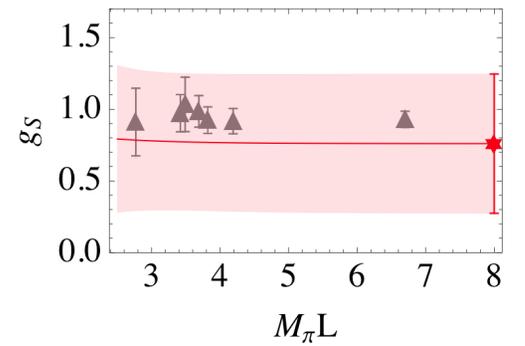
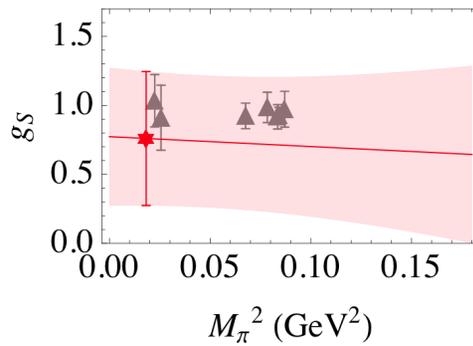
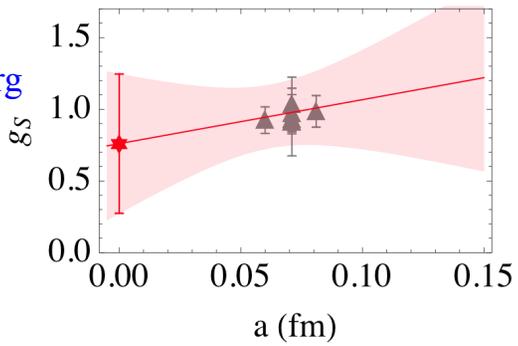


g_S

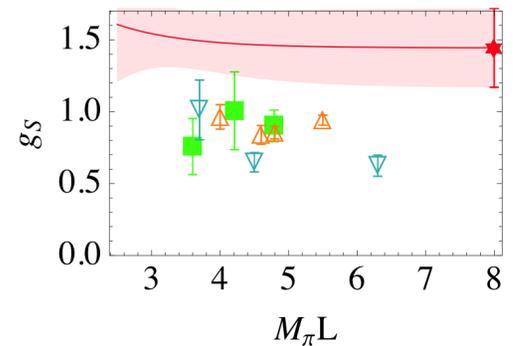
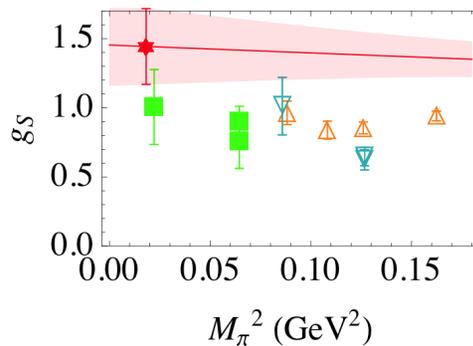
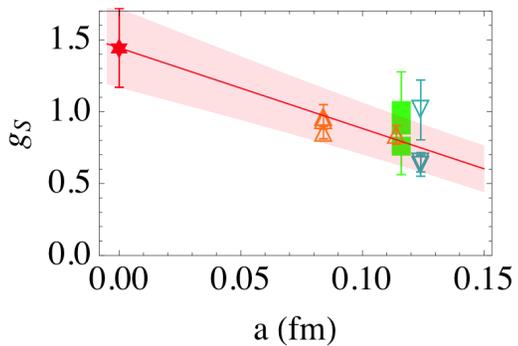
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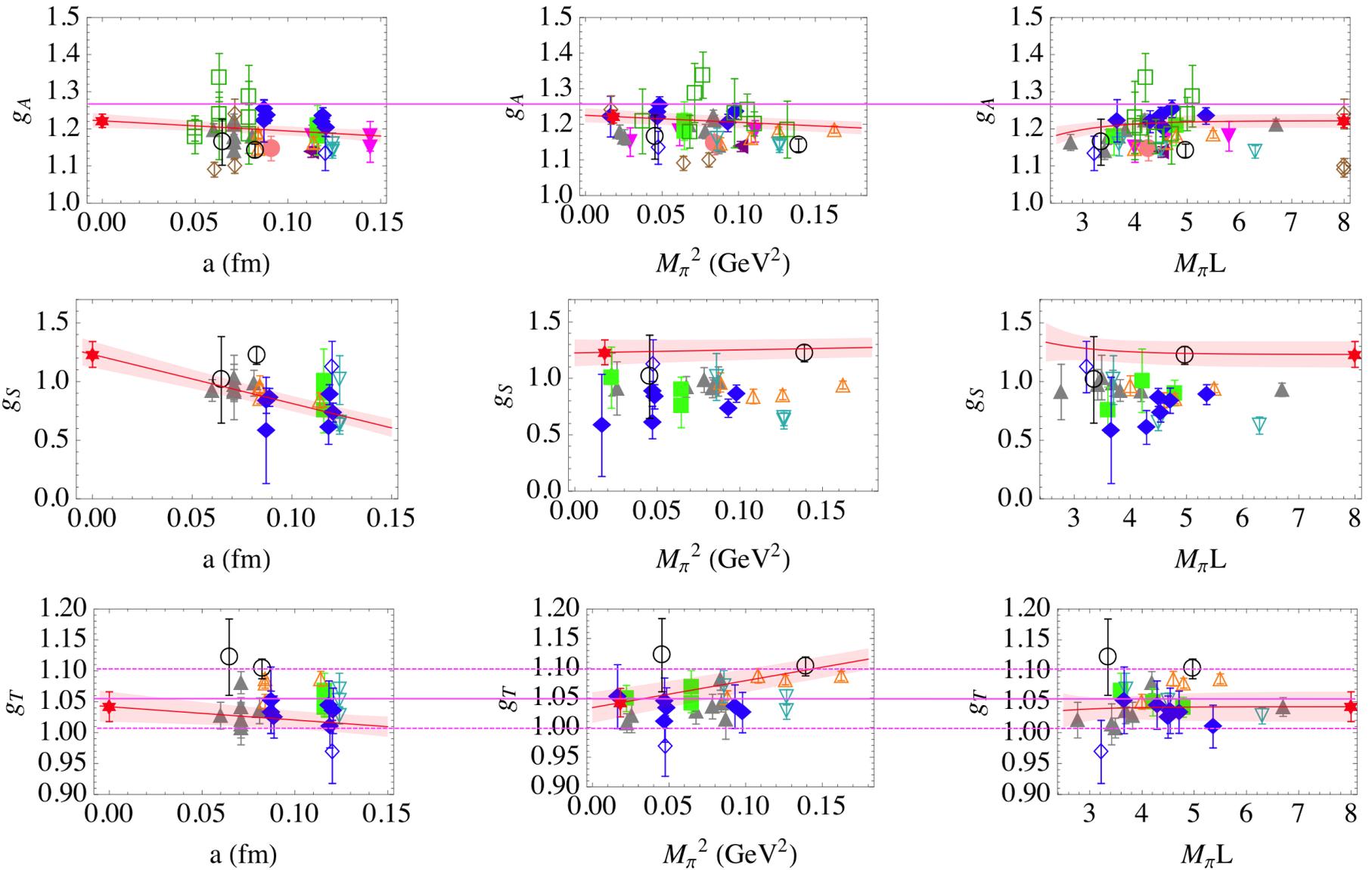
Regensburg



LHPC



lattice data: combined plots



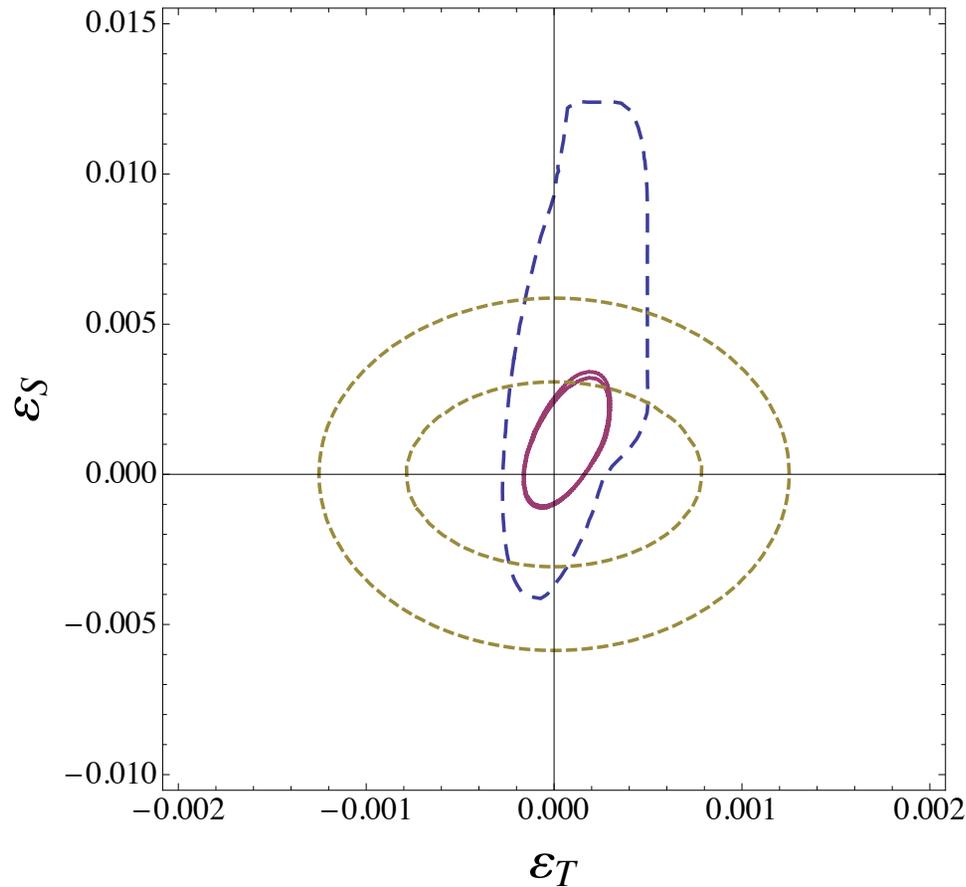
Towards Physical Estimates

- g_T : 1.05(5)
 - Estimate of Z_T is reliable
 - Small dependence on a , M_π^2 , $M_\pi L$ (Caution: see g_A)
- g_A :
 - Estimate of Z_A is reliable
 - Extrapolations in a , M_π^2 , $M_\pi L$ are not yet resolved
- g_S :
 - Statistical errors are large
 - Z_S is not well-determined
 - Extrapolations in a , M_π^2 , $M_\pi L$ are not stable

Constraining ϵ_S and ϵ_T

- Measurements of b and B_I at 10^{-3} to 10^{-4} will probe multi-TeV scale and place stringent constraints on novel scalar and tensor interactions
 - $G_F \epsilon_{S,T} = (1/\Lambda_{S,T})^2$
 - $\epsilon_{S,T} = v^2/\Lambda_{S,T}^2 \sim 10^{-3}$
- $\Lambda_{S,T} \sim 5 \text{ TeV}$
- Constraints on ϵ_S and ϵ_T from [U]CN experiments combined with improved g_S and g_T will be more stringent than existing probes ($0^+ \rightarrow 0^+$; $\pi \rightarrow e\nu\gamma$).
 - Collider experiments are not competitive until $\sqrt{s}=14 \text{ TeV}$ & 100 fb^{-1}

Constraints from β -decay versus LHC



LHC bounds:

- 10 fb^{-1} at 14TeV
- 300 fb^{-1} at 14TeV

Nuclear exp + Theory bounds:

- $b = B_1 = 10^{-3}$
 - g_S and g_T from Herczeg
 - g_S and g_T from Lattice QCD

Improving low energy bounds further depend on improvements in nuclear experiments for b and B_1