

# The Entropy of Hawking Radiation

Netta Engelhardt

MIT

28th Solvay Conference on Physics

# THE ROLE OF QUANTUM INFORMATION IN QUANTUM GRAVITY

- ▶ Connection between QI and quantum gravity: QI appears to be an important ingredient in gravitational physics.
- ▶ E.g. the AdS/CFT correspondence: von Neumann entropy, complexity, quantum error correction, quantum secret sharing, etc. play an important role in the description of space and time in quantum gravity.
- ▶ Very useful in the context of the black hole information paradox and Hawking radiation.

# HAWKING RADIATION AND EVAPORATING BLACK HOLES

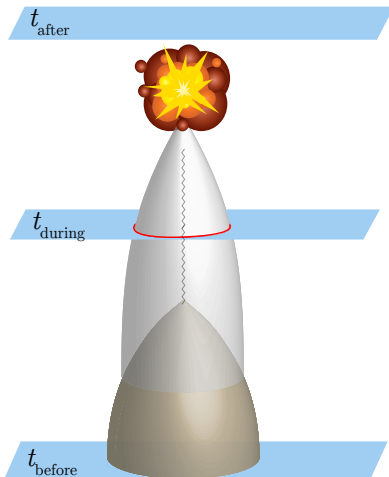
## Black Hole Thermodynamics

Black holes are quantum mechanical systems; in particular they are thermal objects: they have a temperature, and they radiate.

$$T_{BH} \propto M_{BH}^{-1}$$

If  $T_{BH} > T_{\text{surroundings}}$ , the black hole shrinks as it radiates, eventually evaporating altogether.

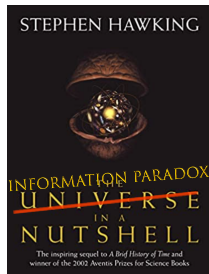
# BLACK HOLE FORMATION AND EVAPORATION



# HAWKING RADIATION IN A NUTSHELL

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Before the black hole forms:

$$S_{\text{vN}}[\rho_{\text{outside BH}}] = S_{\text{vN}}[\rho_{\text{universe}}] = 0$$



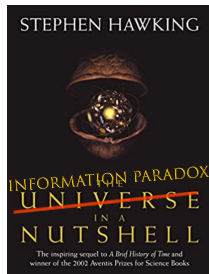
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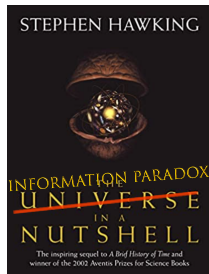
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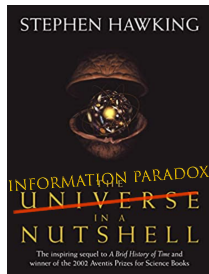
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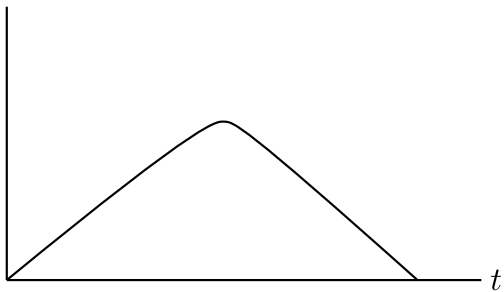
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4. So we expect the Page curve for  $S_{\text{vN}}[\rho_{\text{rad}}]$ .



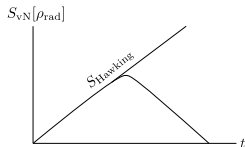


$S_{\text{vN}}[\rho_{\text{rad}}]$

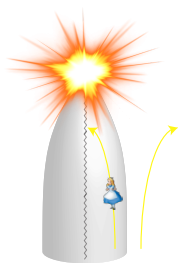


# THE PARADOX

- ▶ To compute  $S_{\text{vN}}[\rho_{\text{rad}}]$ , we need a description of the physics at the black hole event horizon.
- ▶ *Semiclassical gravity*: the limit where large quantum gravity effects are suppressed; should be valid at the horizon.
- ▶ Semiclassical gravity  $\stackrel{\text{Hawking}}{\Rightarrow} -\text{tr} \rho_{\text{rad}} \ln \rho_{\text{rad}}$  does *not* follow the Page curve:



## CARICATURE OF INFORMATION LOSS



Quick argument for info loss:

$$|\psi\rangle = |+\rangle|-\rangle + |-\rangle|+\rangle$$

After evaporation, Alice disappears from the universe – even though the universe is a closed system.

State after evaporation is just  $\rho_{\text{Bob}} = \frac{1}{2}\mathbb{I}$ .

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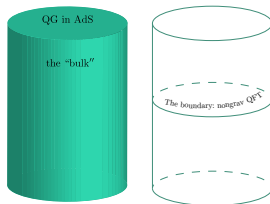
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- ▶ Progress: a semiclassical calculation of the Page curve, but using a different – holographic – formula for  $S_{vN}$ .
- ▶ This formula is a bit of a black box: can be interpreted as indirect input from quantum gravity (or a saddle point analysis of Euclidean quantum gravity)

# HOLOGRAPHY

## Holography

Quantum gravity “in a box” (with Anti-de Sitter boundary conditions) is dual to a lower-dimensional nongravitational QFT.



Colloquially, we call the gravitational theory the bulk, and the non-gravitational theory the boundary.



# HOLOGRAPHY: A BLACK HOLE IS JUST A QUANTUM SYSTEM

This means that a black hole in AdS is just an ordinary, nongravitational quantum system.

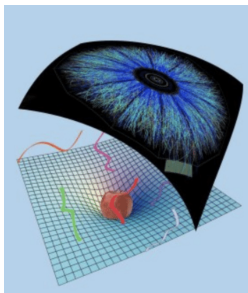


image credit: ESI Programme on AdS Holography and the Quark-Gluon Plasma

... nongravitational (closed) quantum systems evolve unitarily.

## COMPUTING THE PAGE CURVE

We know that  $S_{\text{vN}}[\rho_{\text{rad}}]$  computed in the nongravitational “boundary” theory follows the Page curve.

AdS/CFT is a dictionary: it relates quantities in the boundary theory to the bulk theory.

What quantity in the bulk theory translates to  $S_{\text{vN}}[\rho_{\text{rad}}]$  as computed by the boundary theory?

# HOLOGRAPHIC VON NEUMANN ENTROPY

## $S_{\text{vN}}$ in AdS/CFT

Ryu-Takayanagi, Hubeny-Rangamani-Takayanagi, Faulkner-Lewkowycz-Maldacena, Engelhardt-Wall

The von Neumann entropy of the radiation is given by:

$$S_{\text{vN}}[\rho_{\text{rad}}] = \frac{\text{Area}[\chi]}{4} + S_{\text{vN}}[\rho_{\text{Out}}[\chi]]$$

where  $\chi$  is a “quantum extremal surface”: if you slightly perturb  $\chi$ , the sum  $\frac{\text{Area}[\chi]}{4} + S_{\text{vN}}[\rho_{\text{Out}}[\chi]]$  won't change to leading order in the perturbation.

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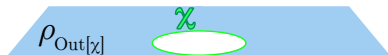
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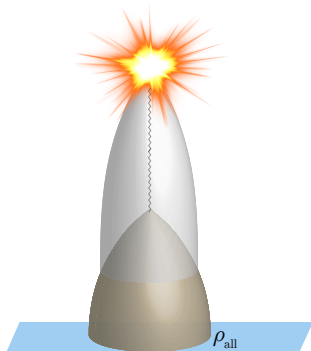
We get  $\rho_{\text{out}}[\chi]$  from a bipartition of the system into interior of  $\chi$  and exterior.



# QUANTUM EXTREMAL SURFACES

Initially, the QES is the empty set.

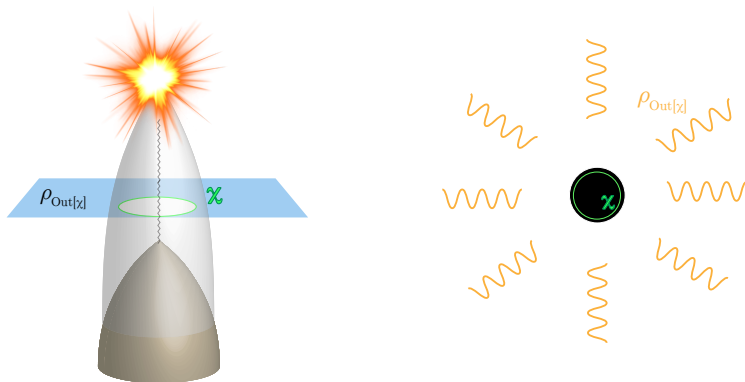
$$S_{vN}[\rho_{bdy}] = \frac{\text{Area}[\emptyset]}{4G\hbar} + S_{vN}[\rho_{all}] = S_{vN}[\rho_{all}]$$



With time  $S_{vN}[\rho_{all}]$  grows.

# QESs IN EVAPORATING BHs ALMHEIRI, ENGELHARDT, MAROLF, MAXFIELD; PENINGTON

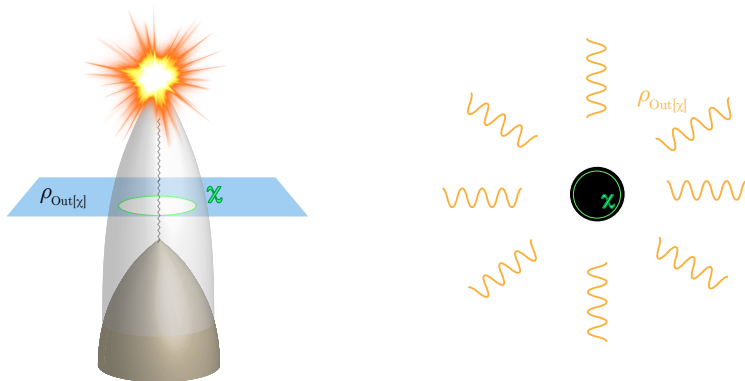
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This alternative computation of  $S_{\text{vN}}$  in gravity *does* yield a unitary Page curve.

# THE ENTANGLEMENT WEDGE

What is the significance of the region  $\text{Out}[\chi]$ ?

## Reconstruction in AdS/CFT?

As Douglas discussed, AdS/CFT is a quantum error correcting code see work beginning with Almheiri, Dong, Harlow. How much of the bulk can we reconstruct from the state  $\rho_{\text{bdy}}$  of the nongravitational dual?



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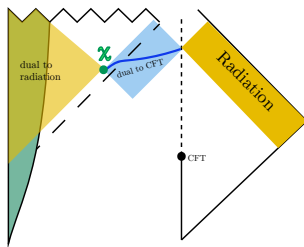
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Everything in  $\text{Out}[\chi]$ .\*

# THE ENTANGLEMENT WEDGE OF THE RADIATION

After the Page time, this is most of the black hole interior: information that fell into the black hole is now available in the radiation.



## DEPENDENCE ON THE CODE SUBSPACE

Recall: “Everything in  $\text{Out}[\chi]^*$ ”

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## The Code Subspace

Less naively: that depends on the code subspace. Given a choice of code subspace, we can reconstruct everything in  $\text{Out}[\chi_M]$ , where  $\chi_M$  is the QES for the maximally mixed state in our code subspace Hayden, Penington.

The larger your code subspace, the less of the bulk you can reconstruct!

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$$V_{EL} |\Psi\rangle_{B'EL} = |i\rangle_1 \otimes |\phi\rangle_{B',2}$$

where  $E$  is the early radiation,  $L$  is the late radiation (after  $|i\rangle$  is thrown in) and  $B'$  is the black hole after  $L$  has radiated (and 1 and 2 are factors of  $\mathcal{H}_{LE}$ ). But:



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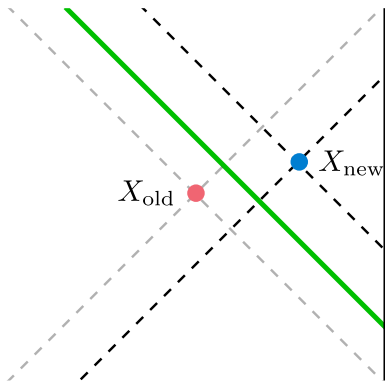
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- ▶ **Harlow-Hayden:** decoding the radiation is exponentially complex in the entropy of the black hole. **Aaronson:** decoding the radiation is polynomially hard  $\Rightarrow$  there are no injective one-way functions.

# HOLOGRAPHIC REALIZATION OF HAYDEN-PRESKILL

If you throw something into an evaporating black hole, it will cross into the entanglement wedge of the radiation (with some delay) as the QES moves outwards.

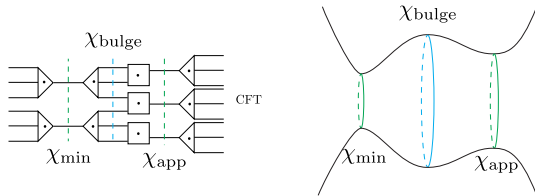


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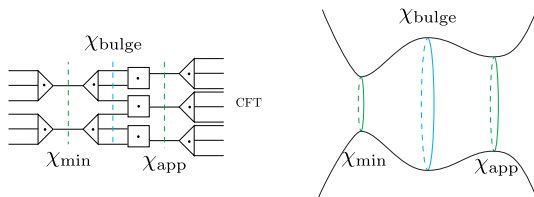
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## Python's Lunch Brown, Gharibyan, Susskind, Penington, Engelhardt, Penington Shahbazi-Moghadam

Reconstruction of a bulk operator is exponentially hard if and only if it lies behind a nonminimal quantum extremal surface. In that case the complexity of reconstruction is proportional to:

$$\text{Exp} \left[ \frac{1}{2} \left( \frac{\text{Area}[\chi_{\text{bulge}}]}{4} + S[\rho_{\text{Out}}[\chi_{\text{bulge}}]] - \frac{\text{Area}[\chi_{\text{app}}]}{4} - S[\rho_{\text{Out}}[\chi_{\text{app}}]] \right) \right]$$